

# Spin-electricity conversion induced by spin pumping into topological insulators

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Phys.Rev.Lett. **113**, 196601 (2014)



# Spin injection into “bulk-insulating” topological insulators

PRL 113, 196601 (2014)

PHYSICAL REVIEW LETTERS

week ending  
7 NOVEMBER 2014

## Spin-Electricity Conversion Induced by Spin Injection into Topological Insulators

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theory

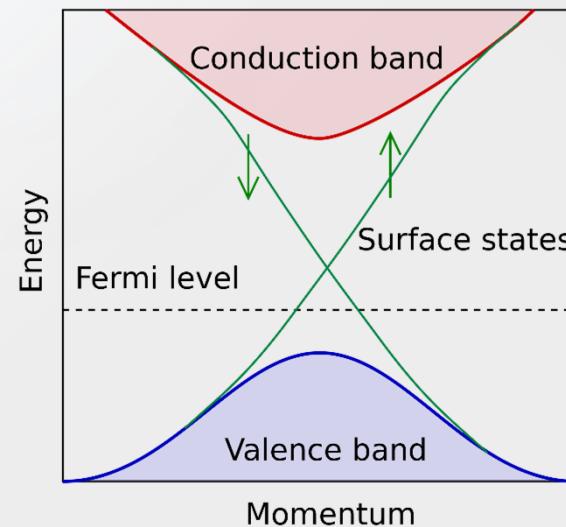
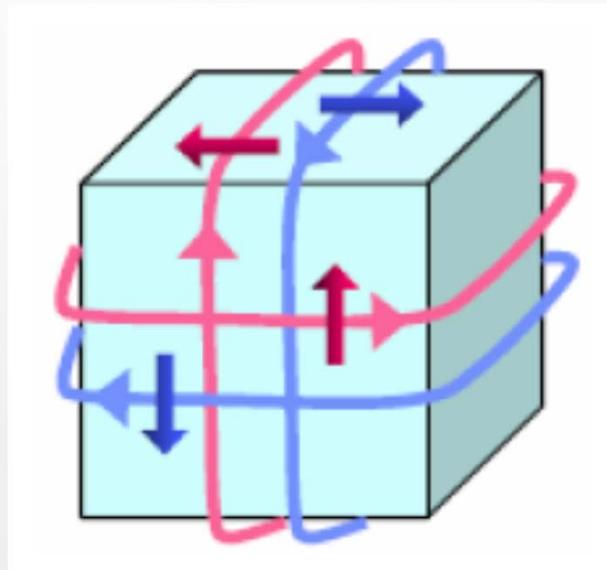
TI sample

boss

arXiv:1312.7091 (2013)  
Phys.Rev.Lett. **113**, 196601 (2014)

TI is a promising material for spintronics application

## Topological Insulator (TI)



Surface: metal  
Interior: Insulator

(1) **spin current  $\neq 0$**   
electric current = 0

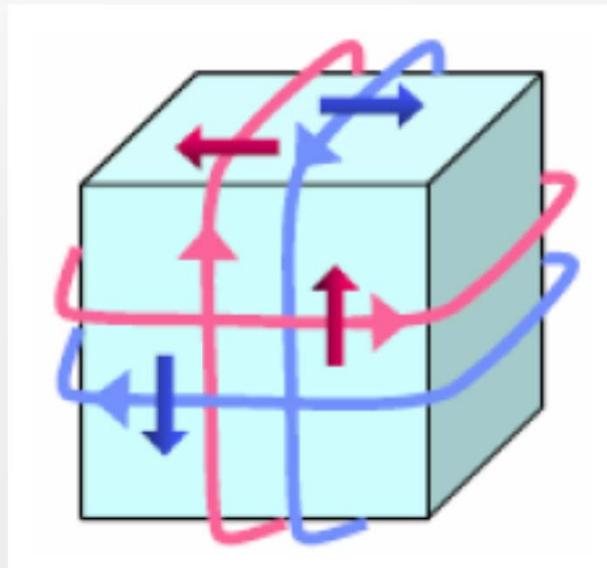
$$\begin{array}{c} \uparrow \\ \longrightarrow \\ \downarrow \end{array} = \begin{array}{c} \uparrow \\ \longrightarrow \end{array}$$

(2) **spin-momentum locking**

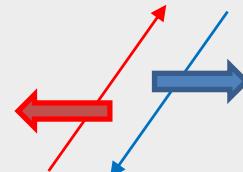
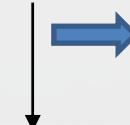
(3) **surface property determined by bulk property**

# Spin injection into surface states of topological insulators

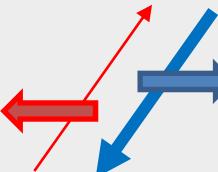
spin injection



spin injection



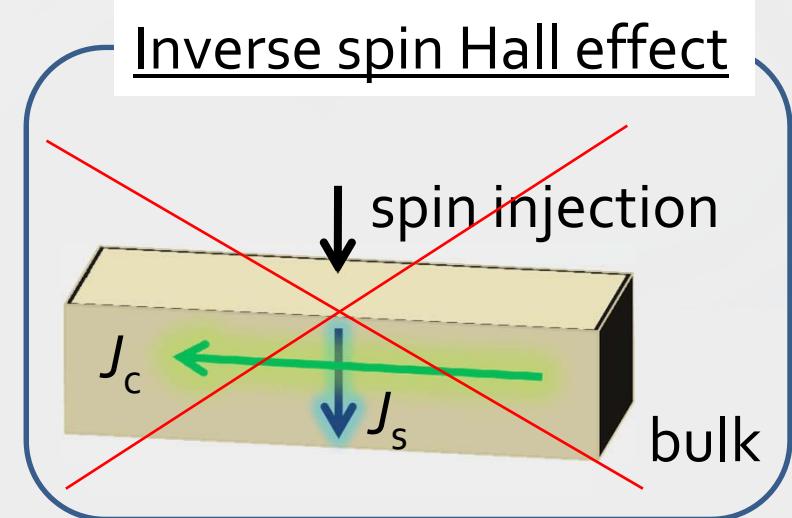
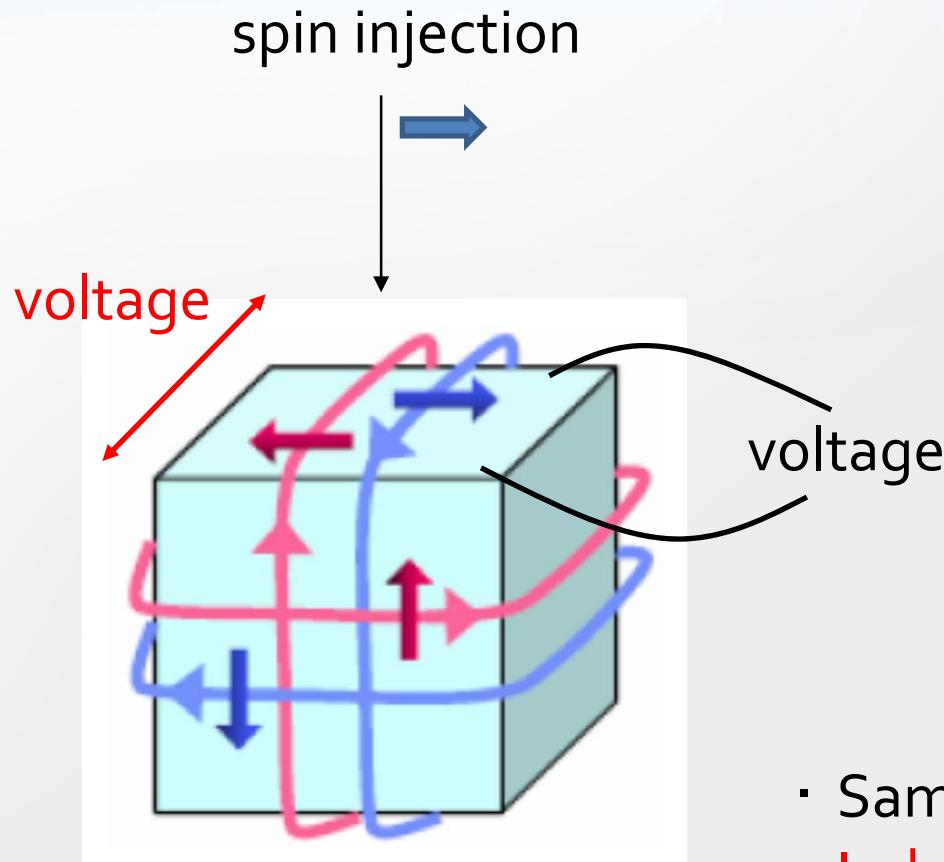
Electric current = 0



$\neq 0$

Electric current induced by spin injection

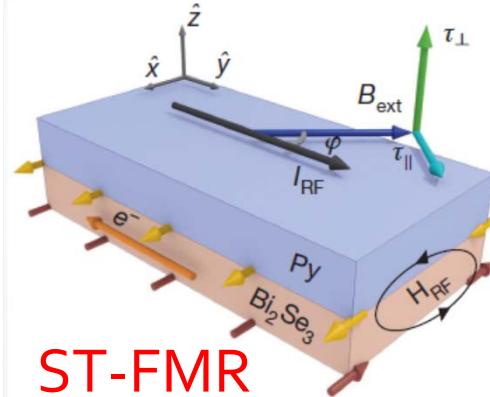
# Spin-electricity conversion induced by spin injection



- Same symmetry as ISHE
- Induced by spin-momentum locking  
(in principle, perfect conversion)

# Some trials for spintronics application of TI

A.R.Mellnik, et al. Nature 511, 449-451 (2014)

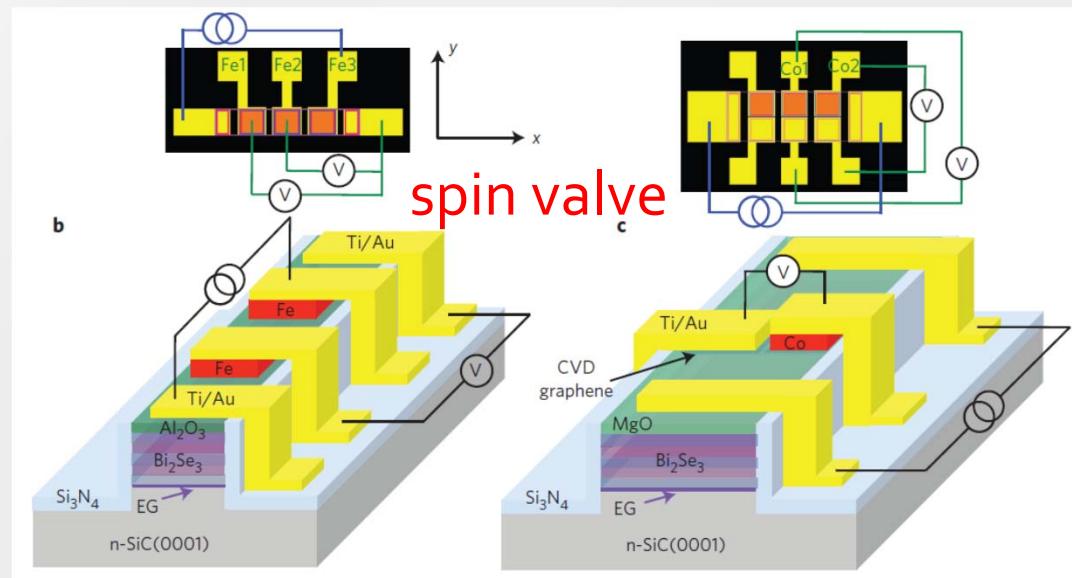


**ST-FMR**

Parameter	Bi <sub>2</sub> Se <sub>3</sub> (this work)
$\theta_{\parallel}$	2.0–3.5
$\sigma_{S,\parallel}$	1.1–2.0

Hall angle > 1

C.H.Li, et al. Nature Nano. 9, 218-224 (2014)

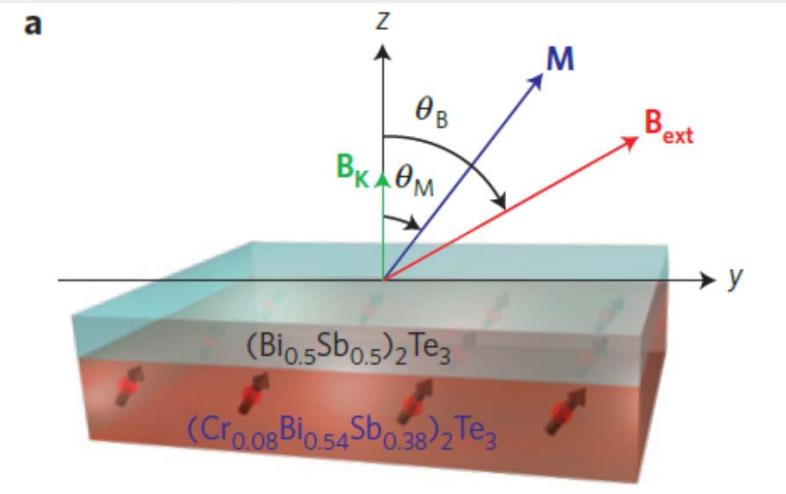


b

spin valve

Y. Fan, et al. Nature Mat. 13, 699-704 (2014)

**Magnetization switching**

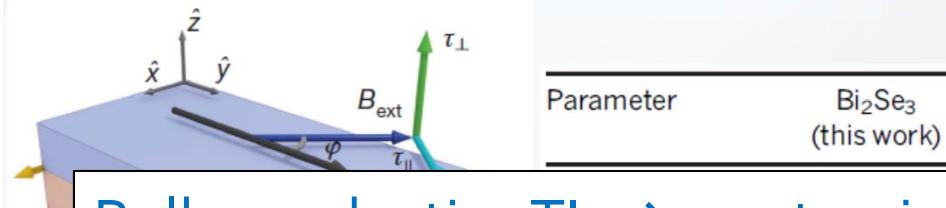


**electric means**

Bulk insulating  
⇒ tiny electric current in TI

# Some trials for spintronics application of TI

A.R.Mellnik, et al. Nature 511, 449-451 (2014)



Y. Fan, et al. Nature Mat. 13, 699-704 (2014)

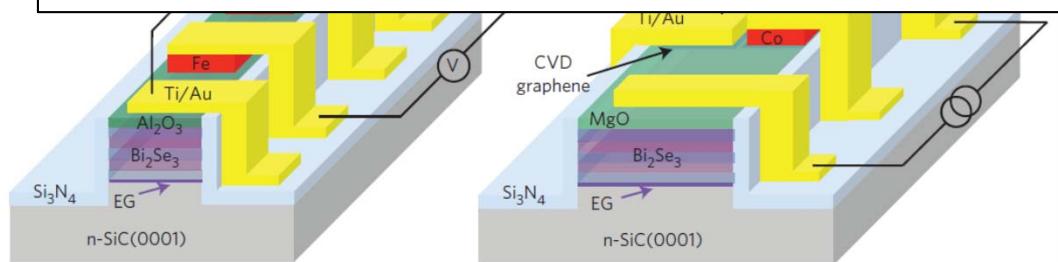
## Magnetization switching

Bulk-conductive TI  $\Rightarrow$  contaminated by bulk carriers  
Bulk-insulating TI  $\Rightarrow$  not-efficient spin transfer (Res. mismatch)

S  
C.H.

How can we observe “pure” surface spin transport of TIs ?

- spin pumping : free from impedance mismatch (K.Ando 2011)
- one more idea to minimize bulk-carrier contribution



Bulk insulating  
 $\Rightarrow$  tiny electric current in TI

# Our experimental setup: Use of “Bulk form” topological insulators

## Bulk-carrier compensated TIs:

$\text{Bi}_{1.5}\text{Sb}_{0.5}\text{Te}_{1.7}\text{Se}_{1.3}$ ,  $\text{BiSbTeSe}_2$ ,  
Sn-doped  $\text{Bi}_2\text{Te}_2\text{Se}$

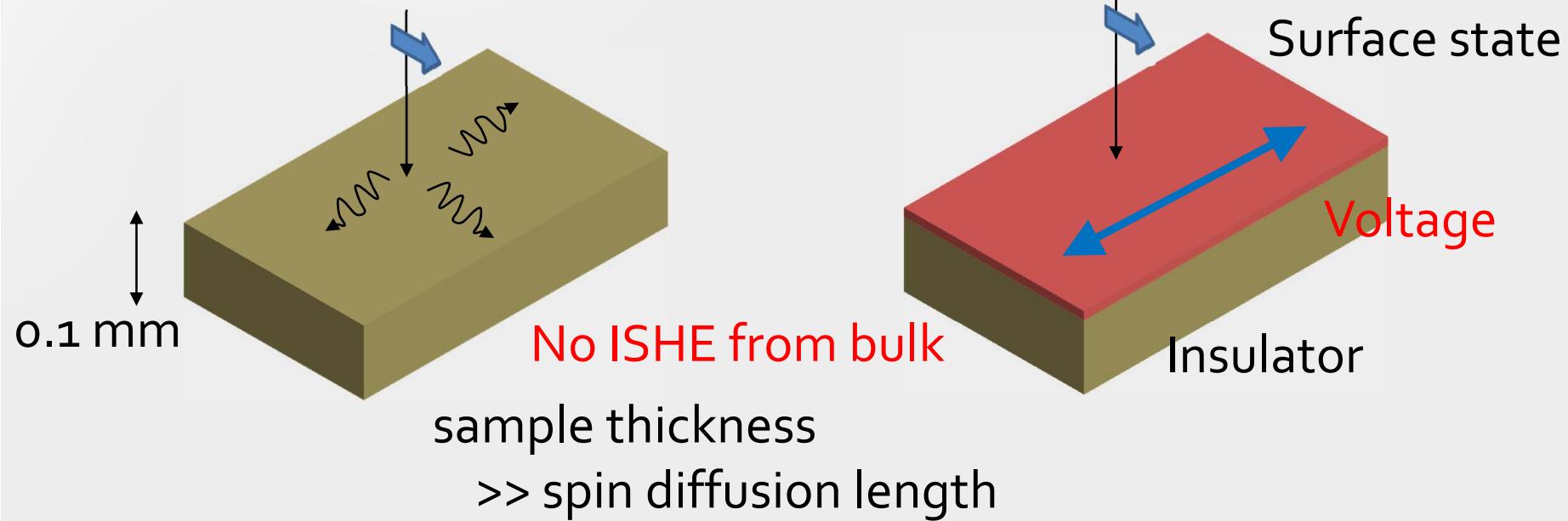
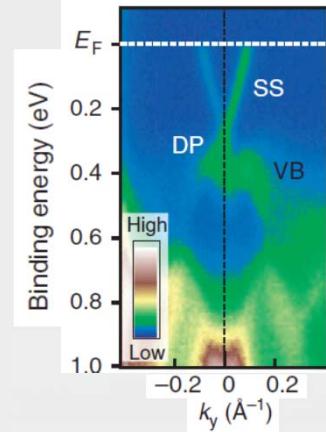
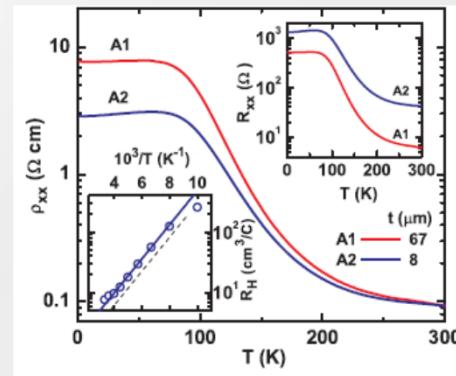
## Bulk-metallic TIs:

$\text{Bi}_2\text{Se}_3$

From Yoichi Ando  
(Osaka Univ.)

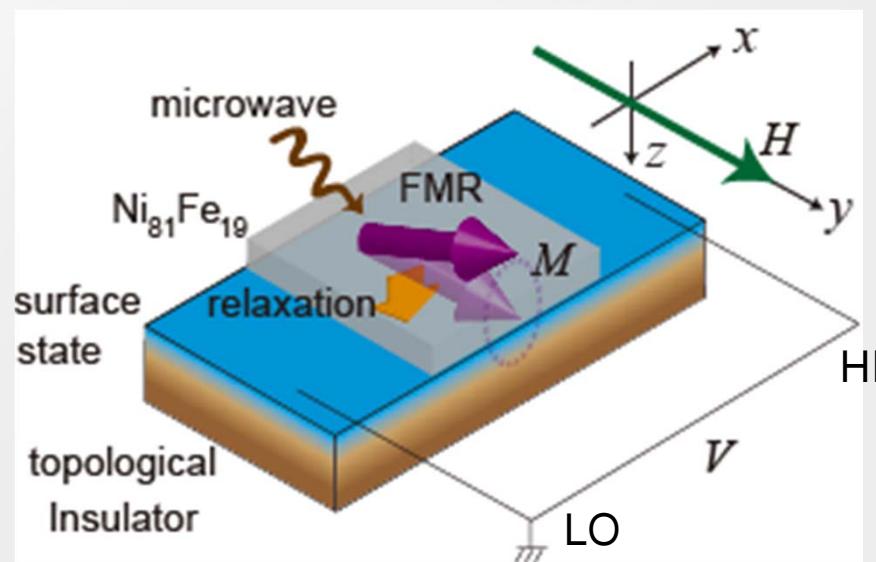
T.Arakane, et al. Nat. Commun. 3, 636 (2011)

A.A.Taskin, et al. PRL 107, 016801 (2011)

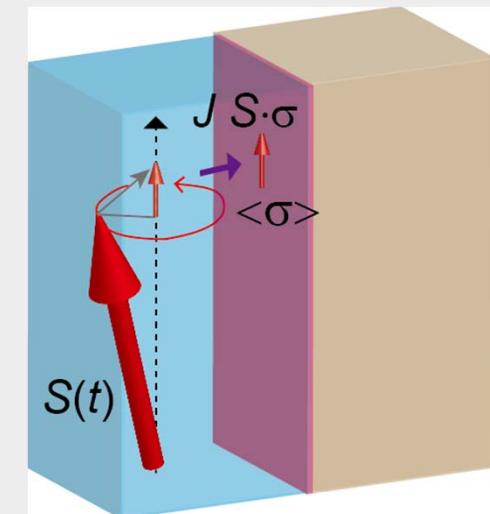


## Methods and samples

- ferromagnet : permalloy ( $\text{Ni}_{81}\text{Fe}_{19}$ ) 20nm thick
- coplanar-type wave guide, network analyser (typically, 5 GHz is used)
- measurement down to 15K (probe station)



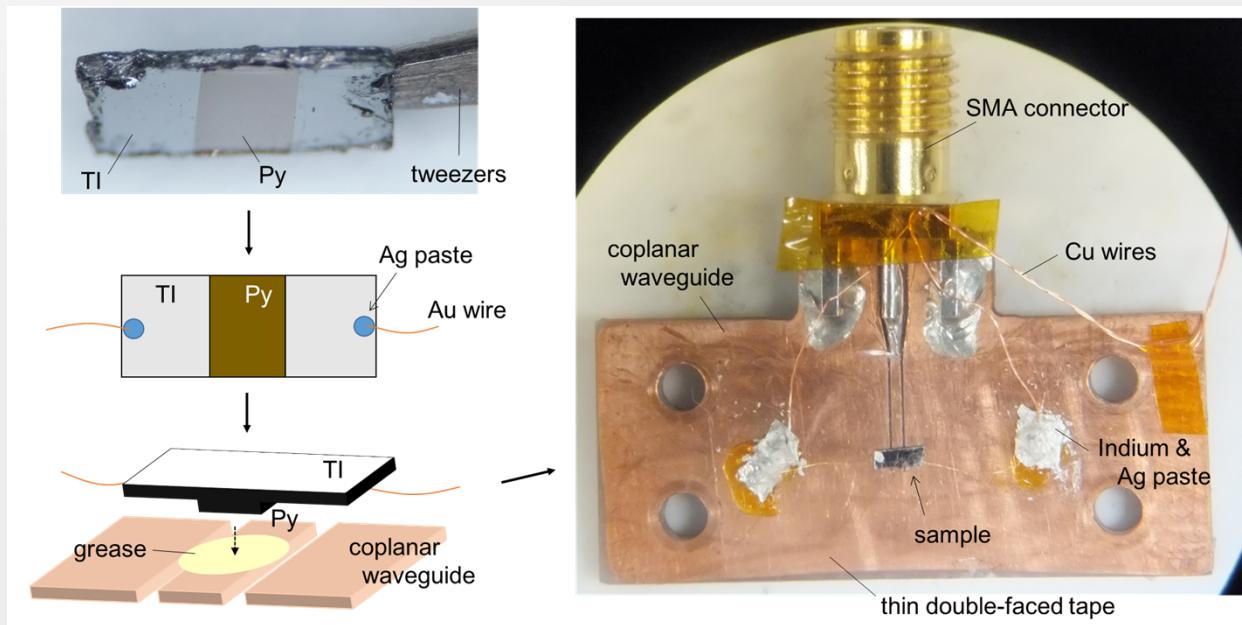
Spin pumping



spin precession  
⇒ spin current

## Methods and samples

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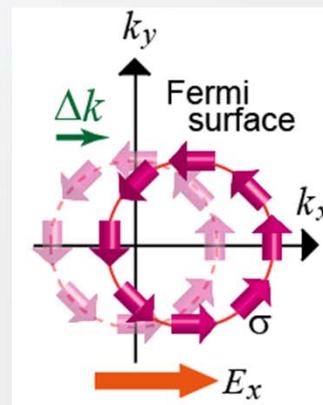
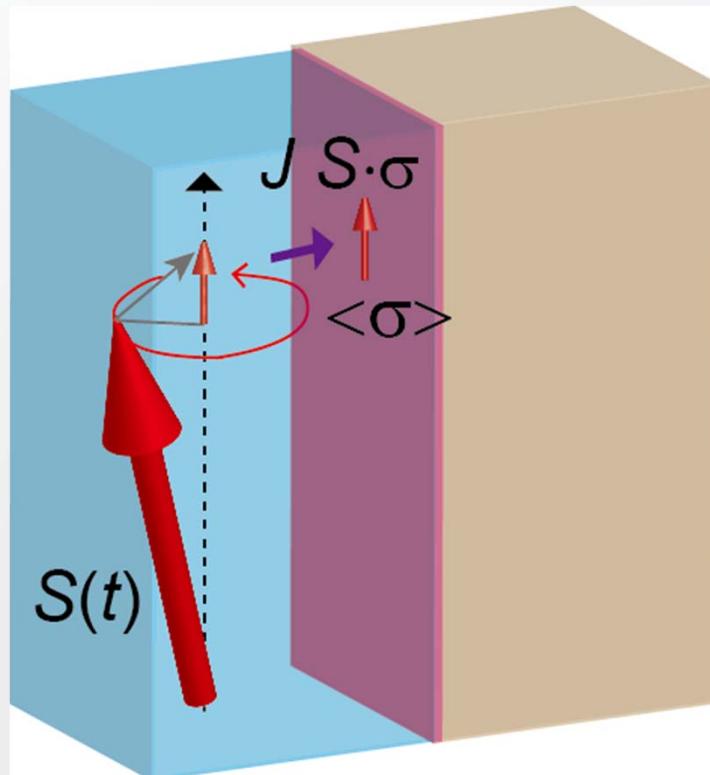


# Spin pumping

Prof. Kentaro Nomura

$$H = v_F (\hat{z} \times \vec{\sigma}) \cdot \vec{p} + J_{sd} \vec{S} \cdot \vec{\sigma} = v_F (\hat{z} \times \vec{\sigma}) \cdot (\vec{p} + e\vec{a}) + J_{sd} S_z \sigma_z$$

## Spin pumping



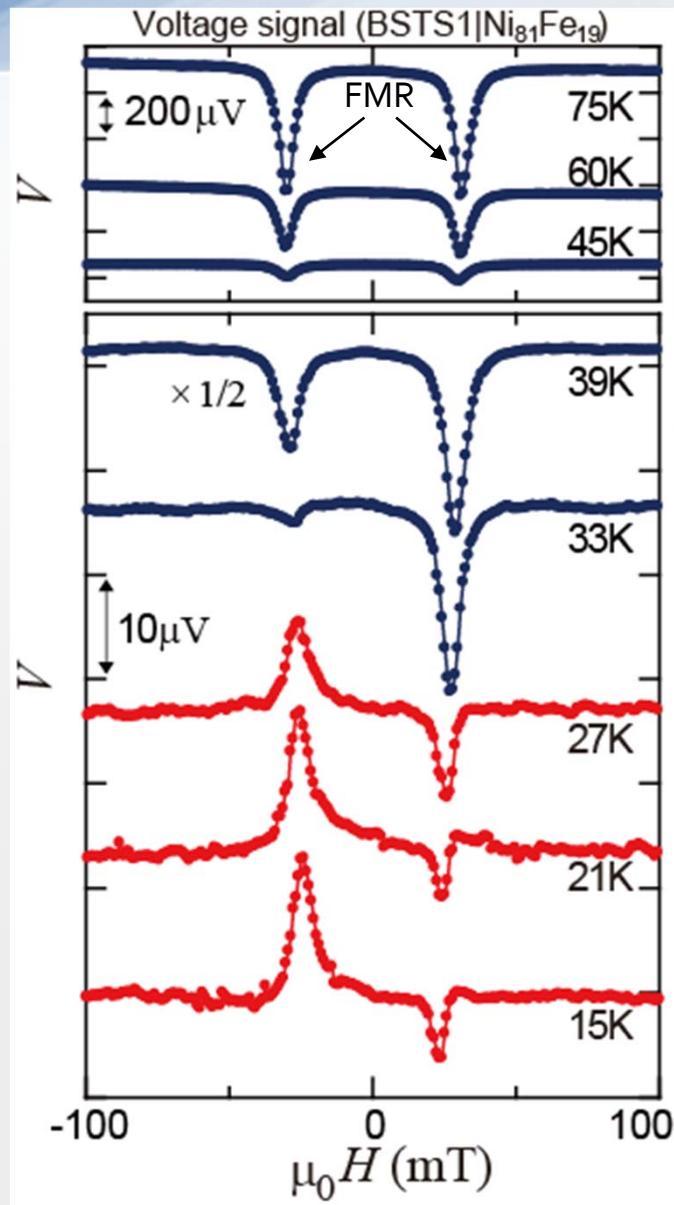
$\dot{\vec{a}}$  = voltage generation

spin exchange interaction  
= torque

$$\vec{T}_{surface} = \gamma \vec{M} \times \vec{J}_{eff} \langle \vec{\sigma} \rangle$$

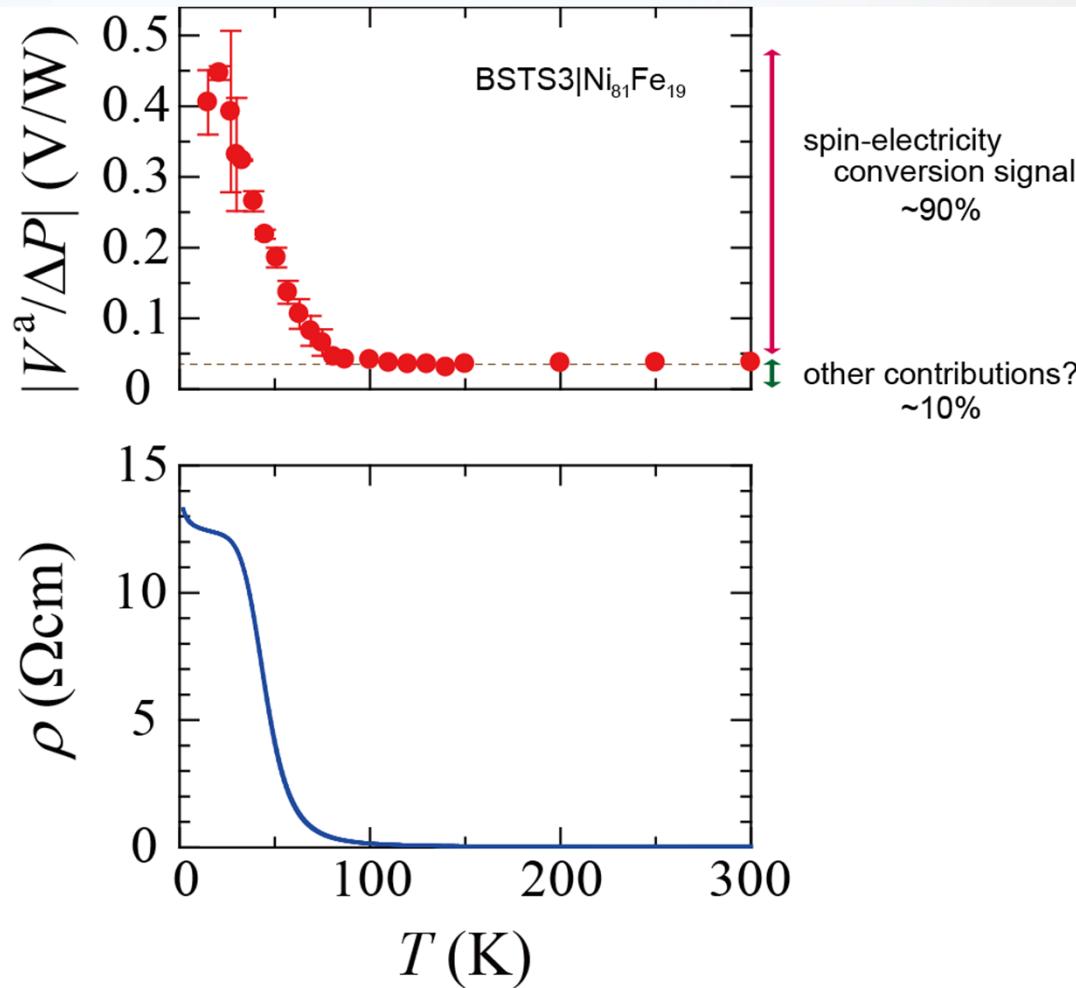
damping enhancement

# Voltage signals at FMR of Py at various temperatures for a $\text{Bi}_{1.5}\text{Sb}_{0.5}\text{Te}_{1.7}\text{Se}_{1.3}$ (BSTS) sample



# Correlation between antisymmetric voltage signals and surface transport

$$V^a = [V(H) - V(-H)]/2$$



$V^a$  increases with decreasing  $T$

$\times 10 !!$

- NOT shunting of AMR

$$\frac{R_{\text{BSTS}}}{R_{\text{Py}} + R_{\text{BSTS}}} V_{\text{AMR}} \quad \times 2$$

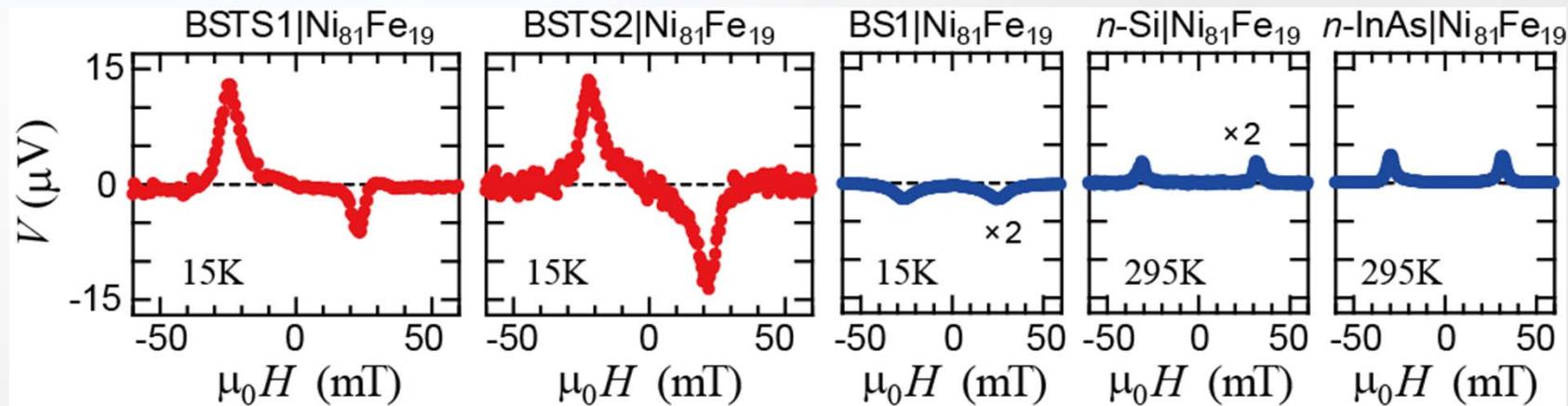
$$R_{\text{BSTS}} \approx R_{\text{py}} (\approx 5\Omega) \text{ at } 300 \text{ K}$$

$$R_{\text{BSTS}} \gg R_{\text{py}} \text{ at low } T_s$$

- NOT Nernst effect

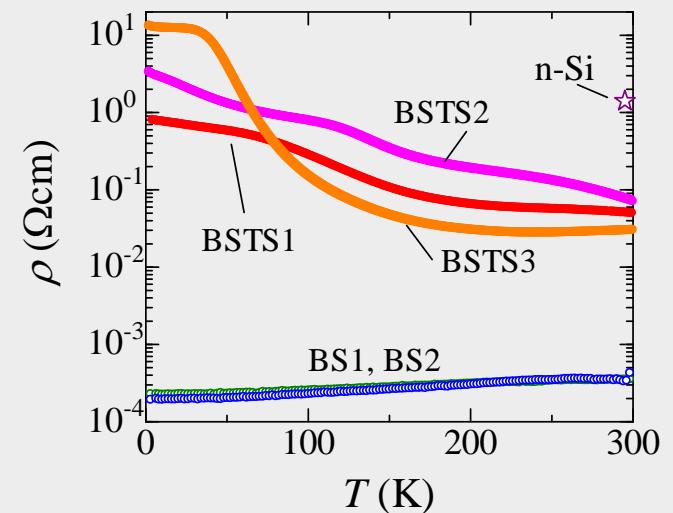
$\propto T$  at low  $T_s$   
(Mott relation)

NO antisymmetric signal in control samples:  
bulk-metallic TIs ( $\text{Bi}_2\text{Se}_3$ ); conventional semiconductors



- Reproducible for all BSTS samples
- not observed for BS1,  $n\text{-Si}$ , or  $n\text{-InAs}$

Anti-symmetric signals arise  
on topological surface states

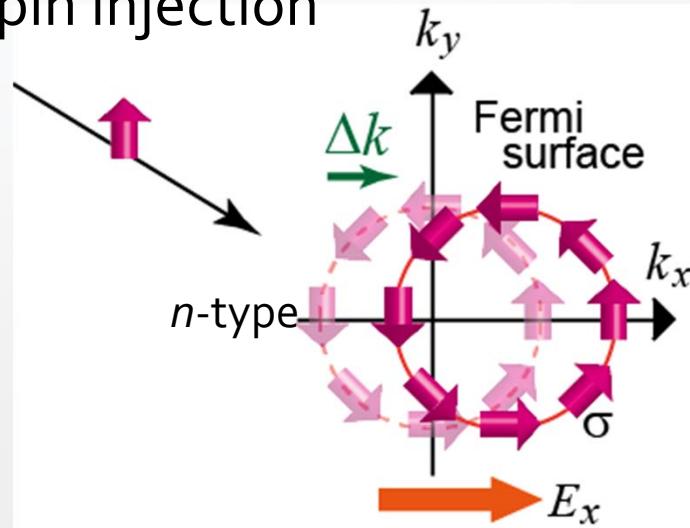


# Mechanism of spin-electricity conversion effect

Prof. Kentaro Nomura

“spin-momentum locking”  $\Rightarrow$  spin-electricity conversion

spin injection



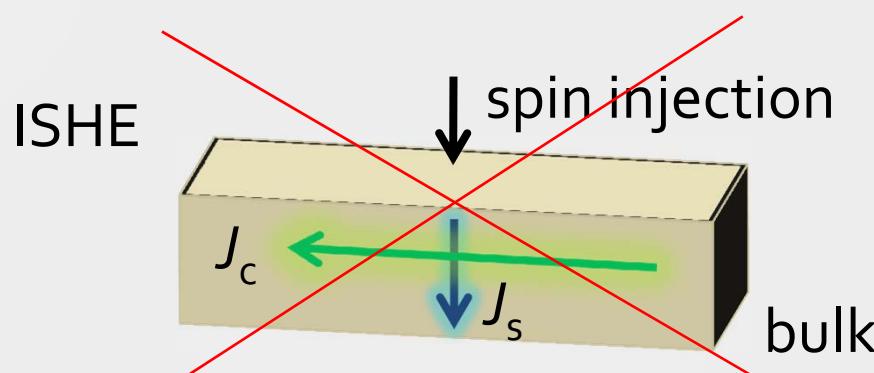
shift of the Fermi circle

$$\Delta k_x = (e\tau/\hbar)E_x$$

$$\frac{\langle \sigma_y \rangle}{A} = \frac{1}{A} \sum_{\vec{k}} f_{\vec{k}} \langle \vec{k},+ | \hat{\sigma}_y | \vec{k},+ \rangle$$

$$E_x = -\frac{4\pi\hbar}{e\tau k_F} \frac{\langle \sigma_y \rangle}{A}$$

same sign as the experiment



if Hall angle  $\theta$  defined,  $\theta \sim 1\%$

## Summary & perspective

Phys.Rev.Lett. **113**, 196601 (2014)

- observation of anti-symmetric signals on millimeter-thick bulk-insulating TIs
- anti-symmetric signals arise on surface states
- build a model : new spin-electricity conversion  
the voltage sign is consistent  
in principle, high efficiency up to 100%, but only ~ 1% now.

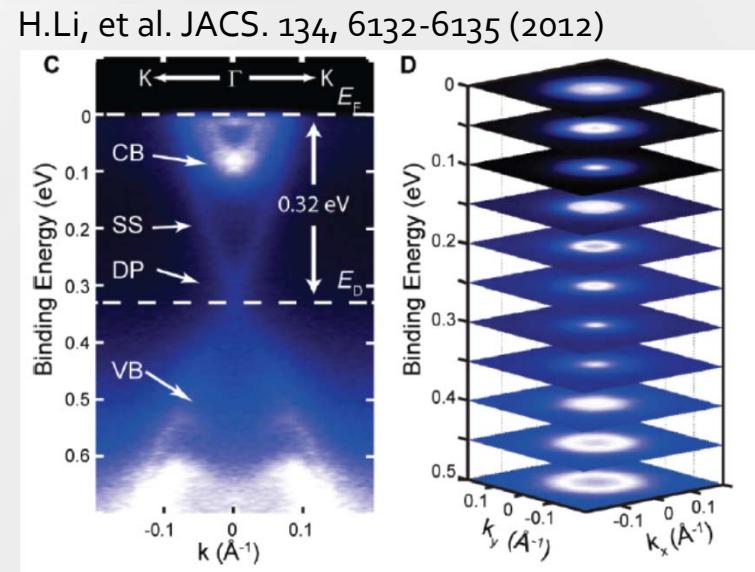
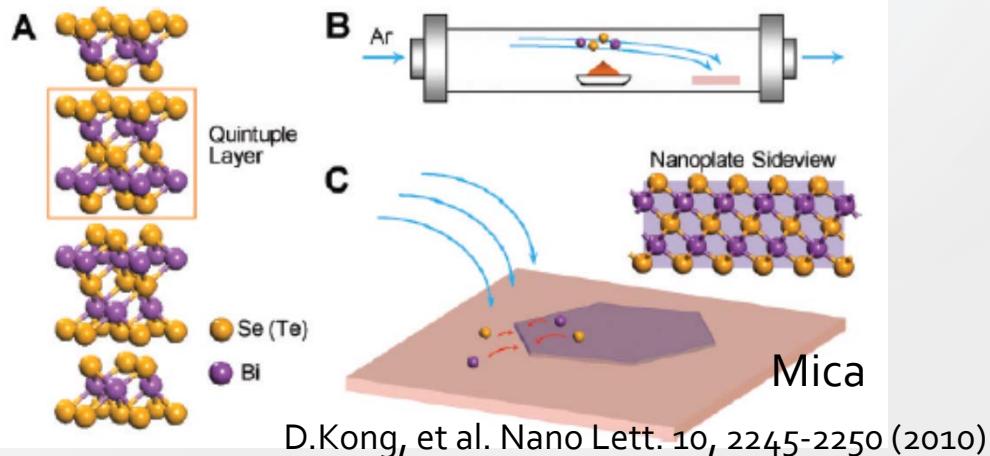


- improve the efficiency  
(improve interface quality, reduce magnetic proximity effect, etc.)
- experiments in BSTS|YIG systems

## Perspective: BSTS on YIG

Growth of BSTS plates on Mica substrate  $\Rightarrow$  transfer onto YIG

## Catalyst-free vapor solid method



BSTS plate (~ 1mm size)  
⇒ transfer onto YIG

By Dr. Tanabe (Tohoku Univ.)



$\text{Bi}_{1.5}\text{Sb}_{0.5}\text{Te}_{1.7}\text{Se}_{1.3}$  plate  
(50 nm thick)

