

Dynamics of one-dimensional Bose liquids in Y-junction and its related system: Andreev-like reflection and absence of the Aharonov-Bohm effect

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- 1. Introduction:
One-dimensional Bose liquid
and Tomonaga-Luttinger liquid**
- 2. Y-junction systems for Bose liquid**
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INTRODUCTION

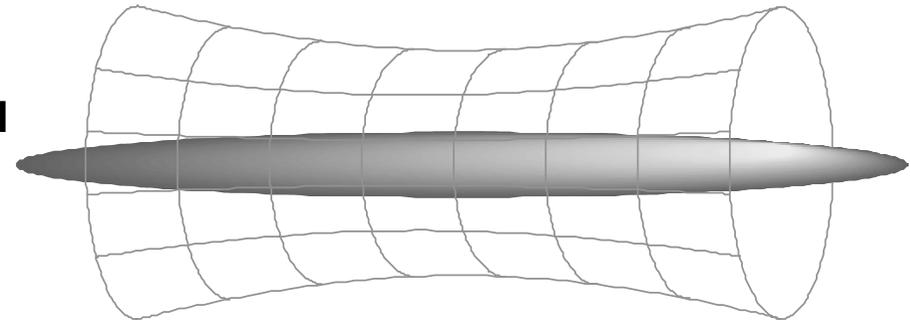
1D bose system

Repulsively interacting bosons

Tight confinement by cylindrical trap potential



Effective 1D shaped systems



Effective 1D Hamiltonian

$$\mathcal{H} = \int dx \left[\frac{\hbar^2}{2m} \partial_x \psi^\dagger(x) \partial_x \psi(x) + \frac{U}{2} \rho(x) \rho(x) \right]$$

Effective 1D interaction for confined systems

$$U = \frac{u}{\pi a_\perp} \left(1 - C \frac{a}{a_\perp} \right)^{-1}$$

M. Olshanii (1998)



Lieb-Liniger model (integrable)

E. H. Lieb, W. Liniger (1963), E. H. Leib (1963)

INTRODUCTION

Low-energy effective theory and Tomonaga-Luttinger liquid for Bose systems

TL liquid is universal description in low-energy physics of 1D quantum systems.

F. D. M. Haldane (1981)

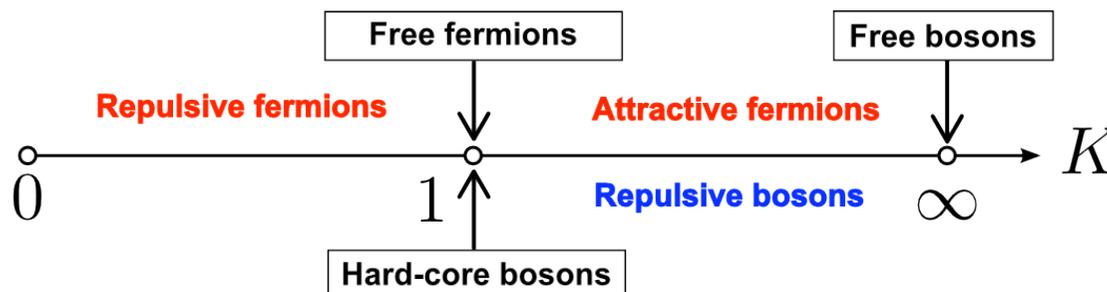
Extract collective excitation

$$\psi(x) \sim \sqrt{\bar{\rho}} e^{i\theta(x)} \quad \rho(x) \sim \bar{\rho} + \frac{1}{\pi} \partial_x \varphi$$

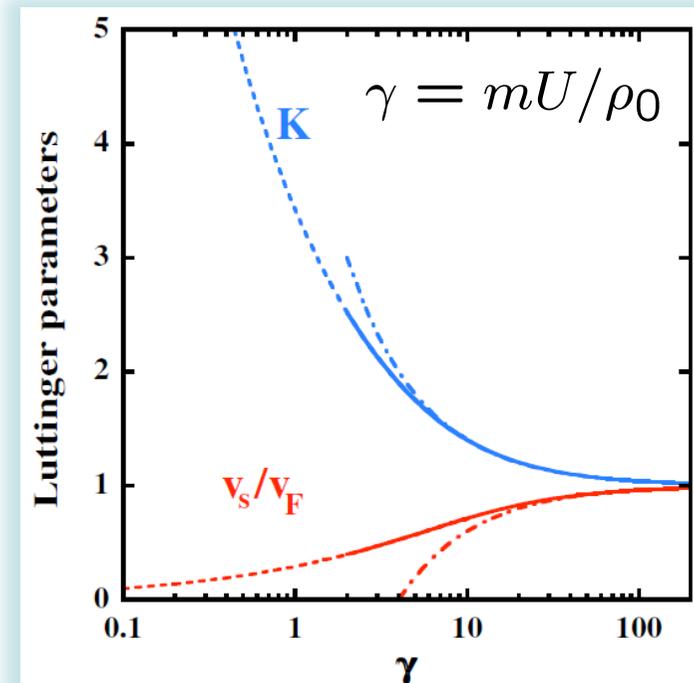
$$\mathcal{H} = \frac{v}{2\pi} \int dx \left[K \left(\frac{\partial \varphi}{\partial x} \right)^2 + \frac{1}{K} \left(\frac{\partial \theta}{\partial x} \right)^2 \right]$$

K : Luttinger parameter

v : velocity of collective mode



1D quantum systems corresponding to TL parameter.



Luttinger parameter for LL model

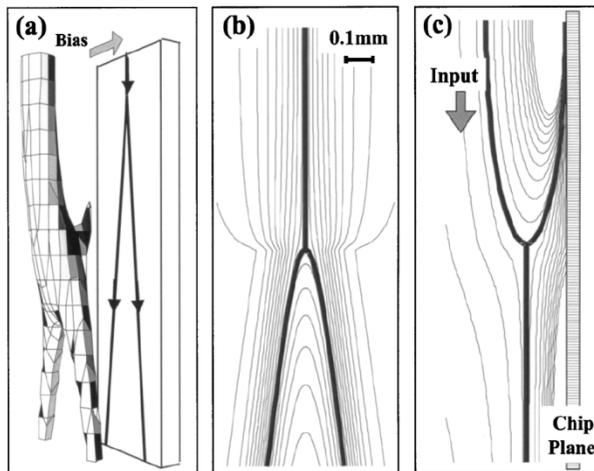
M. A. Cazalilla (2004)

Y-junction system and experiment

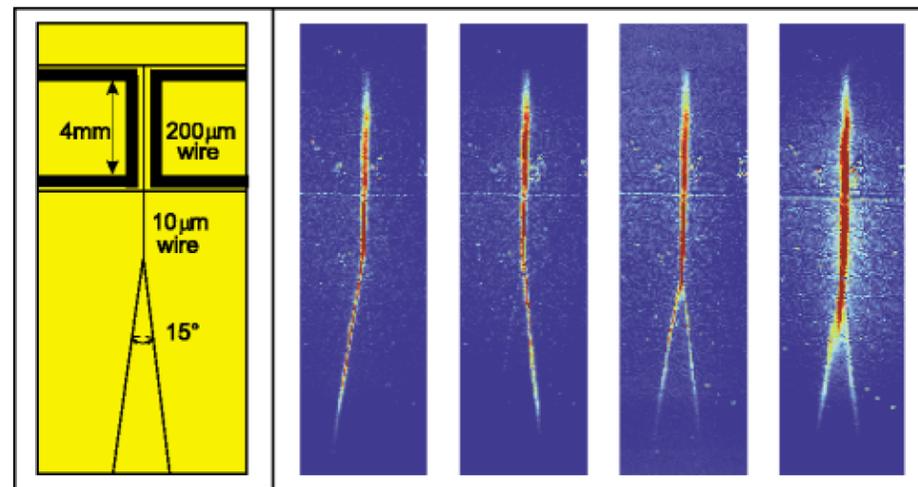
BEC beam splitter for guided trapping atoms.

D. Cassettari et al. (2000)

- Atom chip technology realizes Y-shape trap potential.
- Trapped BEC can be guided as well.
- Application as BEC splitter.



3D imaging and contour map for trap potential.

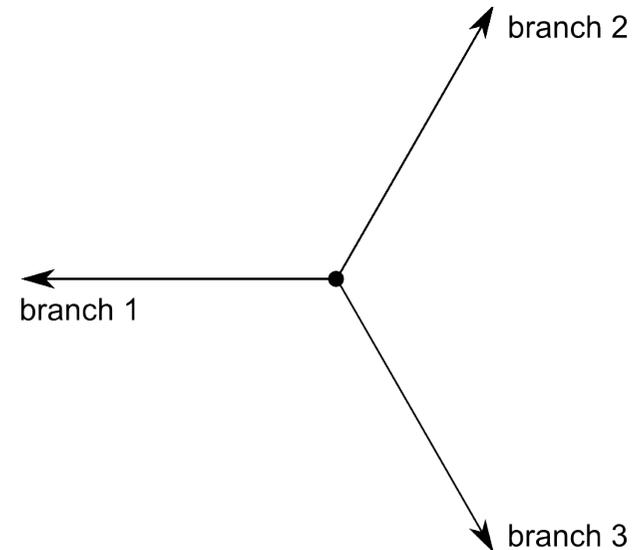
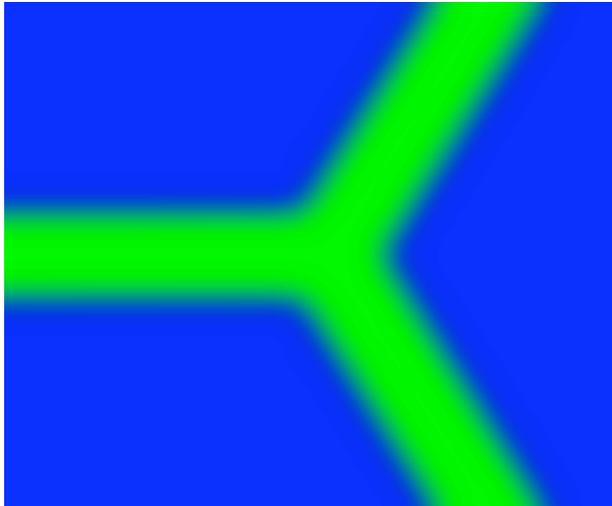


(a) Circuit of Atom chip, and (b) imaging of guiding Y-shaped trapped BEC.

Y-JUNCTION PROBLEM FOR BOSE LIQUID

2008/6/10

Set up the system



Simplify the Y-shaped Bose liquid system.

- Branches are completely 1D.
- Bosons filled in entire system. → TL liquid description is applicable.
- Repulsive interacting bosons. → Luttinger parameter $\mathbf{K} > \mathbf{1}$.

$$\mathcal{H}_j = \int dx \left[-\frac{\hbar^2}{2m} \psi_j^\dagger \frac{\partial^2}{\partial x^2} \psi_j + \frac{U}{2} \rho_j^2 \right]$$

Low-energy physics of Y-junction

$$\mathcal{H} = \mathcal{H}_{\text{bulk}} + \mathcal{H}_{\text{boundary}}$$

◆ Bulk Hamiltonian: TL liquid

$$\mathcal{H}_{\text{bulk}} = \sum_{j=1,2,3} \frac{\hbar v}{2\pi} \int dx \left[K \left(\frac{\partial \theta_j}{\partial x} \right)^2 + \frac{1}{K} \left(\frac{\partial \varphi_j}{\partial x} \right)^2 \right]$$

◆ Boundary Hamiltonian

It is a non-trivial problem to express the concrete form of boundary Hamiltonian.

How to discuss the junction problem.

Transport through a potential barrier. Nayak et al. (1999)

Boundary physics : boundary condition + perturbation

Boundary Condition

Primary boundary condition: current conservation

$$J_1(0, t) + J_2(0, t) + J_3(0, t) = 0$$

What boundary conditions are suitable in low-energy limit?

Y-JUNCTION PROBLEM FOR BOSE LIQUID 2008/6/10

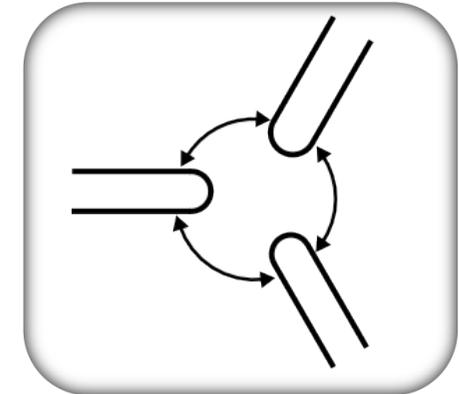
Boundary condition

Simplest boundary condition: decoupled branches

$$J_1(0, t) = J_2(0, t) = J_3(0, t) = 0$$

Perturbation: Tunneling between branches

Scaling dimension = $1/K < 1$ Relevant !!



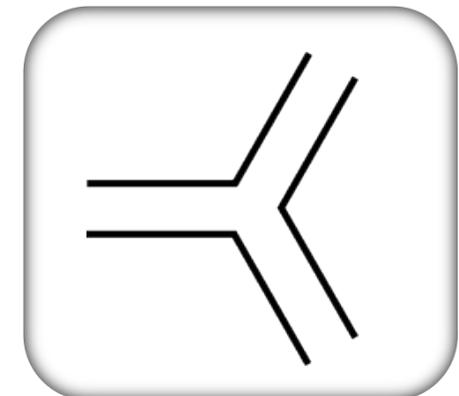
↓
lower energy scale

Arrival point: strongly coupled branches

$$\psi_1(0, t) = \psi_2(0, t) = \psi_3(0, t)$$

Perturbation: Backscattering

Scaling dimension = $4K/3 > 1$ Irrelevant !!



Renormalization group flow goes to the fixed point corresponding to "Strongly coupled limit".

Y-JUNCTION PROBLEM FOR BOSE LIQUID 2008/6/10

Calculation of dynamics

Time evolution

For expectation values of fields \rightarrow classical linear wave.

$$\langle \rho_j(x, t) \rangle = \bar{\rho} + \rho_j^L(x + vt) + \rho_j^R(x - vt).$$

Initial condition

Branch 1

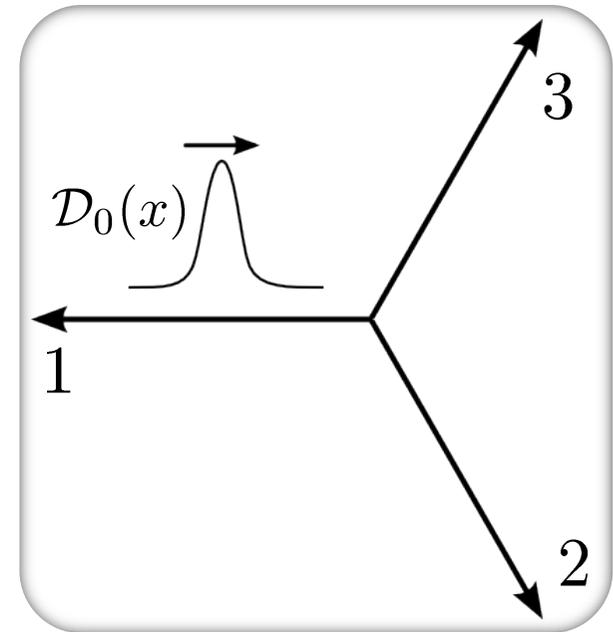
$$\rho_1^L(x, 0) = \mathcal{D}_0(x)$$

$$\rho_1^R(x, 0) = 0$$

Branch 2, 3

$$\rho_{2,3}^L(x, 0) = 0$$

$$\rho_{2,3}^R(x, 0) = 0$$



Boundary condition

Current conservation

$$J_1(0, t) + J_2(0, t) + J_3(0, t) = 0$$

Strongly coupled junction

$$\rho_1(0, t) = \rho_2(0, t) = \rho_3(0, t)$$

Reduced to classical linear wave problem with boundary conditions.

Y-JUNCTION PROBLEM FOR BOSE LIQUID 2008/6/10

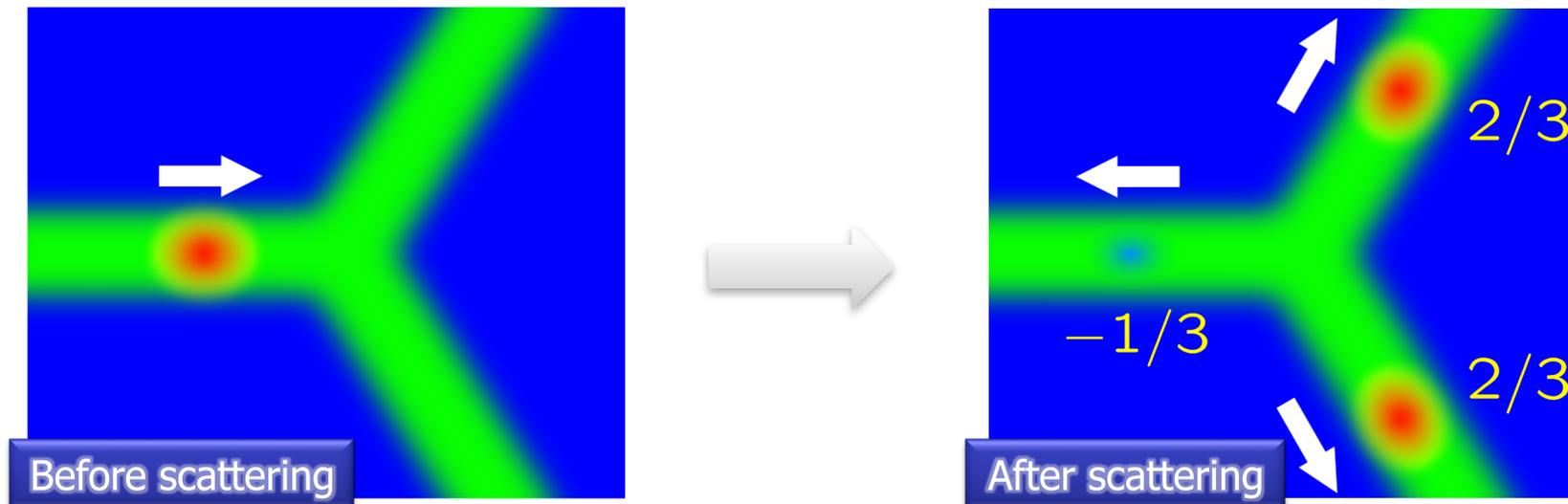
Results

$$\rho_1(x, t) = \rho^0 + \mathcal{D}_0(vt + x) - \frac{1}{3} \mathcal{D}_0(vt - x) \Theta_{\text{step}}(vt - x)$$

Reflection

$$\rho_{2,3}(x, t) = \rho^0 + \frac{2}{3} \mathcal{D}_0(vt - x) \Theta_{\text{step}}(vt - x)$$

Transmission



- Total transmission : **4/3** *Enhanced!*
- Negative reflection : **-1/3** *Hole-like reflection!*

Andreev-like reflection can be also observed in boson systems.

RING-TYPE INTERFEROMETER

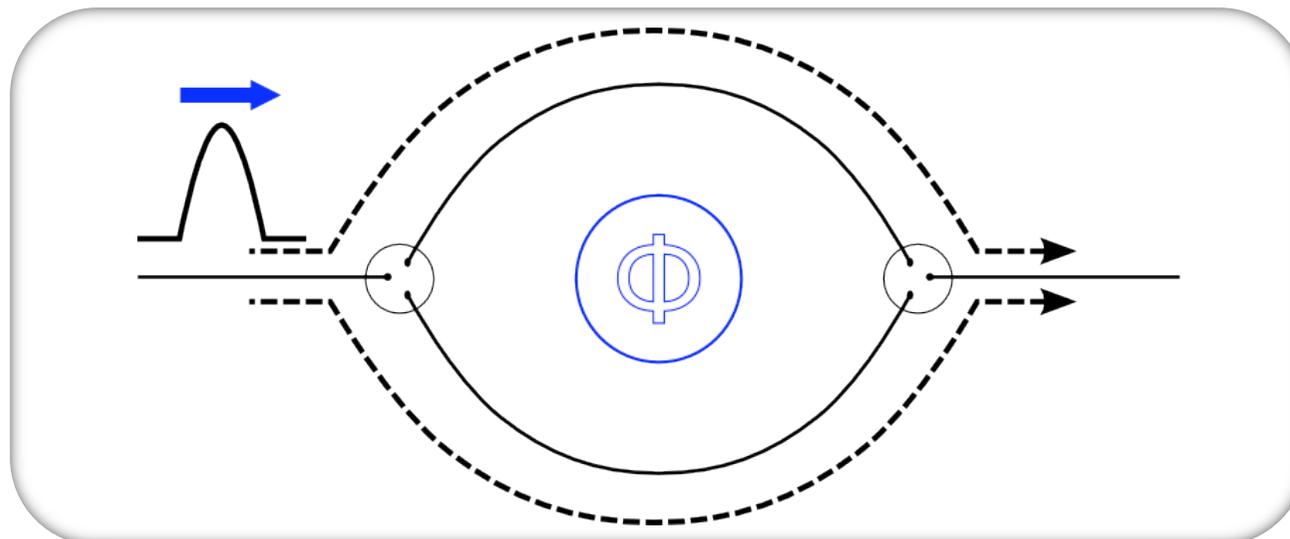
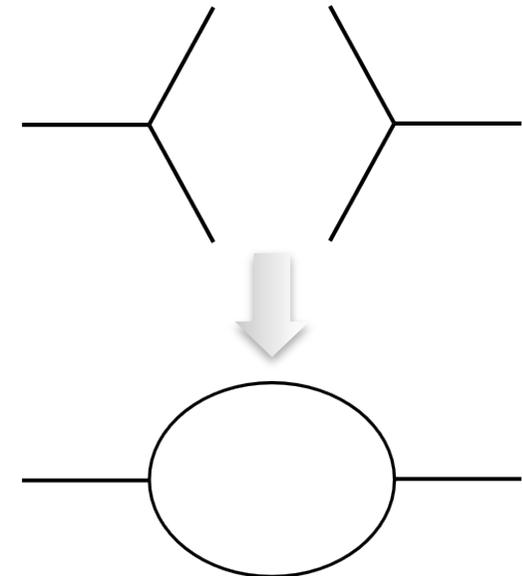
Extension to double Y-junction problem

Symmetrically connect two of branches of each Y-junction.



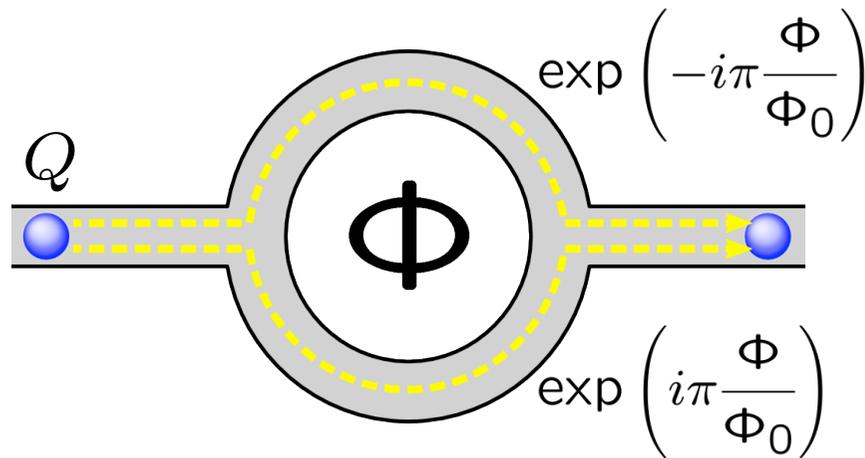
Quantum ring with leads (Ring-type interferometer) system

- "Effective" magnetic flux inside ring.
→ Bosons couple with gauge field.
- Sufficiently large size $2L$ of the ring.
- Bosons are filled in entire system.



RING-TYPE INTERFEROMETER

Example: single particle problem



Probability

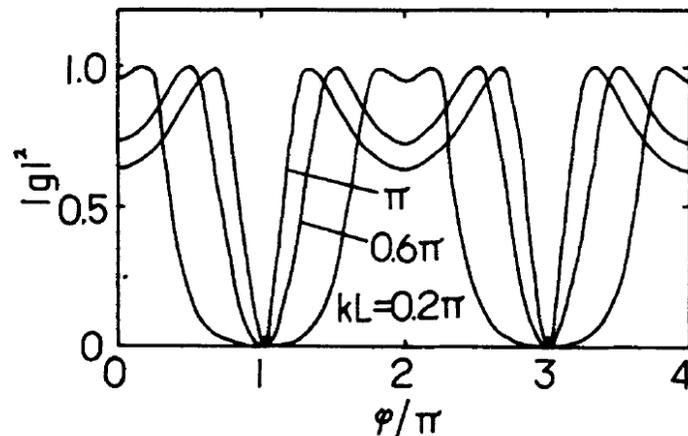
$$|\psi_0 e^{i\pi\frac{\Phi}{\Phi_0}} + \psi_0 e^{-i\pi\frac{\Phi}{\Phi_0}}|^2$$

$$= 2|\psi|^2 \left[1 + \cos 2\pi\frac{\Phi}{\Phi_0} \right]$$

Interference

$$\Phi = \Phi_0 \times (\mathbb{Z} + 1/2)$$

Then, transmission is zero.



Transmission for free fermion

J.-B. Xia (1992)

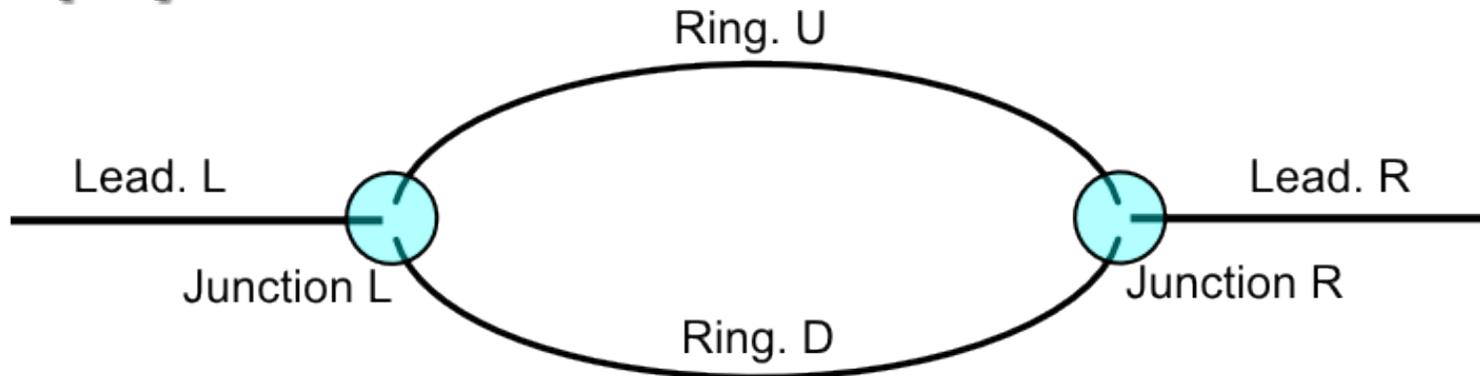
Flux vs Transmission

No transmission at $\Phi = \Phi_0 \times (Z + 1/2)$.

Independent of momentum.

Aharonov-Bohm effect

Set up system



Hamiltonian

$$\mathcal{H} = \mathcal{H}_{ring}^U + \mathcal{H}_{ring}^D + \mathcal{H}_{lead}^L + \mathcal{H}_{lead}^R$$

$$\mathcal{H}_{lead}^{L,R} = \int dx \left[-\frac{1}{2m} \psi_{L,R}^\dagger(x) \frac{\partial^2}{\partial x^2} \psi_{L,R}(x) + \frac{U}{2} \rho_{L,R}(x) \rho_{L,R}(x) \right]$$

$$\mathcal{H}_{ring}^{U,D} = \int_0^L dx \left[-\frac{1}{2m} \tilde{\psi}_{U,D}^\dagger(x) \frac{\partial^2}{\partial x^2} \tilde{\psi}_{U,D}(x) + \frac{U}{2} \rho_{ring}(x) \rho_{ring}(x) \right]$$

$$\tilde{\psi}_{U,D}(x) = \psi_{U,D}(x) \exp\left(\pm i \frac{\pi \Phi}{L \Phi_0} x\right)$$

Boundary condition: Current conservation + strongly coupled limit

Left side boundary condition

$$J_L(0, t) = J_U(0, t) + J_D(0, t)$$

$$\psi_L = \tilde{\psi}_U = \tilde{\psi}_D$$

Right side boundary condition

$$J_R(0, t) = J_U(L, t) + J_D(L, t)$$

$$\psi_R = \tilde{\psi}_U e^{i\pi\Phi/\Phi_0} = \tilde{\psi}_D e^{-i\pi\Phi/\Phi_0}$$

RING-TYPE INTERFEROMETER

Low-energy effective theory

Bosonization formula

$$\psi \sim e^{i\theta}, \quad \rho = \bar{\rho} + \pi^{-1} \partial_x \varphi$$

Branch Hamiltonian

$$\mathcal{H}_{ring}^{U,D} = \frac{v}{2\pi} \int_0^{L/2} dx \left[K \left(\frac{\partial \theta_{U,D}}{\partial x} \right)^2 + \frac{1}{K} \left(\frac{\partial \varphi_{U,D}}{\partial x} \right)^2 \right]$$

$$\mathcal{H}_{lead}^{L,R} = \frac{v}{2\pi} \int dx \left[K \left(\frac{\partial \theta_{L,R}}{\partial x} \right)^2 + \frac{1}{K} \left(\frac{\partial \varphi_{L,R}}{\partial x} \right)^2 \right]$$

Bosonized boundary conditions

Left side boundary condition

$$\partial_x \theta_L = \partial_x \theta_U + \partial_x \theta_D$$

$$\varphi_L = \varphi_U + \varphi_D$$

Right side boundary condition

$$\partial_x \theta_R = \partial_x \theta_U + \partial_x \theta_D$$

$$\varphi_R = \varphi_U + \pi\Phi/\Phi_0 = \varphi_D - \pi\Phi/\Phi_0$$

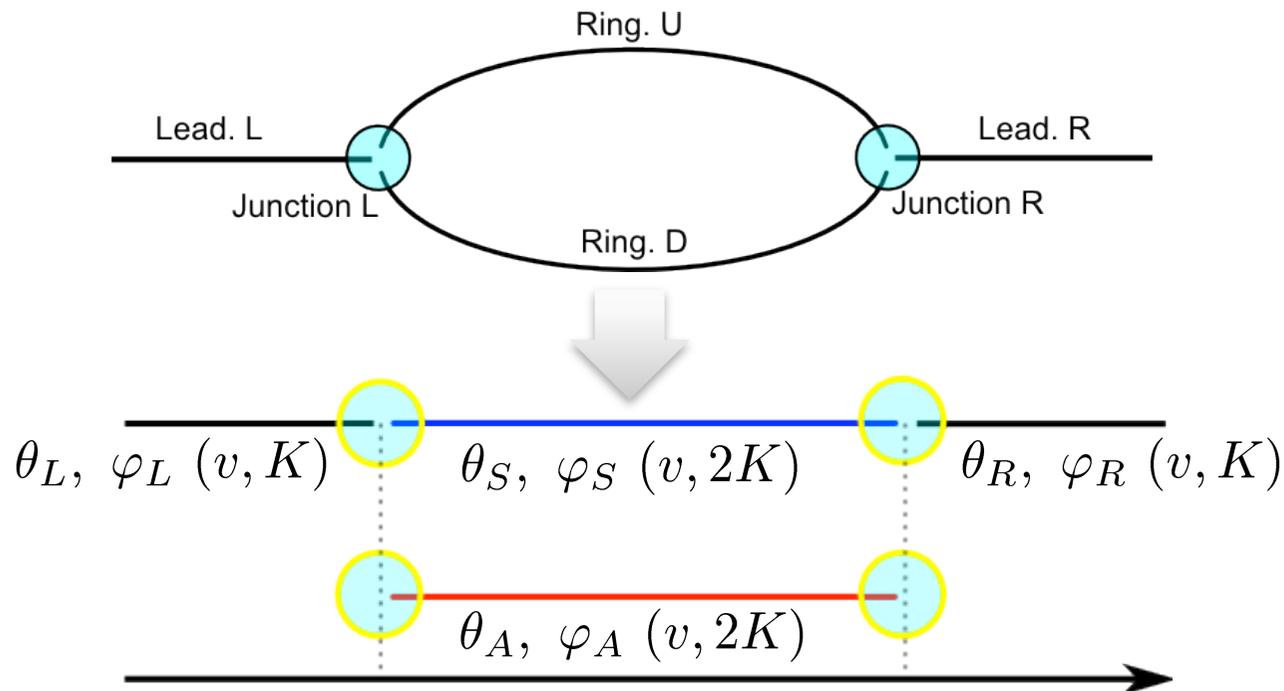
Flux appears only in boundary condition.

→ **absence in boundary conditions for density.**

Low-energy effective theory

Symmetric & anti-symmetric fields

$$\theta_{S,A} = \theta_U \pm \theta_D, \quad \varphi_{S,A} = \frac{\varphi_U \pm \varphi_D}{2}$$



Symmetric & anti-symmetric combination separate the system into two parts.

RING-TYPE INTERFEROMETER

Physics of symmetric part

Analogous to inhomogeneous TL liquid

D. L. Maslov and M. Stone (1995), I. Safi and H. J. Schultz (1995)

Symmetric field: center of mass of packets in the ring.

→ Transport between the leads.

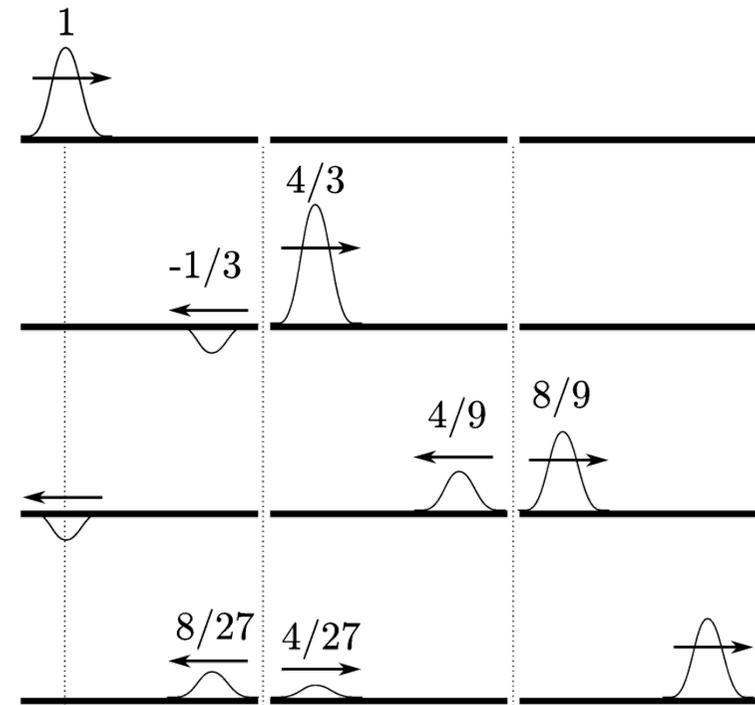
$$\mathcal{H} = \frac{v}{2\pi} \int dx \left[K(x) \left(\frac{\partial \tilde{\theta}}{\partial x} \right)^2 + \frac{1}{K(x)} \left(\frac{\partial \tilde{\varphi}}{\partial x} \right)^2 \right]$$



$$\begin{aligned} \partial_x \theta_L &= \partial_x \theta_S \\ \tilde{\varphi}(0, t) &= \tilde{\varphi}(0, t) \end{aligned}$$

$$\begin{aligned} \partial_x \theta_S &= \partial_x \theta_R \\ \tilde{\varphi}(0, t) &= \tilde{\varphi}(L, t) \end{aligned}$$

In symmetric field + lead sector,
effect of magnetic flux does not appear.
→ Transport between both side leads is
independent of magnetic flux.



Physics of anti-symmetric part

TL liquid with twisted boundary condition.

Quantum ring problem and persistent current.

$$\mathcal{H} = \frac{v}{2\pi} \int_0^L dx \left[2K \left(\frac{\partial \theta_A}{\partial x} \right)^2 + \frac{1}{2K} \left(\frac{\partial \varphi_A}{\partial x} \right)^2 \right]$$



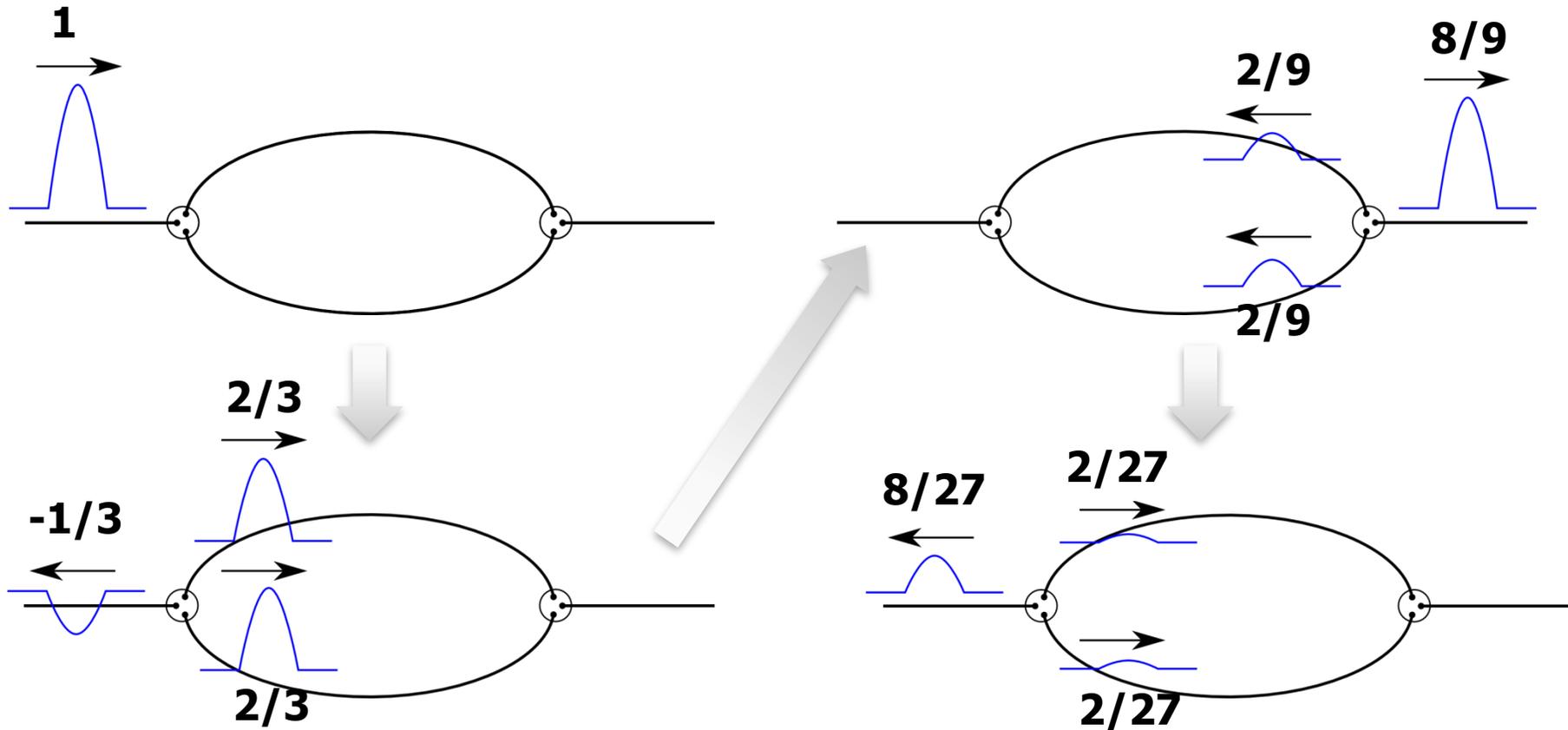
Magnetic flux works as twisting anti-symmetric field.

→ drive persistent current.

RING-TYPE INTERFEROMETER

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Summary of dynamics

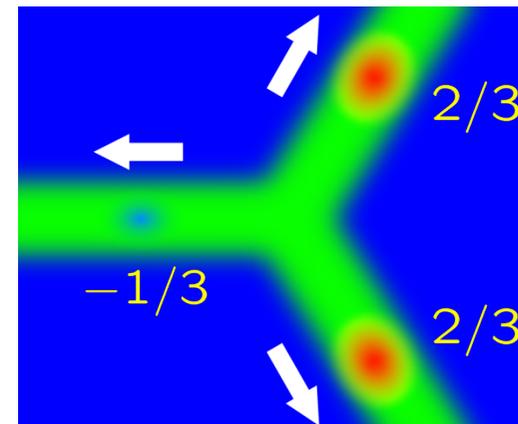


**No interference → Absence of Aharonov-Bohm effect.
Magnetic flux inside ring → Just driving persistent current**

Y-junction for Bose liquid

- Bosons filled in Y-shaped potential.
- Dynamics of density packet in steady boson sea.

Negative reflection & Enhancement of total transmission



Ring-type interferometer for Bose liquid

- Dynamics of density packet in ring-type interferometer.
- Effective magnetic flux inside ring.
- Strongly coupled boundary condition at junctions.

Absence of AB effect.

