

odd interactions in quantum magnets and liquids



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U.S. DEPARTMENT OF
ENERGY

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Science

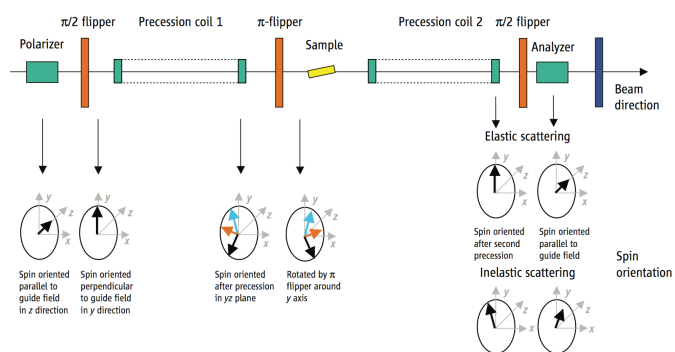


plan

- I. ^4He , roton
- II. XY magnets, lifetime



neutron-scattering spin-echo

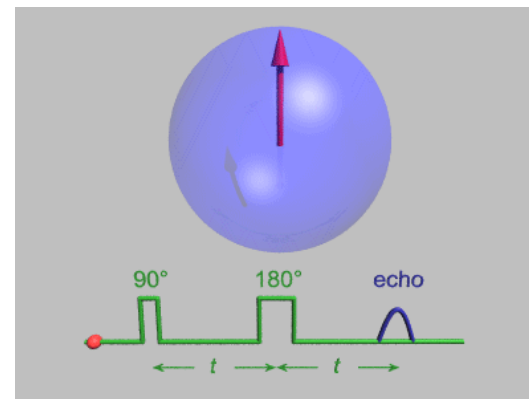


APPLIED PHYSICS 1888 VOL 312 SCIENCE 2006

The Neutron Spin-Echo Technique at Full Strength

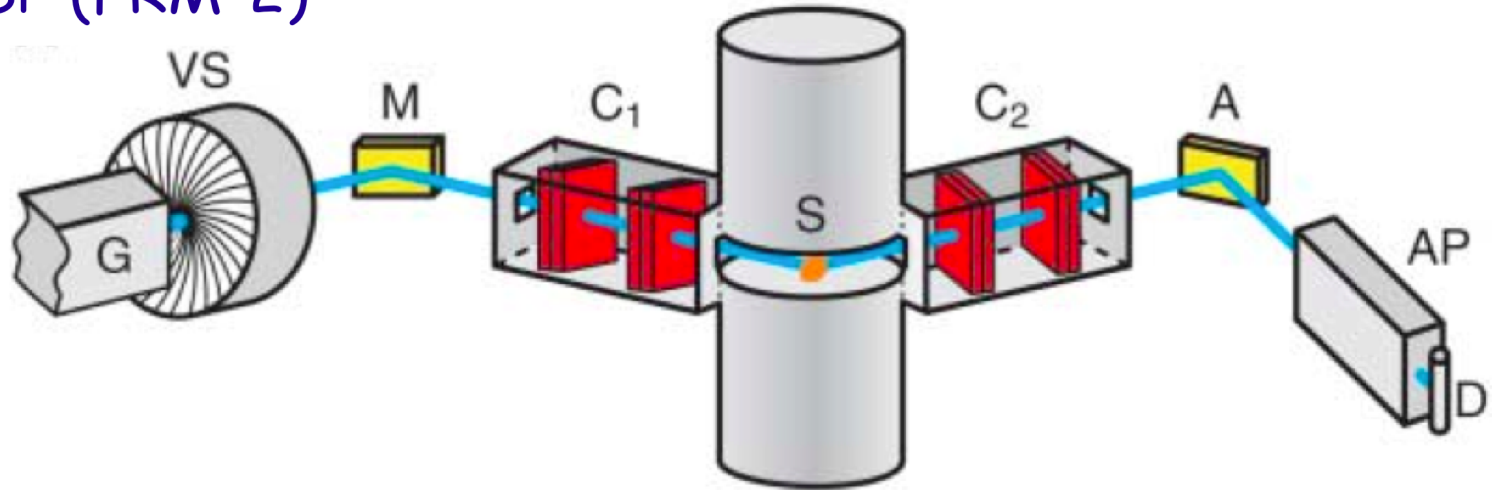
Joël Mesot

- resolution $\sim 1 \mu\text{eV}$ ($\approx 0.01\text{K}$)!

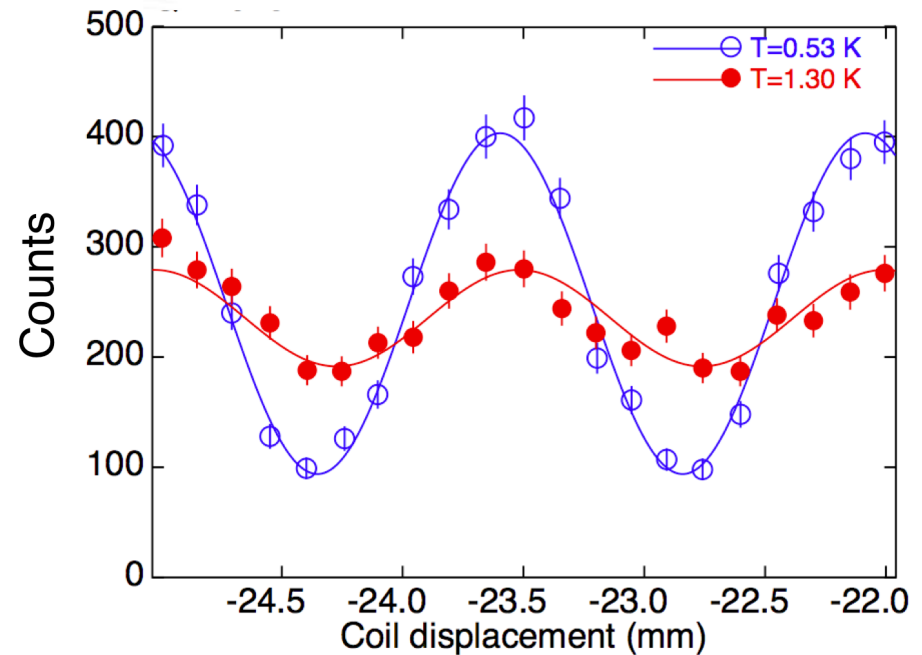


Neutron Resonant Spin Echo

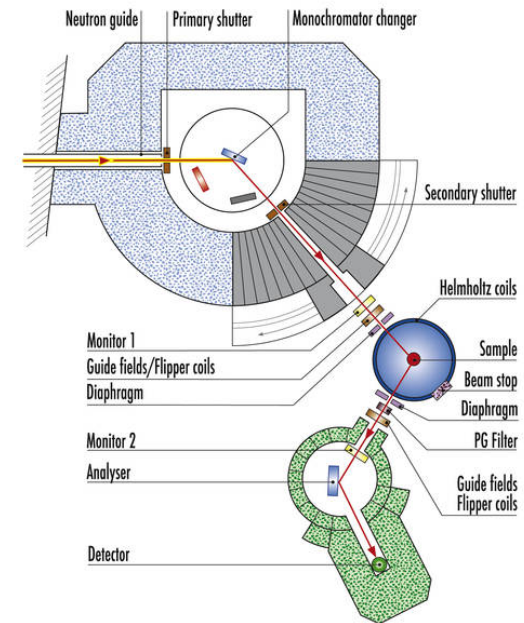
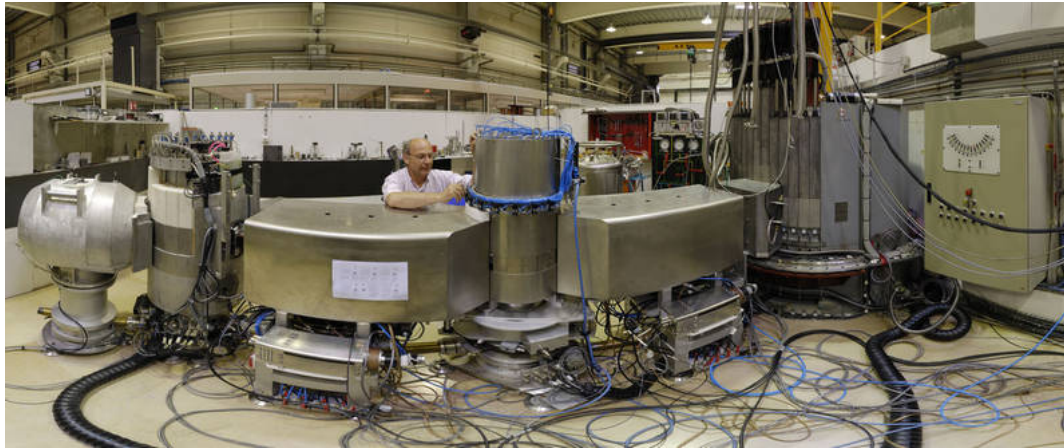
TRISP (FRM-2)



NRSE:
Line widths of dispersive modes
Energy shifts (phase stability)



ILL, IN22, and ZETA



Thermal neutron three-axis spectrometer with polarization analysis IN22

IN22 is a three-axis spectrometer (CRG) equipped for full polarization analysis. The option CRYOPAD and a 15 Tesla cryomagnet are optimised for inelastic scattering. The option ZETA provides neutron resonance spin echo (NRSE).




NEUTRONS[®]
FOR SCIENCE
Institut Laue-Langevin

CRG three-axis spectrometer **IN22** with the **ZETA** resonant neutron spin-echo setup



I. odd interactions in superfluid ^4He



experiments: Björn Fåk (CEA, ILL)
Thomas Keller (Munich, Stuttgart)

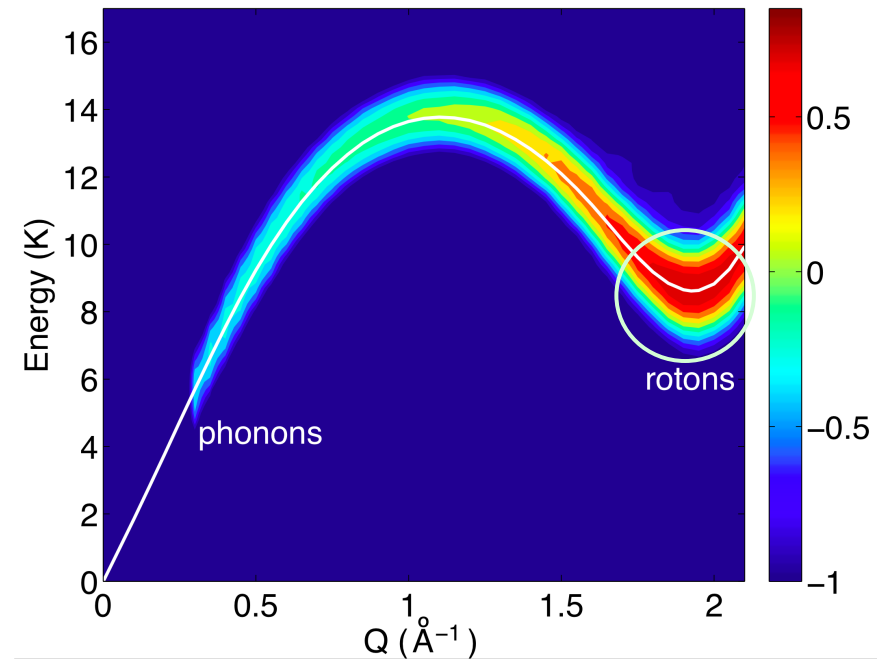
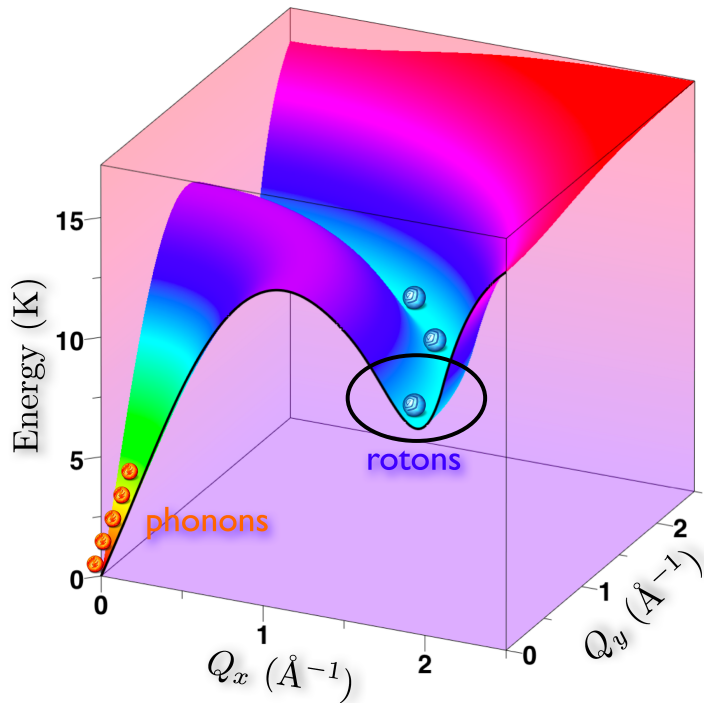
theory: Mike Zhitomirsky (CEA)
Sasha Chernyshev (UC Irvine)

plan

- I. history
- II. experiments
- III. theory/comparison
- IV. conclusions/outlook



spectrum of ^4He

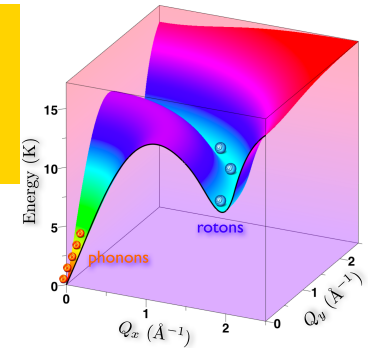


K. H. Andersen *et al.*, J. Phys. Cond. Mat. **6**, 821 (1994).

- T -dependence of roton's:
 - lifetime
 - energy

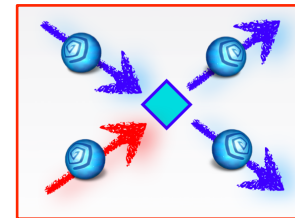


(long) history ...

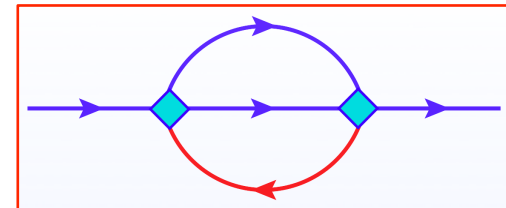


- L. D. Landau and I. M. Khalatnikov, *The theory of the viscosity of helium II: I. Collisions of elementary excitations in helium II*, Zh. Eksp. Teor. Fiz. **19**, 637- 650 (1949).
- roton lifetime (linewidth)

$$\Gamma(T) \propto N_r(T) \propto \sqrt{T} e^{-\frac{\Delta(T)}{T}}$$



- K. Bedell, D. Pines, and A. Zawadowskii, PRB **29**, 102 (1984).
- prefactors + energy shift



$$\delta(T) = \Delta(T) - \Delta_0 \propto N_r(T) \propto \sqrt{T} e^{-\frac{\Delta(T)}{T}}$$

- * **no** three-boson interaction needed/directly involved
- ** Hartree-term gives the same for $\delta(T)$



neutron-scattering spin-echo, ($>1.2\text{K}$)

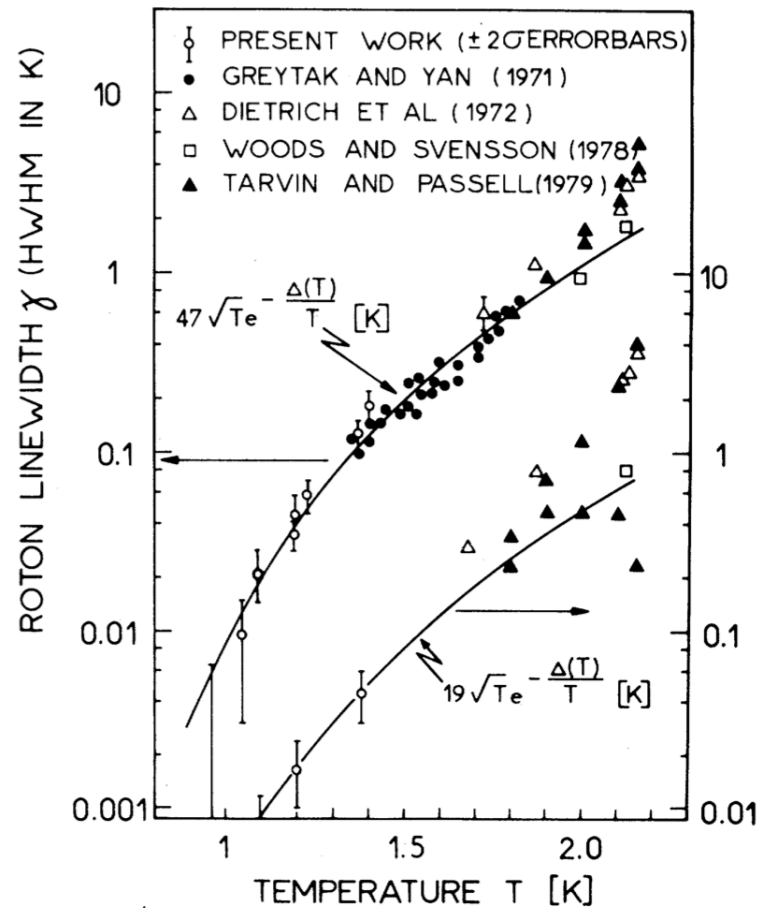
VOLUME 44, NUMBER 24

PHYSICAL REVIEW LETTERS

16 JUNE 1980

High-Resolution Study of Excitations in Superfluid ^4He by the Neutron Spin-Echo Technique

F. Mezei



- LK-theory fits:
 - quantitative



neutron-scattering spin-echo, ($>0.88\text{K}$)

VOLUME 77, NUMBER 19

PHYSICAL REVIEW LETTERS

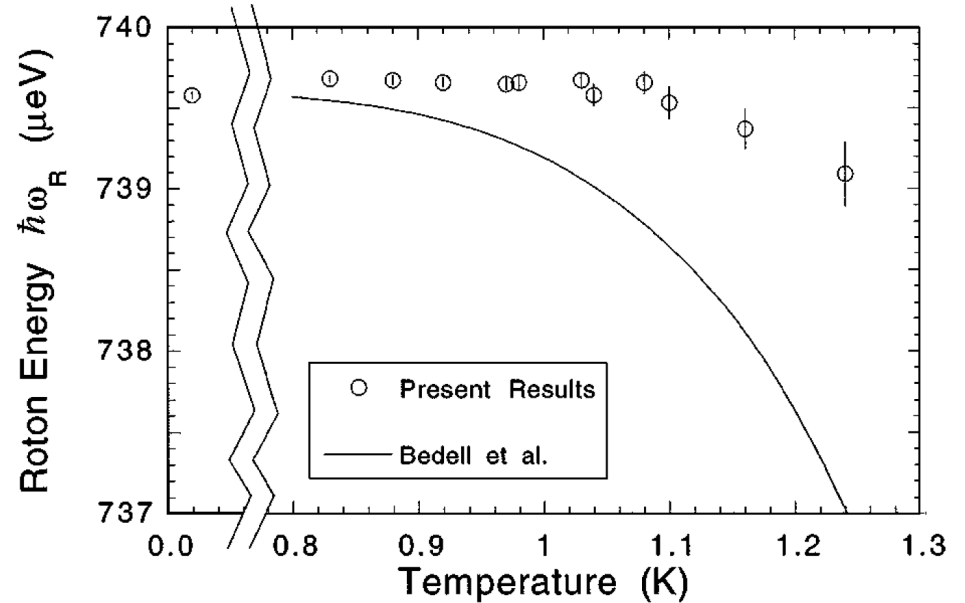
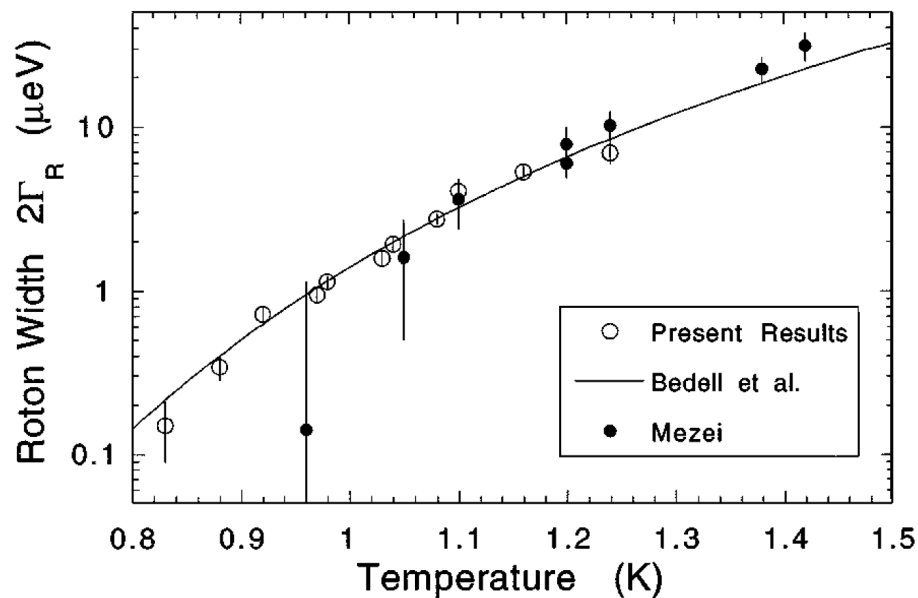
4 NOVEMBER 1996

High-Resolution Measurements of Rotons in ^4He

K. H. Andersen, J. Bossy, J. C. Cook, O. G. Randl, and J.-L. Ragazzoni

Institut Laue-Langevin, B.P. 156, 38042 Grenoble Cedex 9, France

(Received 22 March 1996)



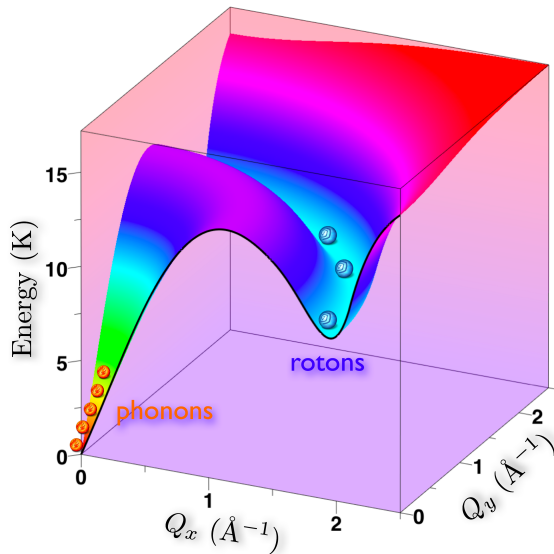
- linewidth:
LK-theory continues to work

$$\Gamma(T) \propto N_r(T) \propto \sqrt{T} e^{-\frac{\Delta(T)}{T}}$$

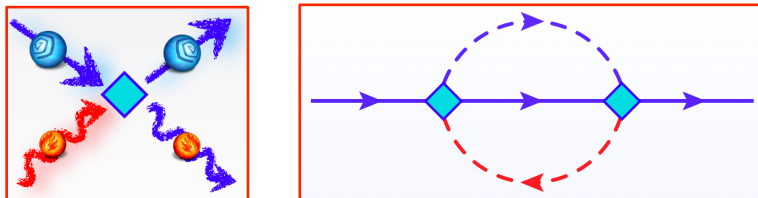
et al. and (2) that the temperature dependence of the roton energy cannot be described by the same theory in this temperature range. In fact, the rather good agreement between theory and experiment previously shown by the NSE technique [3], is now seen to break down below $T = 1.2$ K, while the theory is expected to work best in the $T \rightarrow 0$ limit.



phonons ?



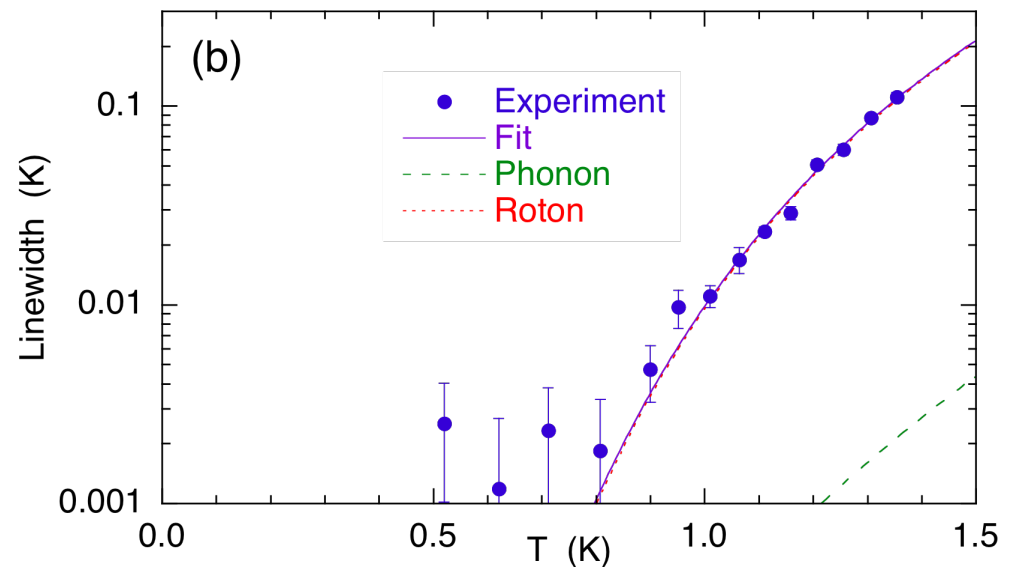
- where are the phonons ?
- roton-*phonon* scattering:



$$\left. \begin{array}{l} \Gamma(T) \\ \delta(T) \end{array} \right\} \propto T^3 \times T^4 = T^7 \left(= \tilde{A} \cdot \frac{T^7}{c^7} \right)$$

population \times Rayleigh scattering

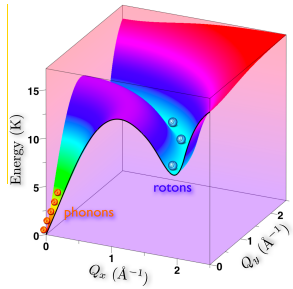
- using numbers: very small effect



- linewidth: LK-theory is fine !

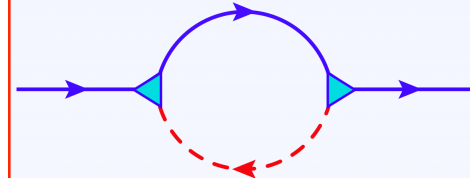
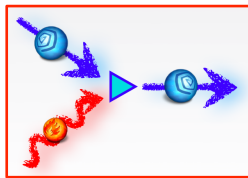
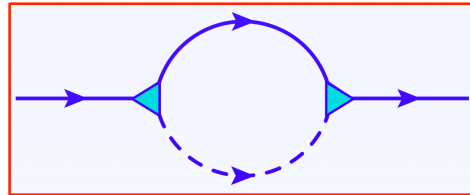
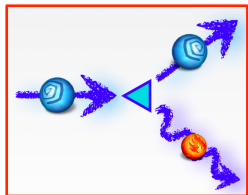
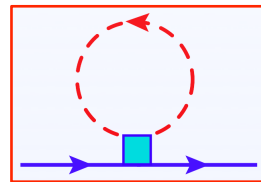
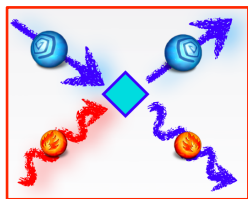
$$\Gamma(T) \propto N_r(T) \propto \sqrt{T} e^{-\frac{\Delta(T)}{T}}$$

* (= r-r at \$T=0.5\$K, where both \$\sim 10^{-6}\$ K)



phonons to the rescue

- **more** roton-*phonon* scatterings:



Hartree

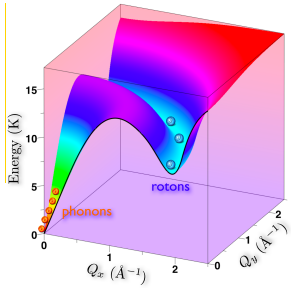
decay

coalescence

$$\left. \begin{array}{l} \text{Hartree} \\ \text{decay} \\ \text{coalescence} \end{array} \right\} \Rightarrow \Gamma \equiv 0 \Rightarrow \delta(T) \propto T^4$$

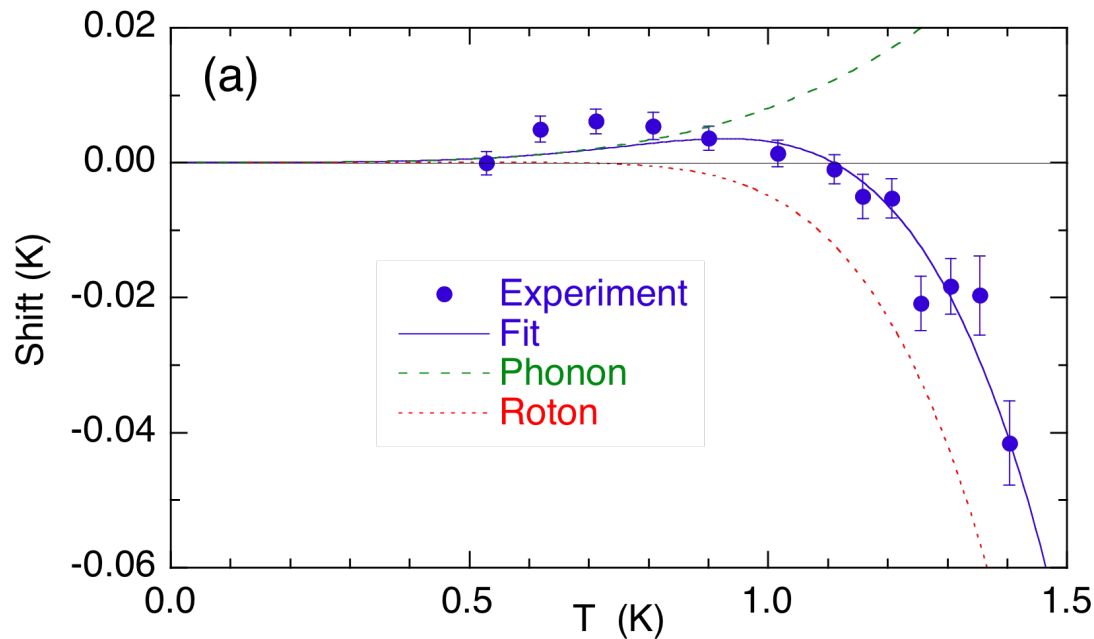
- three-particle (decay, coalescence) are necessarily positive





theory vs experiment

$$\delta(T) = -\delta_r \sqrt{T} \left(1 + \alpha \sqrt{T} \right) e^{-\Delta(T)/T} + \delta_{ph} T^4$$



- theory:
 - quantitatively (!) explains positive shift at low T
 - provides reasonable fit
 - three-particle interactions dominate the new effect



conclusions/outlook

- ✓ clear sign of the odd (3-particle) interactions in the roton energy shift
- ✓ LK ++ [LK-ZC (?)] theory is formulated
- ✓ neutron spin-echo allows to reach new regimes

- ☯ even lower T?
- ☯ phonon-phonon scattering, pressure dependence, etc.



II. **lifetime** of gapped excitations in (collinear) antiferromagnets

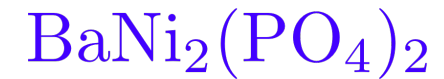
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University of California, Irvine



theory: Sasha Chernyshev (UC Irvine)
Mike Zhitomirsky (CEA)

experiments: Louis-Pierre Regnault,
Nicolas Martin (CEA, ILL)

material



L. P. REGNAULT AND J. ROSSAT-MIGNOD
**PHASE TRANSITIONS IN QUASI TWO-DIMENSIONAL
 PLANAR MAGNETS**

L. J. De Jongh (Ed.), Magnetic Properties of Layered Transition Metal Compounds 271–321.
 © 1990 Kluwer Academic Publishers. Printed in the Netherlands.

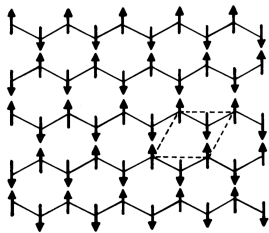
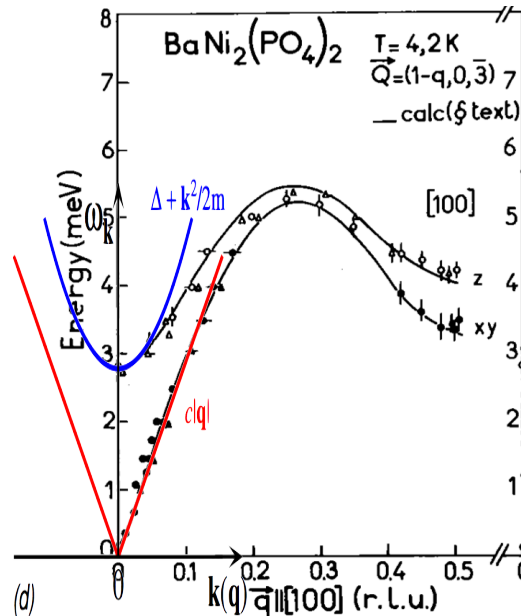
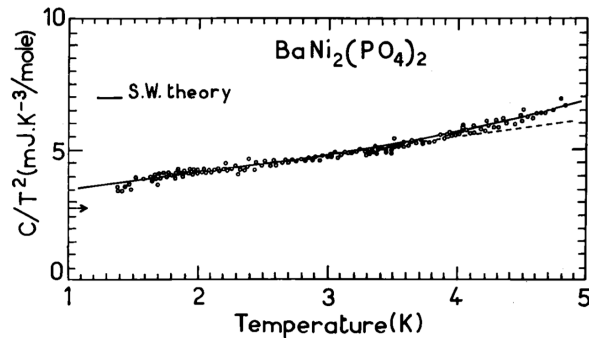
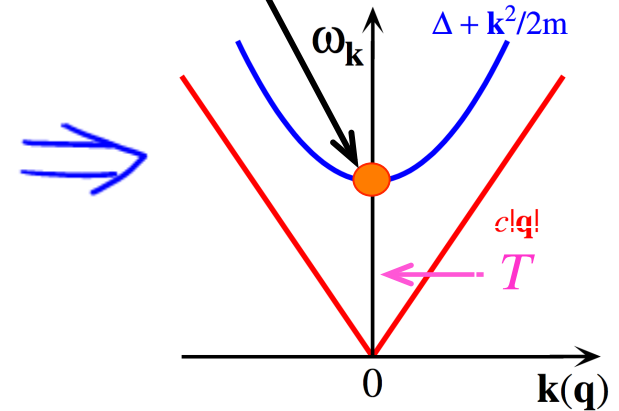


Fig. 3. Antiferromagnetic structure of $\text{BaNi}_2(\text{PO}_4)_2$.



$\Gamma = ???$

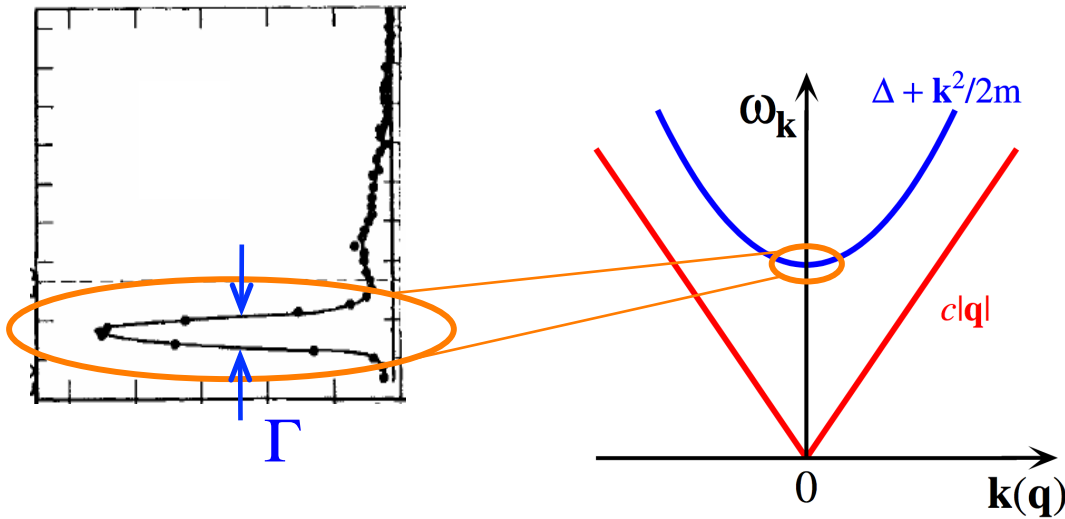


- $S = 1$, 2D, planar (XY)
- has gapped
- and acoustic modes

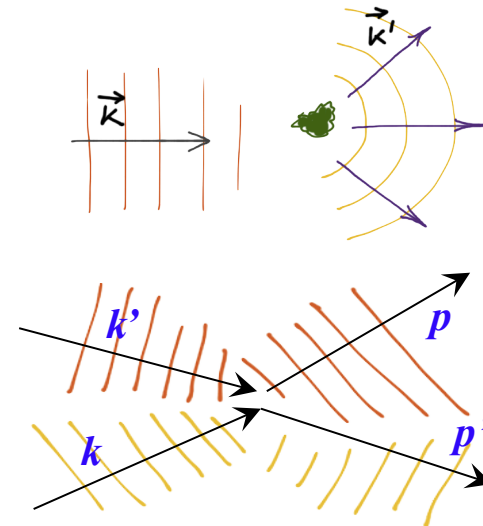
$$\mathcal{H} = \sum_{\langle i,j \rangle} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j + D \sum_i (S^z)^2$$



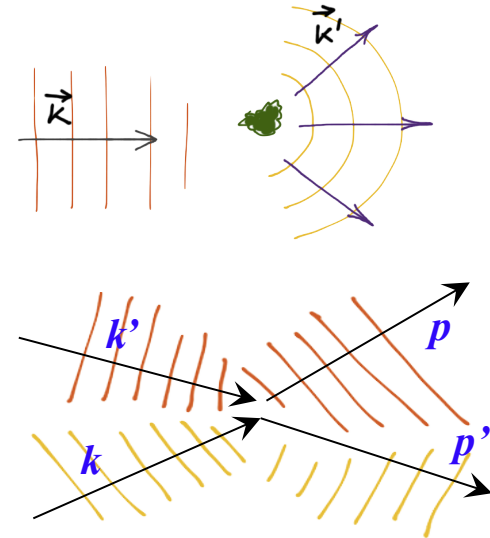
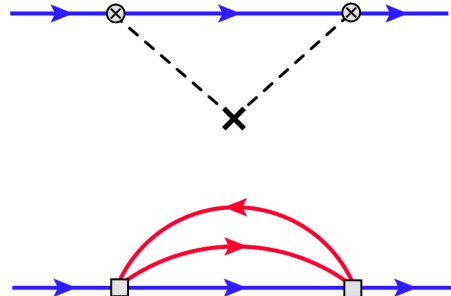
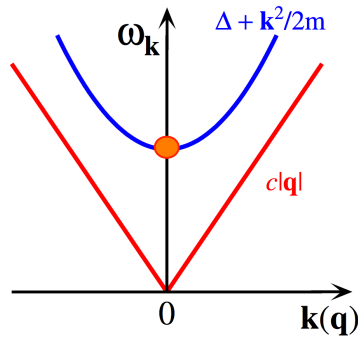
lifetime/linewidth/damping/decay rate



- lifetime⁻¹ = linewidth in "simple" AFs
- spin waves: scattering on?
 - impurities
 - themselves
 - combination of the two (impurity-assisted)



damping, theory expectations, I



- local distortions $\rightarrow \delta D, \delta J \rightarrow$ conventional impurity scattering (2D):

$$\Gamma_{\mathbf{k}}^{\text{imp}} \approx \Gamma_0 \propto n_{\text{imp}} \overline{\delta D}^2 \frac{m \omega_{\text{max}}^2}{\Delta^2}$$

- gapped on thermally excited gapless:

$$\Gamma_{\mathbf{k} \rightarrow 0}^{\text{m-m}} \approx \frac{\pi^3}{15} \frac{\tilde{g}^2}{c} \left(\frac{T}{c} \right)^5$$

- (and on gapped):

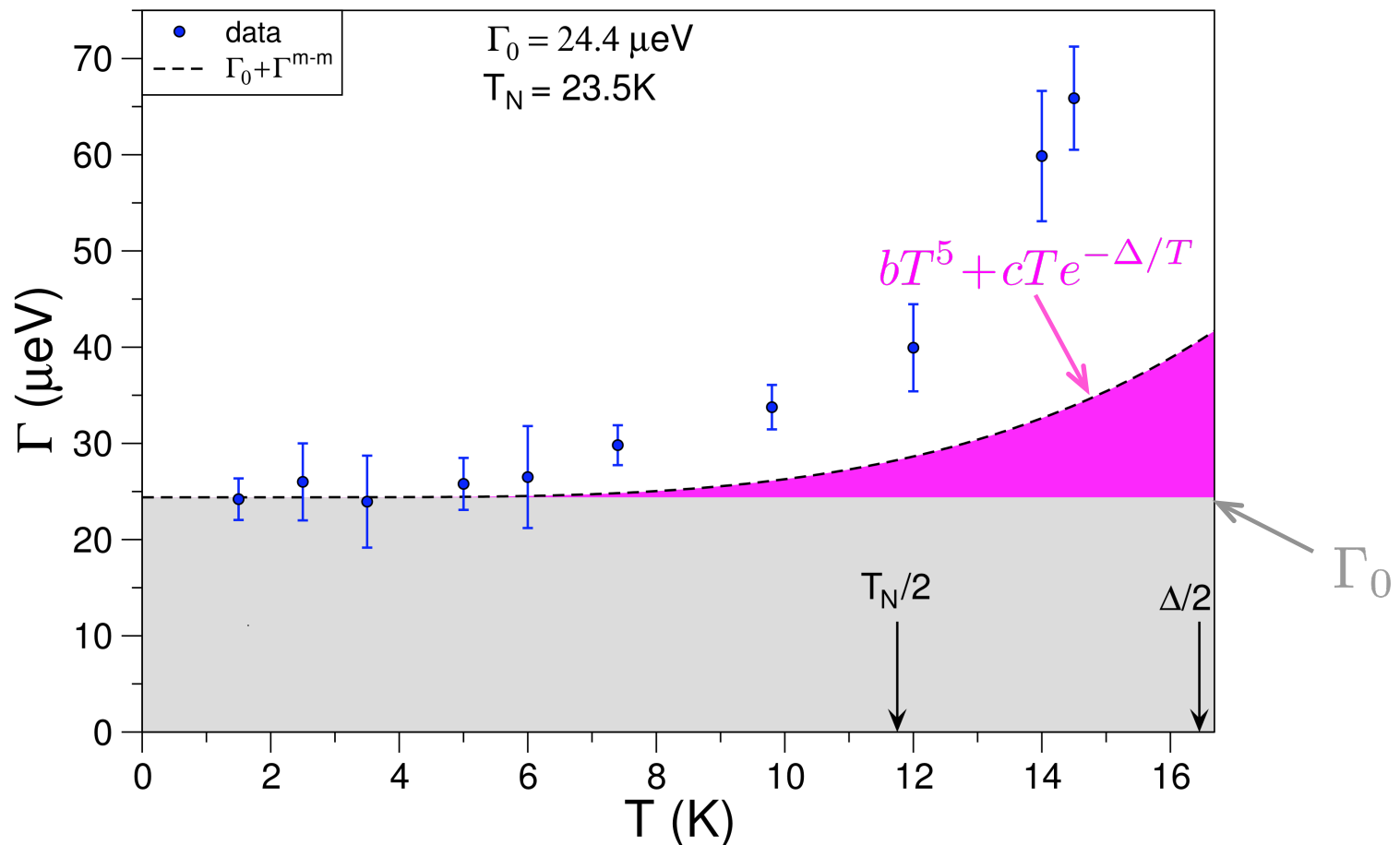
$$\Gamma_{\mathbf{k} \rightarrow 0}^{\beta\beta \rightarrow \beta\beta} \approx \frac{g_{\beta}^2 m^2 T}{4\pi} e^{-\Delta/T}$$

- numbers for m-m scattering are known/derivable!

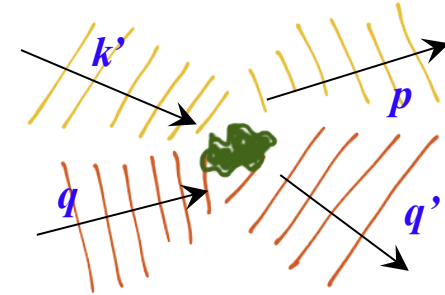
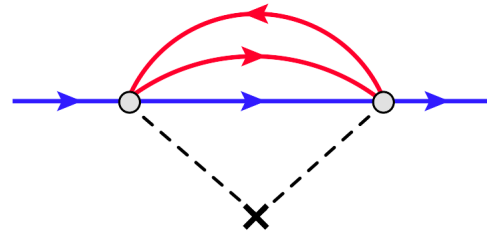
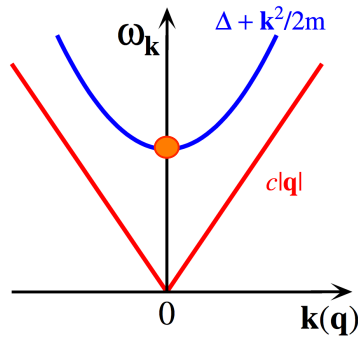


standard lore: $\rho = \rho_0^{\text{imp}} + \rho^{ee}(T)$

$$\Gamma = \underbrace{\Gamma_0}_{\text{imp}} + \underbrace{bT^5 + cTe^{-\Delta/T}}_{\text{m-m}}$$



damping, theory expectations, II



- impurity facilitates stronger m-m scattering

$$V_{\mathbf{k},\mathbf{q};\mathbf{k}',\mathbf{q}'}^{\text{imp}} \approx \tilde{g}_{\text{imp}} / \sqrt{\tilde{q}\tilde{q}'} \quad \text{vs} \quad V_{\mathbf{k}\mathbf{q};\mathbf{k}'\mathbf{q}'}^{\text{m-m}} \propto \sqrt{qq'}$$

- → lower power of T in Γ

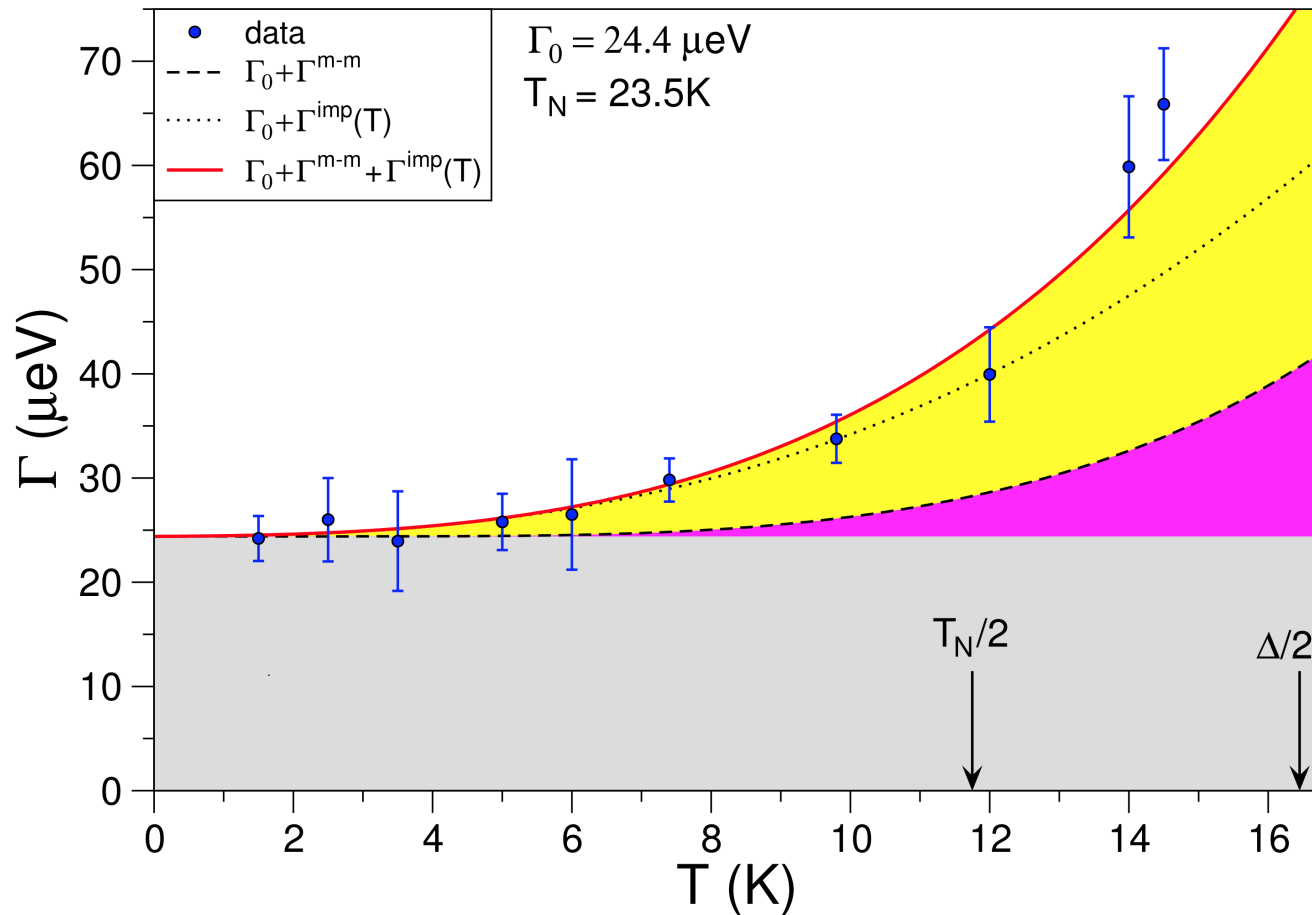
$$\Gamma_{\mathbf{k} \rightarrow 0}^{\text{imp},T} \approx \tilde{A} \left(\frac{T}{c} \right)^2 \left[\left(\ln \frac{T}{\omega_0} \right)^2 + \frac{\pi^2}{3} \right] \quad \tilde{A} \sim n_{\text{imp}} \overline{\delta D}^2 m$$

- reasons for “stronger” potential:
 - m-m interactions are singular but cancel out
 - impurities violate that cancellation
 - dynamically-induced “strong” disorder: optical spin-flip “sits” at impurity, scatters acoustic mode stronger ...

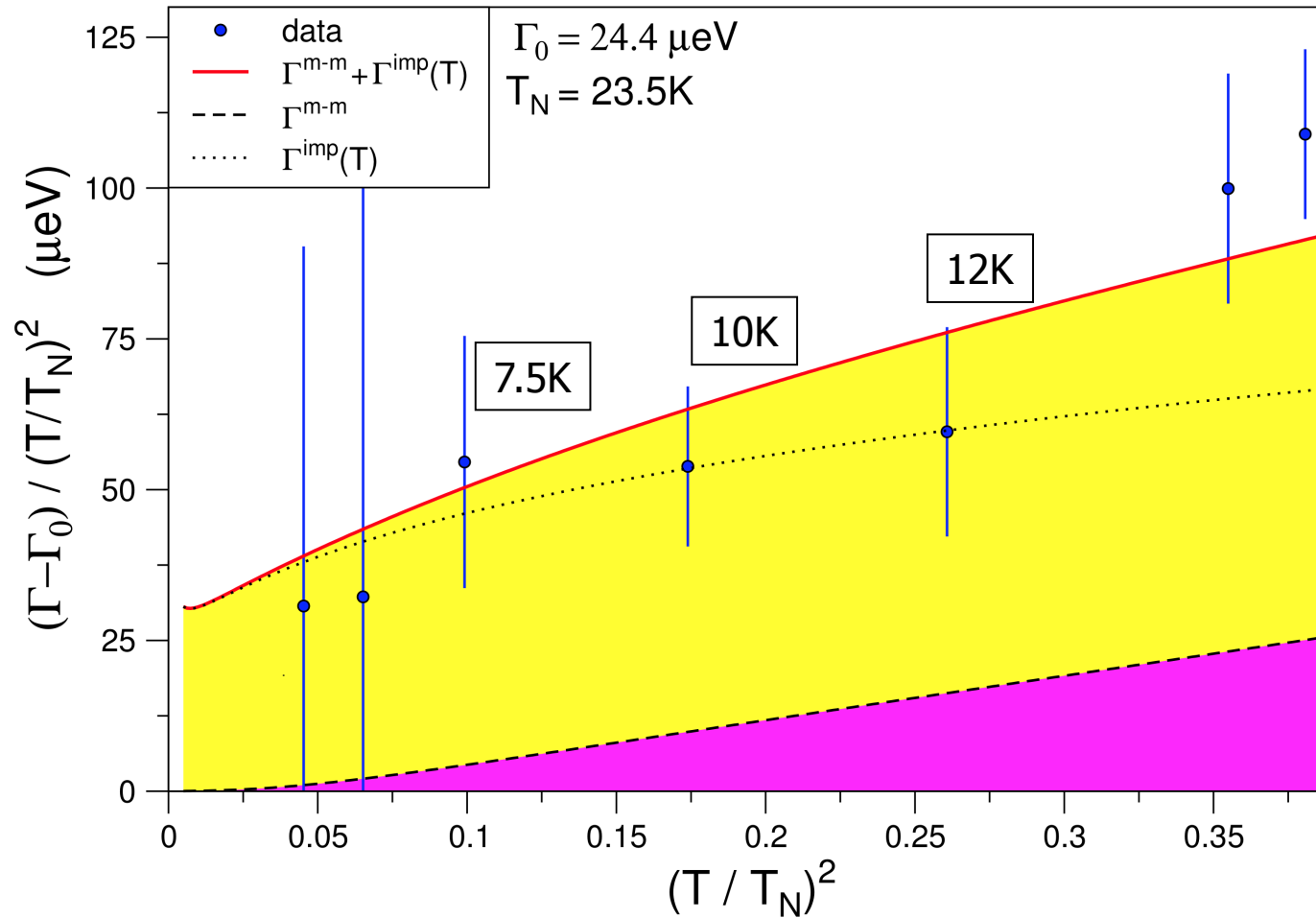


beyond the standard model ...

$$\Gamma = \underbrace{\Gamma_0}_{\text{imp}} + \underbrace{aT^2 \left[\left(\ln \frac{T}{\omega_0} \right)^2 + \frac{\pi^2}{3} \right]}_{\text{imp. finite-T}} + \underbrace{bT^5 + cTe^{-\Delta/T}}_{\text{m-m}}$$



better picture ...



cross-checks, predictions

$T=0$ and $T>0$ impurity terms must be related ($\Gamma_0 \sim \tilde{A} \sim n_{\text{imp}} \overline{\delta D}^2 m$)

✓ true, in our fit: $\Gamma_0 \approx \tilde{A} \approx 25 \mu\text{eV}$

does disorder strength make sense?

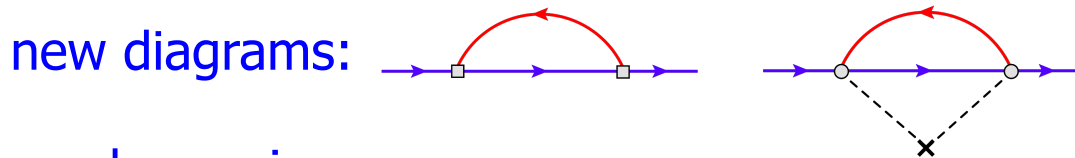
✓ estimate: $n_{\text{imp}} (\overline{\delta D}/D)^2 \approx \Gamma_0/\omega_{\text{max}} \approx 10^{-2}$

translates into a (very reasonable) statement that in $\text{BaNi}_2(\text{PO}_4)_2$, strong modulation of magnetic couplings of order 1 is spread over 1 in 100 unit cells

predictions:

➡ 3D: $\Gamma_{3D}^{\text{imp}} \propto n_{\text{imp}} T^{9/2}$

➡ AFs with non-collinear order \rightarrow 3-magnon coupling $\text{BaCo}_2(\text{AsO}_4)_2$



no change in m-m

lower power of T in impurity-induced:

$$\Gamma_{\mathbf{k} \rightarrow 0}^{\text{imp}, T} \approx \tilde{A}_3 \left(\frac{T}{c} \right) \ln \frac{T}{\omega_0} \quad \tilde{A}_3 \propto n_{\text{imp}} |\delta g_3|^2$$



conclusions

- ✓ general case: low- T lifetime of a magnetic excitation is completely dominated by the effects induced by a simple structural disorder
- ✓ support from experiments
- ✓ further predictions are made
- ✓ should be relevant to other systems

