

Pairing states in Kane-Mele-based systems

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Spin-orbit coupling and pairing symmetry

Cf) chiral p-wave superconductor Sr_2RuO_4

possible spin triplet states in D_{4h}

Sigrist & Ueda RMP ('91)

Γ	$d(\mathbf{k})$
A_{1u}	$\hat{x}\hat{k}_x + \hat{y}\hat{k}_y$
A_{2u}	$\hat{x}\hat{k}_y - \hat{y}\hat{k}_x$
B_{1u}	$\hat{x}\hat{k}_x - \hat{y}\hat{k}_y$
B_{1u}	$\hat{x}\hat{k}_y + \hat{y}\hat{k}_x$
E_u (chiral states)	$\hat{z}(\hat{k}_x \pm i\hat{k}_y)$

- These five states have degenerated condensation energy
- The degeneracy is lifted by SOC

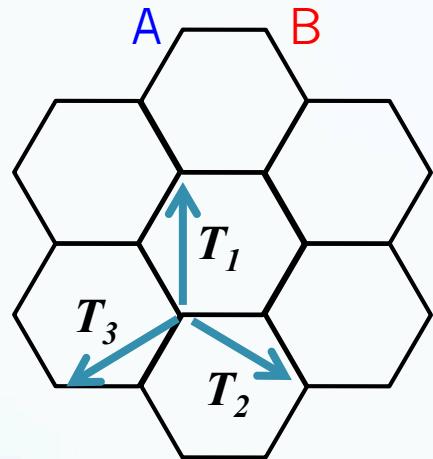


Chiral p-wave

Ng&Sigrist ('99)
Yanase&Ogata('03)

Spin-orbit coupling and band topology

Kane-Mele insulator C.L. Kane and E. Mele PRL('05)('05)

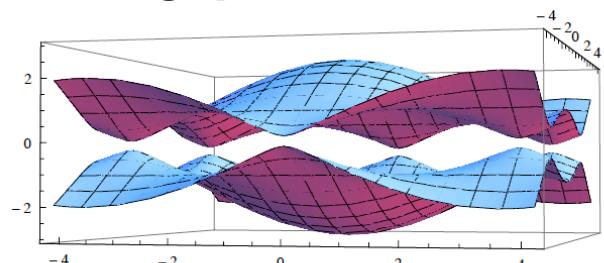


$H = \text{NN hopping} + \text{SOC}$

$$\hat{\mathcal{H}}_{\mathbf{k}}^{\text{SOC}} = \lambda_{SO} \sum_{i=1,2,3} \sin \mathbf{k} \cdot \mathbf{T}_i \hat{\sigma}_z \otimes \hat{s}_z$$

sublattice (A,B) spin (\uparrow, \downarrow)

band gap $\propto \lambda_{SO}$



$$\lambda_{SO}/t = 0.01$$

non-trivial topology

$$\begin{aligned} & \text{Spin Chern \#} \\ &= \int_{BZ} \frac{d^2 k}{(2\pi)^2 i} \text{rot} \{ \langle u_{\mathbf{k}\uparrow} | \nabla_{\mathbf{k}} | u_{\mathbf{k}\uparrow} \rangle - \langle u_{\mathbf{k}\downarrow} | \nabla_{\mathbf{k}} | u_{\mathbf{k}\downarrow} \rangle \} \\ &= \frac{\lambda_{SO}}{|\lambda_{SO}|} \longrightarrow \text{The helical edge state} \end{aligned}$$

QSHE

Our issue:

How does the topological SOC of KM model
affect **the pairing states** ?

1. Cooperon in KM insulator

with S. Tsuchiya, E. Arahata, and M. Sigrist

Cf) Nozieres-Pistolesi ('98)

Band insulator + attractive U

Increasing U \Rightarrow Cooperon excitations

Increasing more... \Rightarrow SC due to Cooperon condensation

The KM SOC favors the helical p-wave Cooperon channel

(2. SrPtAs: an example of KM metal and its superconductivity)

with M. Fischer, A. Schnyder, T. Neupert, R. Thomale...

Y. Imai, and M. Sigrist

Cooperons in Kane-Mele model

With S. Tsuchiya, E. Arahata, and M. Sigrist

$$H_{KM} = \sum_{\mathbf{k}} \left\{ \Psi_{\mathbf{k}\uparrow}^\dagger \begin{pmatrix} \lambda_{\mathbf{k}} & t_{\mathbf{k}} \\ t_{\mathbf{k}}^* & -\lambda_{\mathbf{k}} \end{pmatrix} \Psi_{\mathbf{k}\uparrow} + \Psi_{\mathbf{k}\downarrow}^\dagger \begin{pmatrix} -\lambda_{\mathbf{k}} & t_{\mathbf{k}} \\ t_{\mathbf{k}}^* & \lambda_{\mathbf{k}} \end{pmatrix} \Psi_{\mathbf{k}\downarrow} \right\}$$

Fermion OP

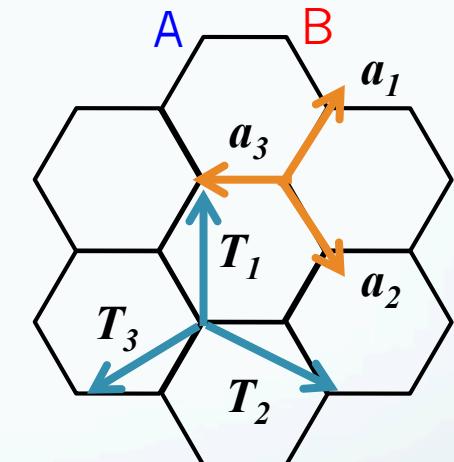
$$\Psi_{\mathbf{k}s} = \begin{pmatrix} A_{\mathbf{k}s} \\ B_{\mathbf{k}s} \end{pmatrix}$$

NN hopping

$$t_{\mathbf{k}} = t \sum_i e^{i\mathbf{k} \cdot \mathbf{a}_i}$$

SOC

$$\lambda_{\mathbf{k}} = \lambda_{SO} \sum_i \sin \mathbf{k} \cdot \mathbf{T}_i$$



When we introduce...

$$H_{on} = -U \sum_i n_{i\uparrow} n_{i\downarrow}$$

U>Uc
→

S-wave
condensate

Cooperons in Kane-Mele model

With S. Tsuchiya, E. Arahata, and M. Sigrist

$$H_{KM} = \sum_{\mathbf{k}} \left\{ \Psi_{\mathbf{k}\uparrow}^\dagger \begin{pmatrix} \lambda_{\mathbf{k}} & t_{\mathbf{k}} \\ t_{\mathbf{k}}^* & -\lambda_{\mathbf{k}} \end{pmatrix} \Psi_{\mathbf{k}\uparrow} + \Psi_{\mathbf{k}\downarrow}^\dagger \begin{pmatrix} -\lambda_{\mathbf{k}} & t_{\mathbf{k}} \\ t_{\mathbf{k}}^* & \lambda_{\mathbf{k}} \end{pmatrix} \Psi_{\mathbf{k}\downarrow} \right\}$$

Fermion OP

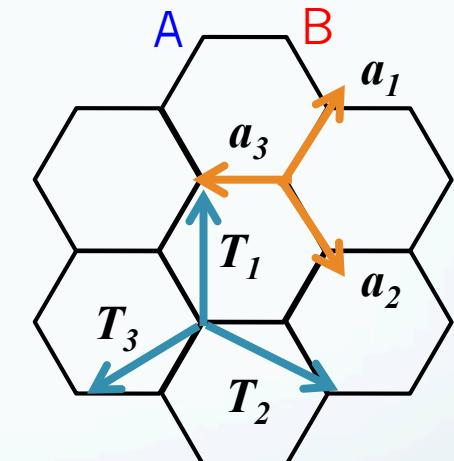
$$\Psi_{\mathbf{k}s} = \begin{pmatrix} A_{\mathbf{k}s} \\ B_{\mathbf{k}s} \end{pmatrix}$$

NN hopping

$$t_{\mathbf{k}} = t \sum_i e^{i\mathbf{k} \cdot \mathbf{a}_i}$$

SOC

$$\lambda_{\mathbf{k}} = \lambda_{SO} \sum_i \sin \mathbf{k} \cdot \mathbf{T}_i$$



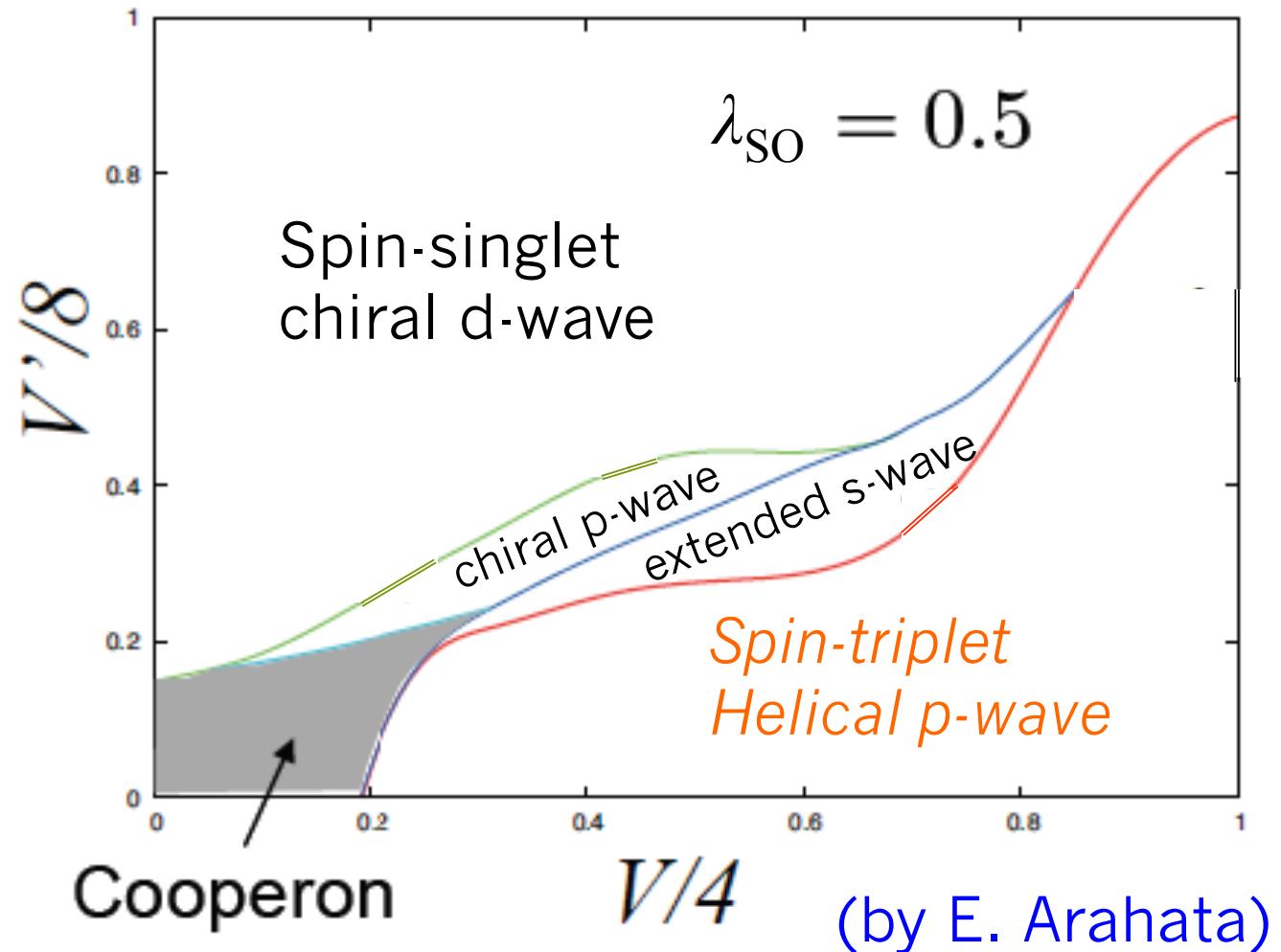
When we introduce...

$$H_{\text{int}} = -V \sum_{\langle i,j \rangle} n_i n_j - V' \sum_{\langle\langle i,j \rangle\rangle} n_i n_j$$

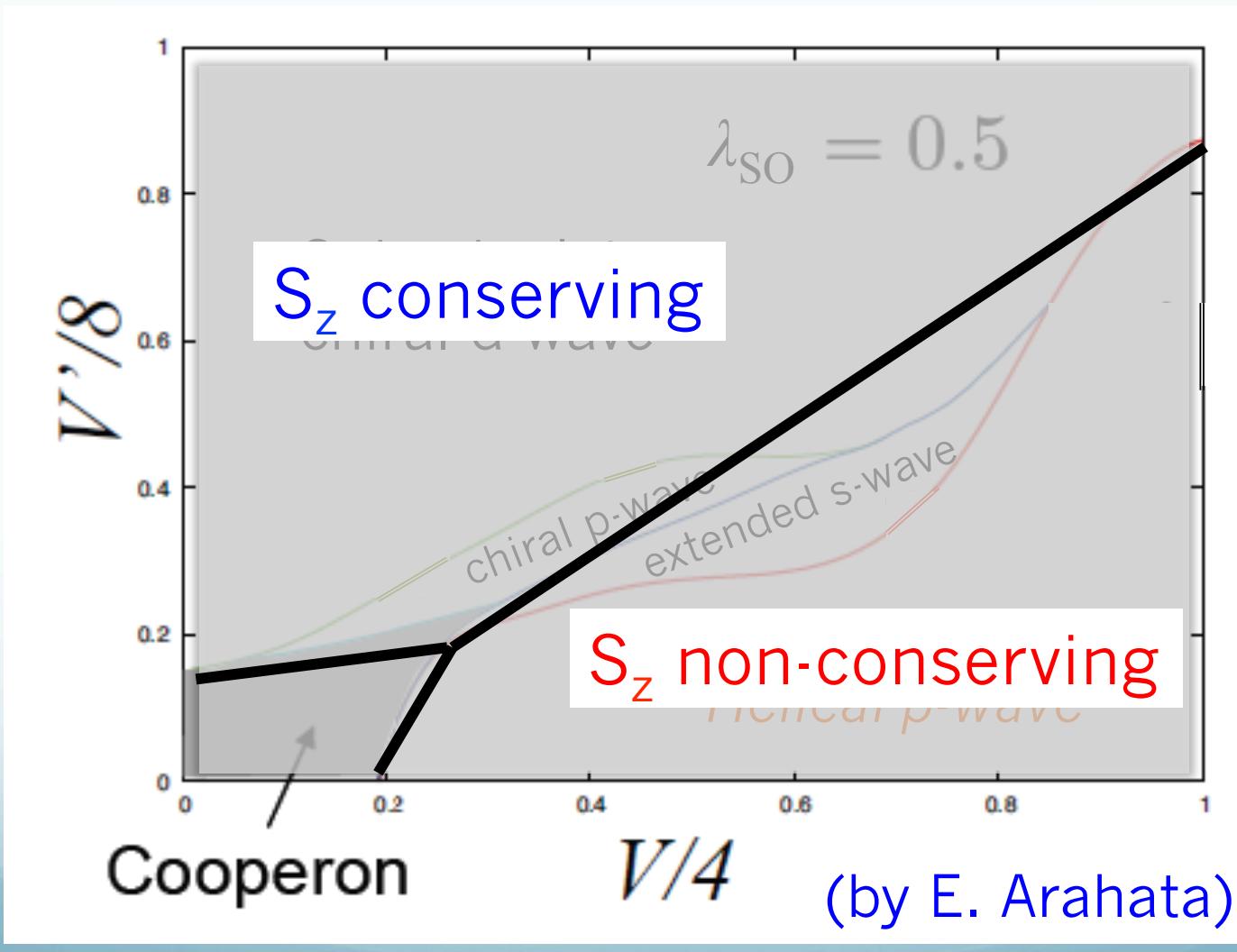
$V > V_c$
 $V > V_{c'}$

?

Analysis of the Gap equation at T=0 with $V(NN)$ and $V'(NNN)$ attraction



Analysis of the Gap equation at T=0 with $V(NN)$ and $V'(NNN)$ attraction



Eigenstates of H_{KM}

$$\begin{aligned}\epsilon_{\mathbf{k}}^{(+)} &= +\sqrt{|t_{\mathbf{k}}|^2 + \lambda_{\mathbf{k}}^2} \\ &\equiv \epsilon_{\mathbf{k}}\end{aligned}$$

$$\begin{aligned}\epsilon_{\mathbf{k}}^{(-)} &= -\sqrt{|t_{\mathbf{k}}|^2 + \lambda_{\mathbf{k}}^2} \\ &= -\epsilon_{\mathbf{k}}\end{aligned}$$

$$|\mathbf{k} \uparrow +\rangle = \begin{pmatrix} u_{\mathbf{k}} e^{i\phi_{\mathbf{k}}} \\ v_{\mathbf{k}} \end{pmatrix}$$

$$|\mathbf{k} \uparrow -\rangle = \begin{pmatrix} v_{\mathbf{k}} \\ -u_{\mathbf{k}} e^{-i\phi_{\mathbf{k}}} \end{pmatrix}$$

$$|\mathbf{k} \downarrow +\rangle = \begin{pmatrix} v_{\mathbf{k}} \\ u_{\mathbf{k}} e^{-i\phi_{\mathbf{k}}} \end{pmatrix}$$

$$|\mathbf{k} \downarrow -\rangle = \begin{pmatrix} u_{\mathbf{k}} e^{i\phi_{\mathbf{k}}} \\ -v_{\mathbf{k}} \end{pmatrix}$$

2-fold (spin) degeneracy

$$u_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 + \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$$

$$v_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 - \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$$

$$\phi_{\mathbf{k}} = \arg(t_{\mathbf{k}})$$

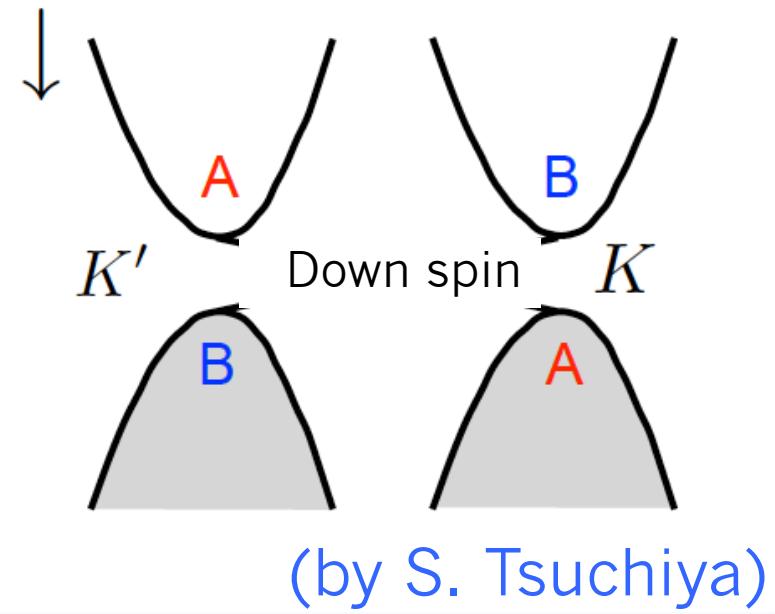
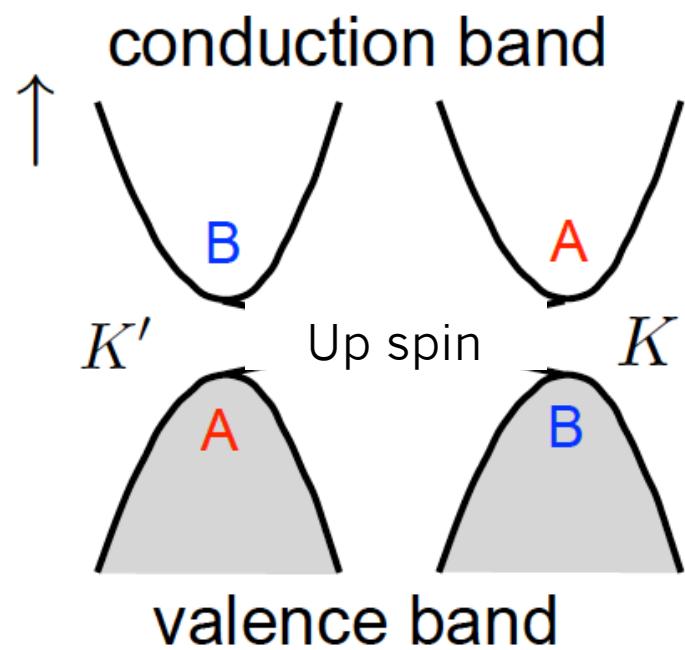
- Pairing between two states around K and K' points would be dominant due to the narrow gap

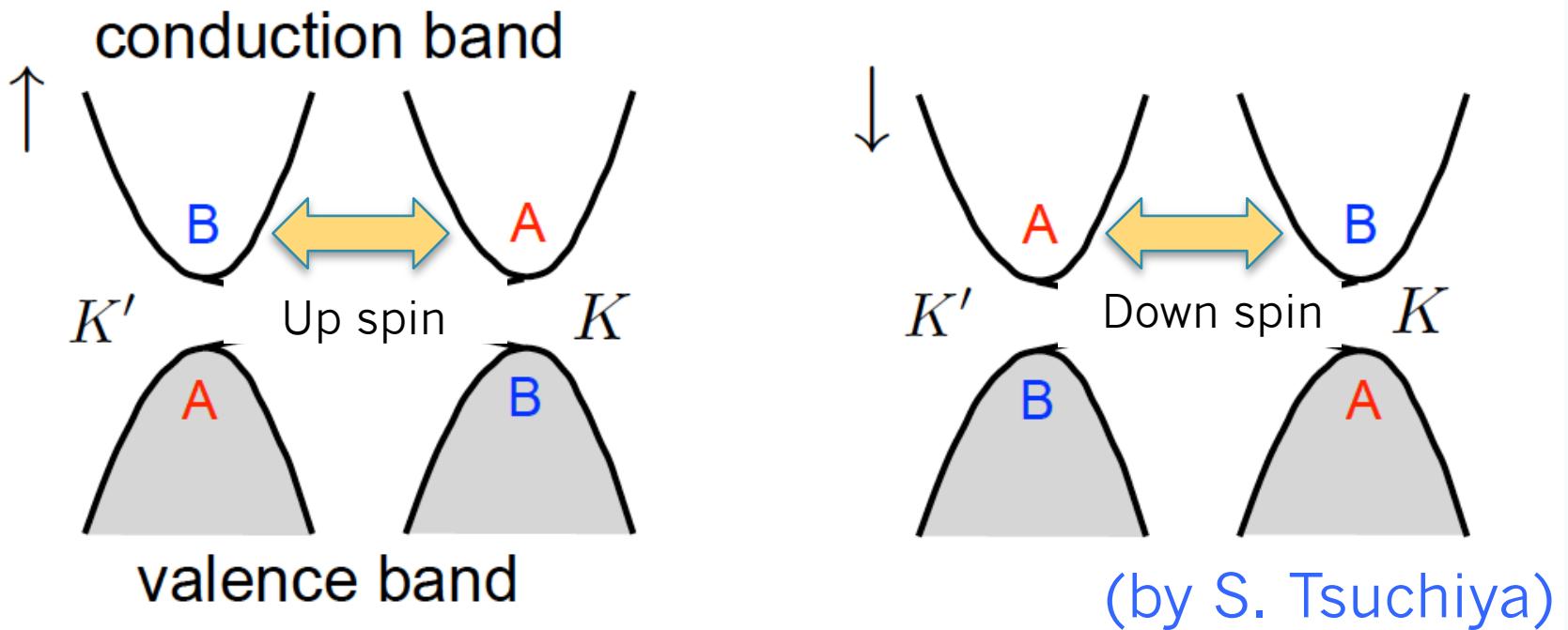


focus on the behavior around K and K' points

	around K	around K'
$ t_{\mathbf{k}} $	0	0
$\phi_{\mathbf{k}} = \arg(t_{\mathbf{k}})$	$\tan^{-1} \frac{\delta k_y}{\delta k_x} \equiv \theta_{\delta \mathbf{k}}$	$-\tan^{-1} \frac{\delta k_y}{\delta k_x} = -\theta_{\delta \mathbf{k}}$
$\lambda_{\mathbf{k}}$	λ_{SO}	$-\lambda_{SO}$
$\epsilon_{\mathbf{k}}$	$ \lambda_{SO} $	$ \lambda_{SO} $
	$\cdot \lambda_{SO} > 0$	$\cdot \lambda_{SO} > 0$
$u_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 + \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$	1	0
$v_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 - \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$	0	1

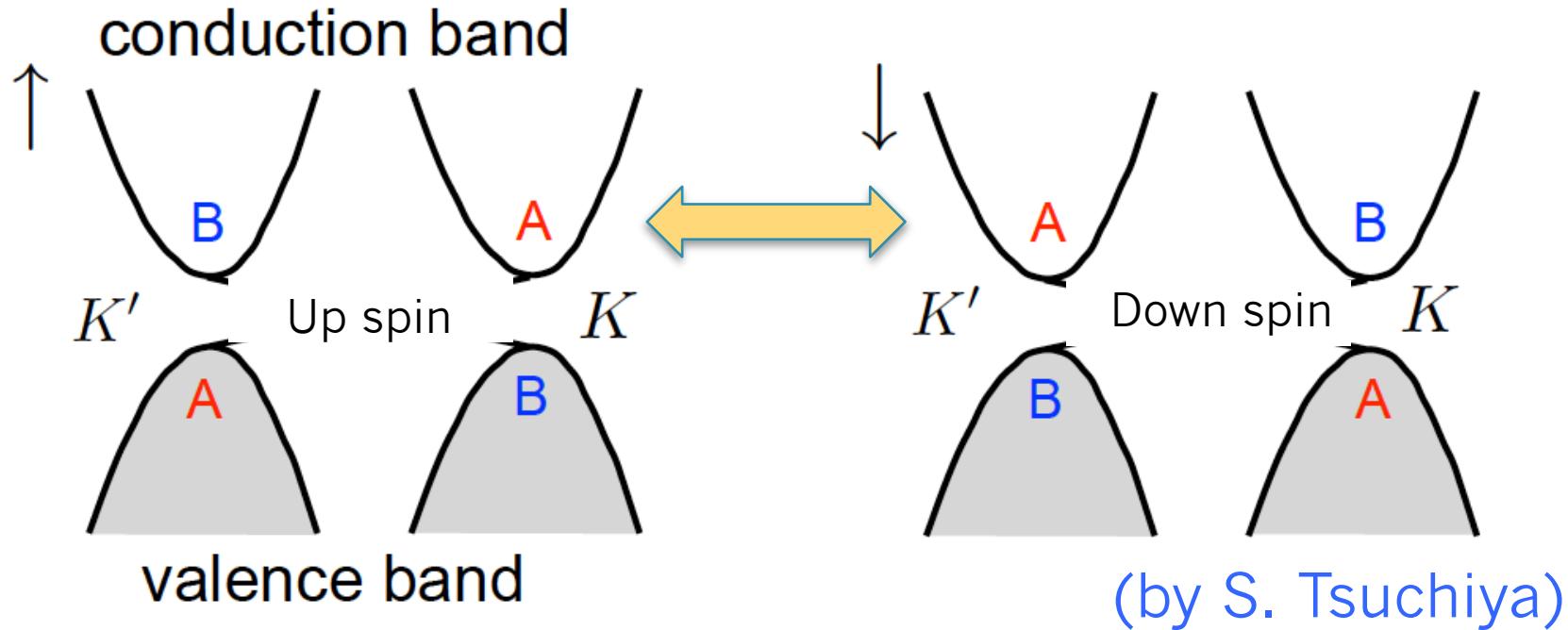
	around K	around K'
$ t_{\mathbf{k}} $	0	0
$\phi_{\mathbf{k}} = \arg(t_{\mathbf{k}})$	$\tan^{-1} \frac{\delta k_y}{\delta k_x} \equiv \theta_{\delta \mathbf{k}}$	$-\tan^{-1} \frac{\delta k_y}{\delta k_x} = -\theta_{\delta \mathbf{k}}$
$\lambda_{\mathbf{k}}$	λ_{SO}	$-\lambda_{SO}$
$\epsilon_{\mathbf{k}}$	$ \lambda_{SO} $	$ \lambda_{SO} $
	$\cdot \lambda_{SO} > 0$	$\cdot \lambda_{SO} > 0$
$u_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 + \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$	1	0
$v_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 - \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$	0	1
$ \mathbf{k} \uparrow +\rangle = \begin{pmatrix} u_{\mathbf{k}} e^{i\phi_{\mathbf{k}}} \\ v_{\mathbf{k}} \end{pmatrix}$	$\begin{pmatrix} e^{-i\theta} \\ 0 \end{pmatrix}$	$\begin{pmatrix} 0 \\ 1 \end{pmatrix}$
$ \mathbf{k} \downarrow +\rangle = \begin{pmatrix} v_{\mathbf{k}} \\ u_{\mathbf{k}} e^{-i\phi_{\mathbf{k}}} \end{pmatrix}$	$\begin{pmatrix} 0 \\ e^{i\theta} \end{pmatrix}$	$\begin{pmatrix} 1 \\ 0 \end{pmatrix}$
$ \mathbf{k} \uparrow -\rangle = \begin{pmatrix} v_{\mathbf{k}} \\ -u_{\mathbf{k}} e^{-i\phi_{\mathbf{k}}} \end{pmatrix}$	$\begin{pmatrix} 0 \\ -e^{i\theta} \end{pmatrix}$	$\begin{pmatrix} 1 \\ 0 \end{pmatrix}$
$ \mathbf{k} \downarrow -\rangle = \begin{pmatrix} u_{\mathbf{k}} e^{i\phi_{\mathbf{k}}} \\ -v_{\mathbf{k}} \end{pmatrix}$	$\begin{pmatrix} e^{-i\theta} \\ 0 \end{pmatrix}$	$\begin{pmatrix} 0 \\ -1 \end{pmatrix}$
also sublattice eigenstates		





NN (inter-sublattice) attraction
⇒ spin-polarized components of pairing

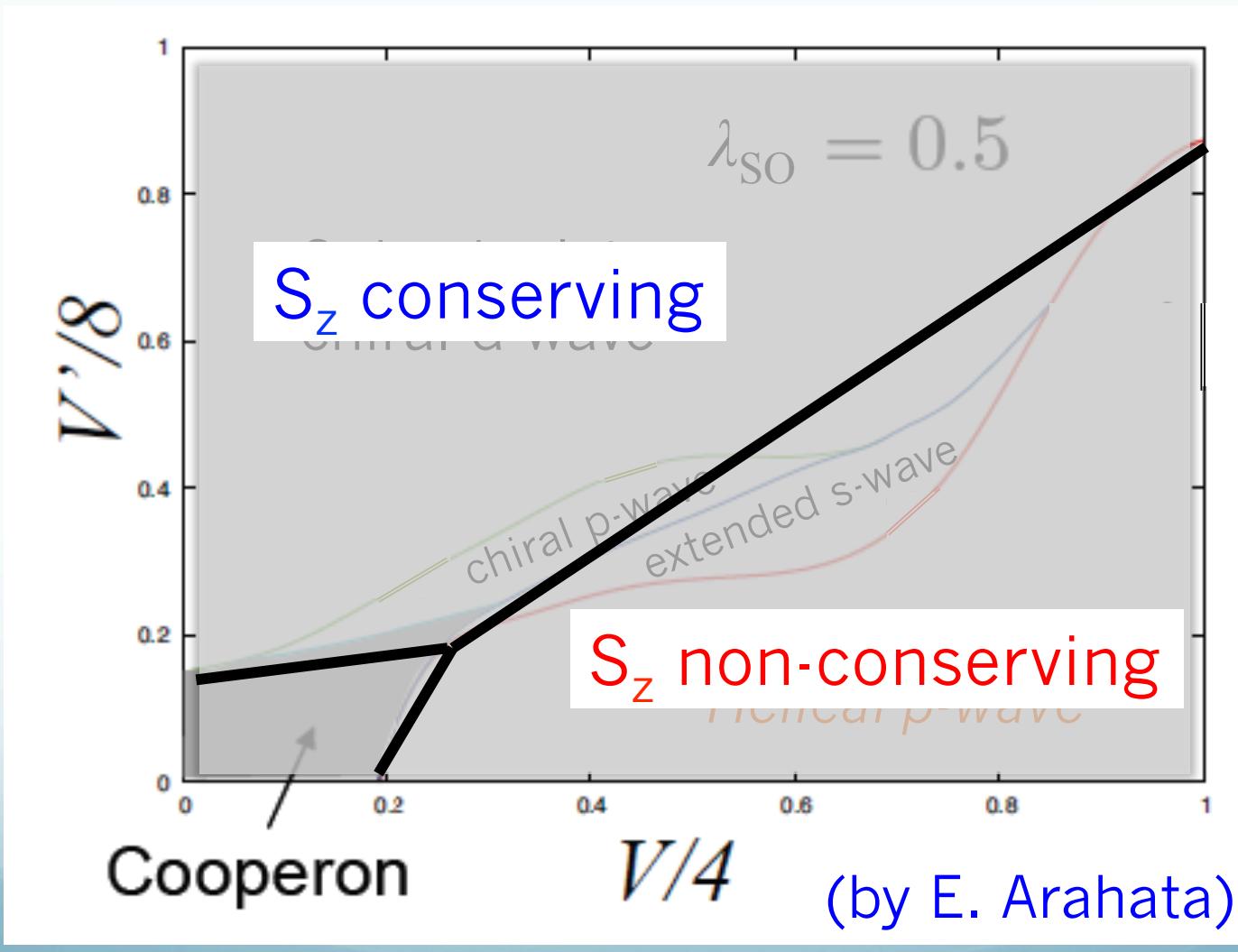
i.e. spin-triplet pairing



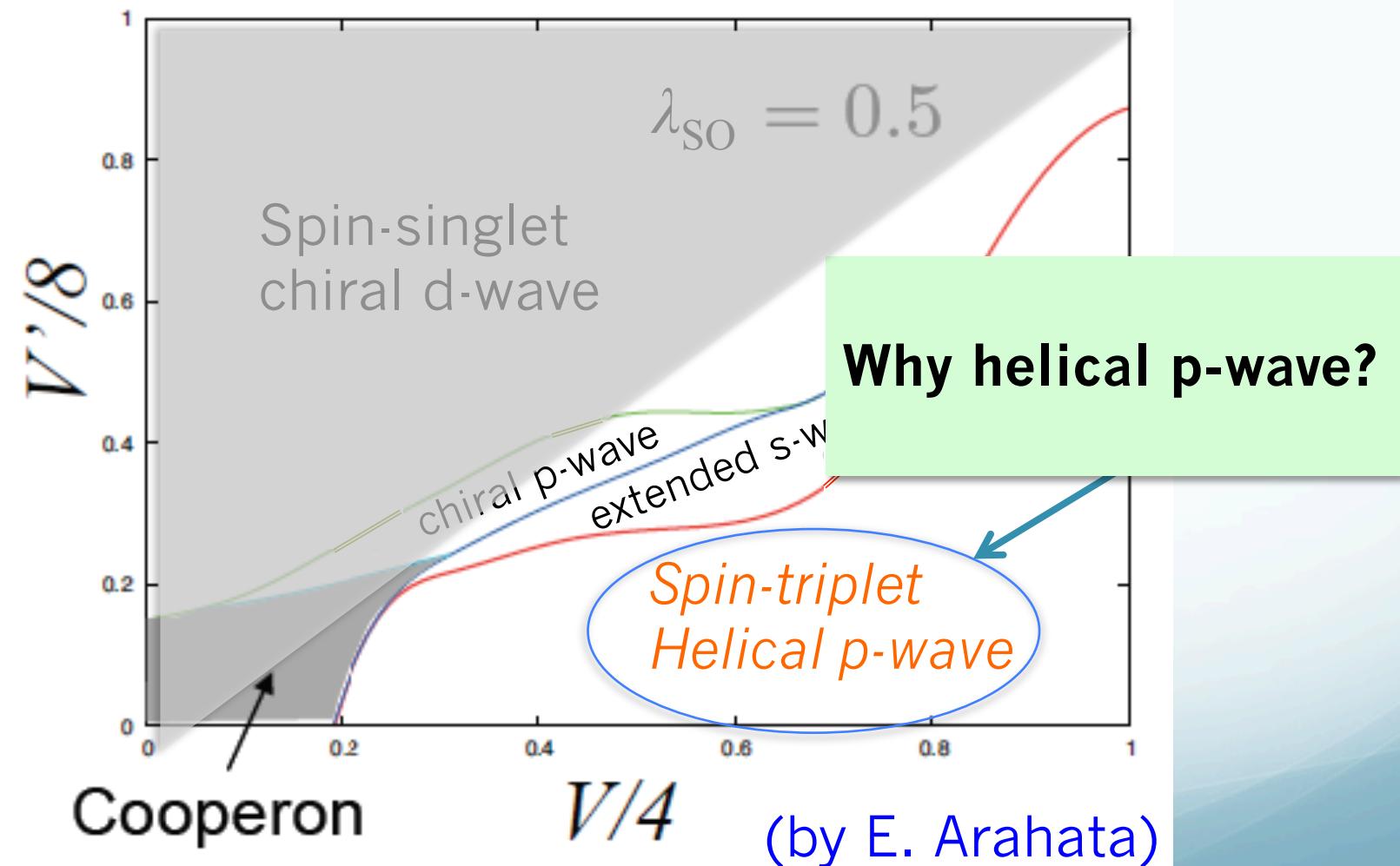
NNN (intra-sublattice) attraction
 \Rightarrow spin-unpolarized components of pairing

i.e. spin-singlet or S_z -conserving spin-triplet pairing

Analysis of the Gap equation at T=0 with $V(NN)$ and $V'(NNN)$ attraction



Analysis of the Gap equation at T=0 with $V(NN)$ and $V'(NNN)$ attraction

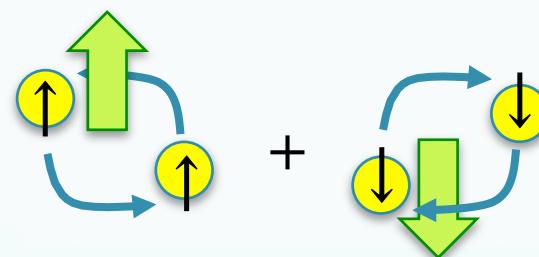


Helical p-wave state

2D version of superfluid $^3\text{He-B}$

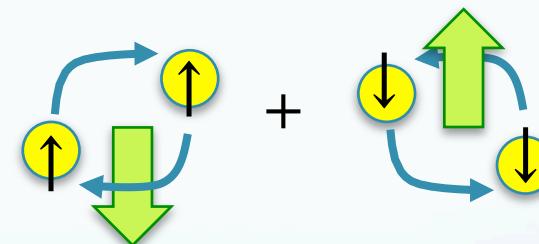
- Superposition of Kramers states
- Two different helicity (parallel and anti-parallel)

Paralell



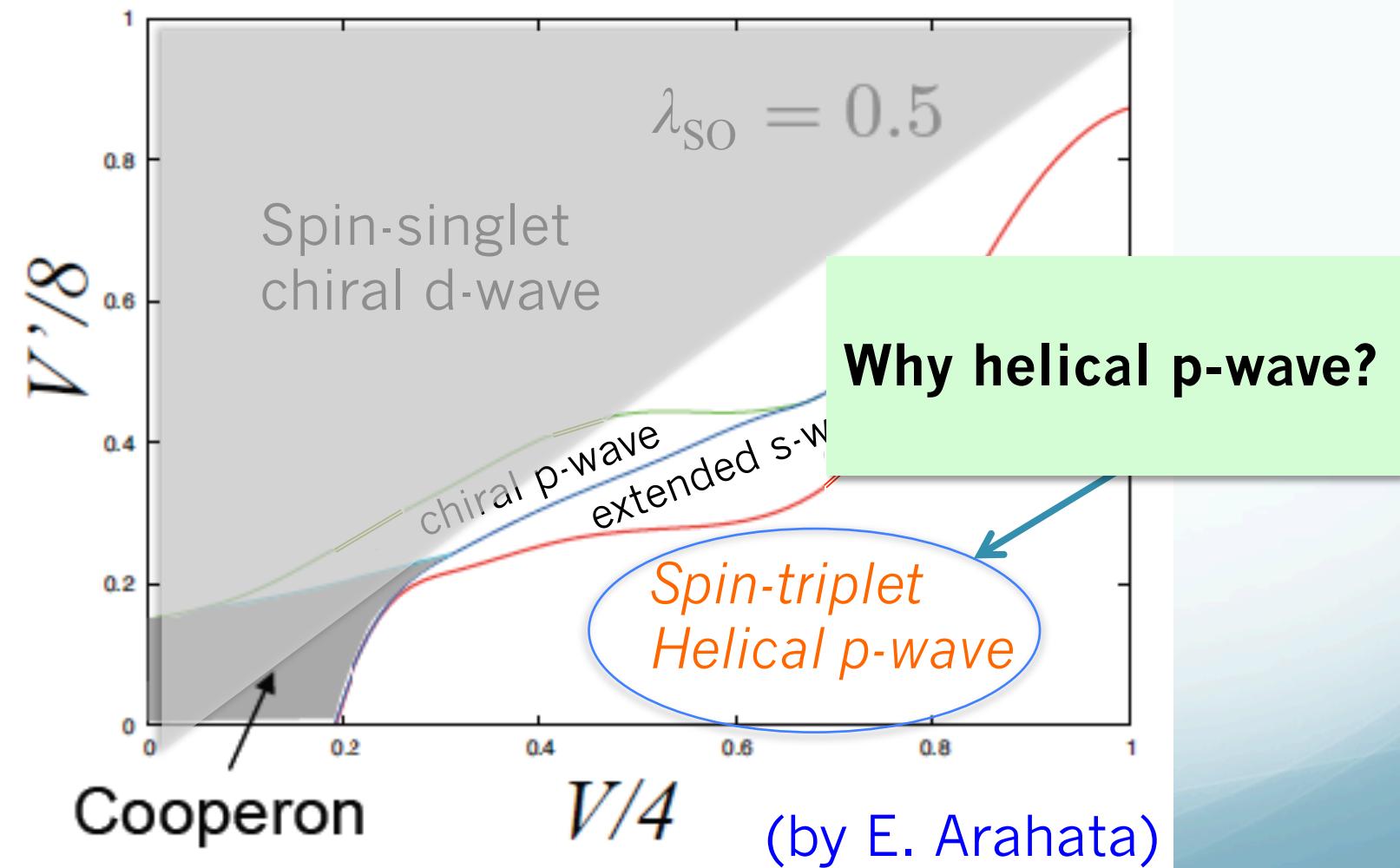
$$L_z = 1 \\ S_z = 1$$

Anti-parallel



$$L_z = -1 \\ S_z = 1$$

Analysis of the Gap equation at T=0 with $V(NN)$ and $V'(NNN)$ attraction



Diagonalization of H_{KM}

$$\begin{pmatrix} A_{\mathbf{k}s} \\ B_{\mathbf{k}s} \end{pmatrix} = U(\mathbf{k}, s) \begin{pmatrix} \tilde{c}_{\mathbf{k}s} \\ \tilde{d}_{\mathbf{k}s} \end{pmatrix}$$

Ann. of conduction electron

Ann. of valence electron

$$U(\mathbf{k}, \uparrow) = \begin{pmatrix} u_{\mathbf{k}} e^{i\phi_{\mathbf{k}}} & v_{\mathbf{k}} \\ v_{\mathbf{k}} & -u_{\mathbf{k}} e^{-i\phi_{\mathbf{k}}} \end{pmatrix}$$

$$U(\mathbf{k}, \downarrow) = \begin{pmatrix} v_{\mathbf{k}} & u_{\mathbf{k}} e^{i\phi_{\mathbf{k}}} \\ u_{\mathbf{k}} e^{-i\phi_{\mathbf{k}}} & -v_{\mathbf{k}} \end{pmatrix}$$

Pairing Hamiltonian for the conduction electrons with NN interaction

$$H^c = H_{\text{KM}}^c + H_{nn}^c$$

$$H_{\text{KM}}^c = \sum_{\mathbf{k}s} \epsilon_{\mathbf{k}} \tilde{c}_{\mathbf{k}s}^\dagger \tilde{c}_{\mathbf{k}s}$$

$$H_{nn}^c = -\frac{V}{M} \sum_{\mathbf{k}, \mathbf{p}} t_{\mathbf{k}-\mathbf{p}} e^{-i(\phi_{\mathbf{k}} - \phi_{\mathbf{p}})} u_{\mathbf{k}}^2 u_{\mathbf{p}}^2 \tilde{c}_{\mathbf{k}\uparrow}^\dagger \tilde{c}_{-\mathbf{k}\uparrow}^\dagger \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow}$$

$$-\frac{V}{M} \sum_{\mathbf{k}, \mathbf{p}} t_{\mathbf{k}-\mathbf{p}} e^{-i(\phi_{\mathbf{k}} - \phi_{\mathbf{p}})} v_{\mathbf{k}}^2 v_{\mathbf{p}}^2 \tilde{c}_{\mathbf{k}\downarrow}^\dagger \tilde{c}_{-\mathbf{k}\downarrow}^\dagger \tilde{c}_{-\mathbf{p}\downarrow} \tilde{c}_{\mathbf{p}\downarrow}$$

$$-\frac{V}{M} \sum_{\mathbf{k}, \mathbf{p}} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{k}} v_{\mathbf{k}} u_{\mathbf{p}} v_{\mathbf{p}} \tilde{c}_{\mathbf{k}\uparrow}^\dagger \tilde{c}_{-\mathbf{k}\downarrow}^\dagger \tilde{c}_{-\mathbf{p}\downarrow} \tilde{c}_{\mathbf{p}\uparrow}$$

$$-\frac{V}{M} \sum_{\mathbf{k}, \mathbf{p}} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{k}} v_{\mathbf{k}} u_{\mathbf{p}} v_{\mathbf{p}} \tilde{c}_{\mathbf{k}\downarrow}^\dagger \tilde{c}_{-\mathbf{k}\uparrow}^\dagger \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\downarrow},$$

M: # of
the unit cells

BCS mean field (Gap matrix)

$$\hat{\Delta}(\mathbf{k}) = \begin{pmatrix} \Delta_{\uparrow\uparrow}(\mathbf{k}) & \Delta_{\uparrow\downarrow}(\mathbf{k}) \\ \Delta_{\downarrow\uparrow}(\mathbf{k}) & \Delta_{\downarrow\downarrow}(\mathbf{k}) \end{pmatrix}$$

$$\Delta_{\uparrow\uparrow}(\mathbf{k}) = -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \sum_{\mathbf{p}} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{k}'\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle$$

$$\Delta_{\downarrow\downarrow}(\mathbf{k}) = -\frac{V}{M} v_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \sum_{\mathbf{p}} t_{\mathbf{k}-\mathbf{p}} v_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\downarrow} \tilde{c}_{\mathbf{p}\downarrow} \rangle$$

$$\Delta_{\uparrow\downarrow}(\mathbf{k}) = -\frac{V}{M} u_{\mathbf{k}} v_{\mathbf{k}} \sum_{\mathbf{p}} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{k}'} v_{\mathbf{p}} \langle \tilde{c}_{-\mathbf{p}\downarrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle$$

$$\Delta_{\downarrow\uparrow}(\mathbf{k}) = -\frac{V}{M} u_{\mathbf{k}} v_{\mathbf{k}} \sum_{\mathbf{p}} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{p}} v_{\mathbf{p}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\downarrow} \rangle$$

$$\Delta_{\uparrow\uparrow}(\mathbf{k}) = -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \sum_{\mathbf{p}} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle$$

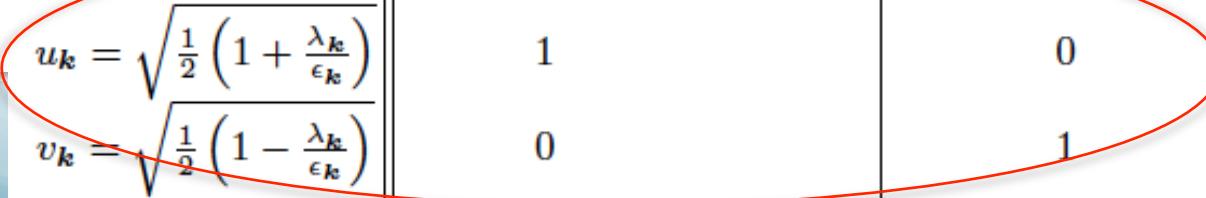
$$\mathbf{k}, \mathbf{p} \in K, K'$$

would be dominant

$$\begin{aligned}
\Delta_{\uparrow\uparrow}(\boldsymbol{k}) &= -\frac{V}{M} u_{\boldsymbol{k}}^2 e^{-i\phi_{\boldsymbol{k}}} \sum_{\boldsymbol{p}} t_{\boldsymbol{k}-\boldsymbol{p}} u_{\boldsymbol{p}}^2 e^{i\phi_{\boldsymbol{p}}} \langle \tilde{c}_{-\boldsymbol{p}\uparrow} \tilde{c}_{\boldsymbol{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\boldsymbol{k}}^2 e^{-i\phi_{\boldsymbol{k}}} \left\{ \sum_{\boldsymbol{p} \in K} + \sum_{\boldsymbol{p} \in K'} \right\} t_{\boldsymbol{k}-\boldsymbol{p}} u_{\boldsymbol{p}}^2 e^{i\phi_{\boldsymbol{p}}} \langle \tilde{c}_{-\boldsymbol{p}\uparrow} \tilde{c}_{\boldsymbol{p}\uparrow} \rangle
\end{aligned}$$

$$\begin{aligned}
\Delta_{\uparrow\uparrow}(\mathbf{k}) &= -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \sum_{\mathbf{p}} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \left\{ \sum_{\mathbf{p} \in K} + \sum_{\mathbf{p} \in K'} \right\} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle
\end{aligned}$$

	around K	around K'
$ t_{\mathbf{k}} $	0	0
$\phi_{\mathbf{k}} = \arg(t_{\mathbf{k}})$	$\tan^{-1} \frac{\delta k_y}{\delta k_x} \equiv \theta_{\delta \mathbf{k}}$	$-\tan^{-1} \frac{\delta k_y}{\delta k_x} = -\theta_{\delta \mathbf{k}}$
$\lambda_{\mathbf{k}}$	λ_{SO}	$-\lambda_{SO}$
$\epsilon_{\mathbf{k}}$	$ \lambda_{SO} $	$ \lambda_{SO} $
$u_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 + \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$ $v_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 - \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$		
	$\cdot \lambda_{SO} > 0$ 1 0	$\cdot \lambda_{SO} > 0$ 0 1



$$\begin{aligned}
\Delta_{\uparrow\uparrow}(\boldsymbol{k}) &= -\frac{V}{M} u_{\boldsymbol{k}}^2 e^{-i\phi_{\boldsymbol{k}}} \sum_{\boldsymbol{p}} t_{\boldsymbol{k}-\boldsymbol{p}} u_{\boldsymbol{p}}^2 e^{i\phi_{\boldsymbol{p}}} \langle \tilde{c}_{-\boldsymbol{p}\uparrow} \tilde{c}_{\boldsymbol{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\boldsymbol{k}}^2 e^{-i\phi_{\boldsymbol{k}}} \left\{ \sum_{\boldsymbol{p} \in K} + \sum_{\boldsymbol{p} \in K'} \right\} t_{\boldsymbol{k}-\boldsymbol{p}} u_{\boldsymbol{p}}^2 e^{i\phi_{\boldsymbol{p}}} \langle \tilde{c}_{-\boldsymbol{p}\uparrow} \tilde{c}_{\boldsymbol{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\boldsymbol{k}}^2 e^{-i\phi_{\boldsymbol{k}}} \sum_{\boldsymbol{p} \in K} t_{\boldsymbol{k}-\boldsymbol{p}} u_{\boldsymbol{p}}^2 e^{i\phi_{\boldsymbol{p}}} \langle \tilde{c}_{-\boldsymbol{p}\uparrow} \tilde{c}_{\boldsymbol{p}\uparrow} \rangle
\end{aligned}$$

$$\begin{aligned}
\Delta_{\uparrow\uparrow}(\boldsymbol{k}) &= -\frac{V}{M} u_{\boldsymbol{k}}^2 e^{-i\phi_{\boldsymbol{k}}} \sum_{\boldsymbol{p}} t_{\boldsymbol{k}-\boldsymbol{p}} u_{\boldsymbol{p}}^2 e^{i\phi_{\boldsymbol{p}}} \langle \tilde{c}_{-\boldsymbol{p}\uparrow} \tilde{c}_{\boldsymbol{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\boldsymbol{k}}^2 e^{-i\phi_{\boldsymbol{k}}} \left\{ \sum_{\boldsymbol{p} \in K} + \sum_{\boldsymbol{p} \in K'} \right\} t_{\boldsymbol{k}-\boldsymbol{p}} u_{\boldsymbol{p}}^2 e^{i\phi_{\boldsymbol{p}}} \langle \tilde{c}_{-\boldsymbol{p}\uparrow} \tilde{c}_{\boldsymbol{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\boldsymbol{k}}^2 e^{-i\phi_{\boldsymbol{k}}} \sum_{\boldsymbol{p} \in K} t_{\boldsymbol{k}-\boldsymbol{p}} u_{\boldsymbol{p}}^2 e^{i\phi_{\boldsymbol{p}}} \langle \tilde{c}_{-\boldsymbol{p}\uparrow} \tilde{c}_{\boldsymbol{p}\uparrow} \rangle
\end{aligned}$$

$$\boldsymbol{k},\boldsymbol{p}\in K$$

$$t_{\boldsymbol{k}-\boldsymbol{p}} = \sum_{i=1}^3 t e^{i(\boldsymbol{k}-\boldsymbol{p})\cdot \boldsymbol{a}_i} \simeq t_{\boldsymbol{k}-\boldsymbol{p}\simeq\boldsymbol{0}} = 3t$$

$$\begin{aligned}
\Delta_{\uparrow\uparrow}(\mathbf{k}) &= -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \sum_{\mathbf{p}} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \left\{ \sum_{\mathbf{p} \in K} + \sum_{\mathbf{p} \in K'} \right\} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \sum_{\mathbf{p} \in K} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \sum_{\mathbf{p} \in K} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle
\end{aligned}$$

$$\begin{aligned}
\Delta_{\uparrow\uparrow}(\mathbf{k}) &= -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \sum_{\mathbf{p}} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \left\{ \sum_{\mathbf{p} \in K} + \sum_{\mathbf{p} \in K'} \right\} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \sum_{\mathbf{p} \in K} t_{\mathbf{k}-\mathbf{p}} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle \\
&\simeq -\frac{V}{M} u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}} \sum_{\mathbf{p} \in K} u_{\mathbf{p}}^2 e^{i\phi_{\mathbf{p}}} \langle \tilde{c}_{-\mathbf{p}\uparrow} \tilde{c}_{\mathbf{p}\uparrow} \rangle \\
&= \Delta u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}}
\end{aligned}$$

$$\Delta_{\uparrow\uparrow}(\mathbf{k}) =$$

\approx

\approx

\approx

$$= \Delta u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}}$$

	around K	around K'
$ t_{\mathbf{k}} $	0	0
$\phi_{\mathbf{k}} = \arg(t_{\mathbf{k}})$	$\tan^{-1} \frac{\delta k_y}{\delta k_x} \equiv \theta_{\delta\mathbf{k}}$	$-\tan^{-1} \frac{\delta k_y}{\delta k_x} = -\theta_{\delta\mathbf{k}}$
$\lambda_{\mathbf{k}}$	λ_{SO}	$-\lambda_{SO}$
$\epsilon_{\mathbf{k}}$	$ \lambda_{SO} $	$ \lambda_{SO} $
	$\cdot \lambda_{SO} > 0$	$\cdot \lambda_{SO} > 0$
$u_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 + \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$	1	0
$v_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 - \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$	0	1

$$\Delta_{\uparrow\uparrow}(\mathbf{k}) \rightarrow \begin{cases} \Delta e^{-i\theta_{\delta\mathbf{k}}} & (\mathbf{k} \in K) \\ 0 & (\mathbf{k} \in K') \end{cases}$$

	around K	around K'
$ t_{\mathbf{k}} $	0	0
$\phi_{\mathbf{k}} = \arg(t_{\mathbf{k}})$	$\tan^{-1} \frac{\delta k_y}{\delta k_x} \equiv \theta_{\delta \mathbf{k}}$	$-\tan^{-1} \frac{\delta k_y}{\delta k_x} = -\theta_{\delta \mathbf{k}}$
$\lambda_{\mathbf{k}}$	λ_{SO}	$-\lambda_{\text{SO}}$
$\epsilon_{\mathbf{k}}$	$ \lambda_{\text{SO}} $	$ \lambda_{\text{SO}} $
	$\cdot \lambda_{\text{SO}} > 0$	$\cdot \lambda_{\text{SO}} > 0$
$u_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 + \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$	1	0
$v_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 - \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$	0	1
$\Delta_{\uparrow\uparrow}(\mathbf{k}) \propto u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}}$	$\Delta e^{-i\theta_{\delta \mathbf{k}}}$	0
$\Delta_{\downarrow\downarrow}(\mathbf{k}) \propto v_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}}$	0	$\Delta e^{i\theta_{\delta \mathbf{k}}}$
$\Delta_{\uparrow\downarrow}(\mathbf{k}) \propto u_{\mathbf{k}} v_{\mathbf{k}}$	0	0
$\Delta_{\downarrow\uparrow}(\mathbf{k}) \propto u_{\mathbf{k}} v_{\mathbf{k}}$	0	0

	around K		around K'
$ t_{\mathbf{k}} $	0		0
$\phi_{\mathbf{k}} = \arg(t_{\mathbf{k}})$	$\tan^{-1} \frac{\delta k_y}{\delta k_x} \equiv \theta_{\delta \mathbf{k}}$		$-\tan^{-1} \frac{\delta k_y}{\delta k_x} = -\theta_{\delta \mathbf{k}}$
$\lambda_{\mathbf{k}}$	λ_{SO}		$-\lambda_{\text{SO}}$
$\epsilon_{\mathbf{k}}$	$ \lambda_{\text{SO}} $		$ \lambda_{\text{SO}} $
		$\cdot \lambda_{\text{SO}} > 0$	$\cdot \lambda_{\text{SO}} < 0$
$u_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 + \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$	1	0	0
$v_{\mathbf{k}} = \sqrt{\frac{1}{2} \left(1 - \frac{\lambda_{\mathbf{k}}}{\epsilon_{\mathbf{k}}} \right)}$	0	1	1
$\Delta_{\uparrow\uparrow}(\mathbf{k}) \propto u_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}}$	$\Delta e^{-i\theta_{\delta \mathbf{k}}}$	0	0
$\Delta_{\downarrow\downarrow}(\mathbf{k}) \propto v_{\mathbf{k}}^2 e^{-i\phi_{\mathbf{k}}}$	0	$\Delta e^{-i\theta_{\delta \mathbf{k}}}$	$\Delta e^{i\theta_{\delta \mathbf{k}}}$
$\Delta_{\uparrow\downarrow}(\mathbf{k}) \propto u_{\mathbf{k}} v_{\mathbf{k}}$	0	0	0
$\Delta_{\downarrow\uparrow}(\mathbf{k}) \propto u_{\mathbf{k}} v_{\mathbf{k}}$	0	0	0

$$\lambda_{\text{SO}} > 0$$

$$e^{-i\theta_{\delta k}} |\uparrow\uparrow\rangle + e^{i\theta_{\delta k}} |\downarrow\downarrow\rangle$$

winding -2π

$S_z = +1$

winding $+2\pi$ (anti-parallel helicity)

$S_z = -1$

$$\lambda_{\text{SO}} < 0$$

$$e^{i\theta_{\delta k}} |\uparrow\uparrow\rangle + e^{-i\theta_{\delta k}} |\downarrow\downarrow\rangle$$

(parallel helicity)

winding 2π

$S_z = +1$

winding -2π

$S_z = -1$

- similar to the helical p-wave state, and would induce it
- helicity depends on $\text{sgn } \lambda_{\text{SO}}$ (=spin Chern # of ins. phase)

Summary

The role of the KM spin-orbit coupling for the Cooperon condensation

- S_z -conserving state is stabilized by NNN (intra-sublattice) int.
- S_z -non-conserving helical p-wave state is stabilized by NN (inter-sublattice) int.
- Helicity of helical p-wave state would be related to the band topology (i.e. spin Chern #) induced by the SOC