

Studying BEC with a New Self-Consistent Approximation

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Introduction and Summary

We study properties of a homogeneous dilute Bose gas based on a self-consistent perturbation expansion that satisfies Noether's theorem and Goldstone's theorem simultaneously[1]. This formalism predicts that there should be a new class of Feynman diagrams for the self-energy characteristic of BEC that has been overlooked so far, which will be shown to modify standard results based on the Bogoliubov theory[2] substantially. First, these diagrams, which may be classified as "one-particle-reducible"(1PR), add an extra constant $c_{ip} \approx O(1)$ to the well-known expressions of the ground-state energy per particle E/N and condensate density n_0/n reported by Lee, Huang, and Yang[3] as

$$\frac{E}{N} = \frac{2\pi\hbar^2 a n}{m} \left[1 + \frac{16}{5} \left(\frac{8}{3\sqrt{\pi}} + c_{ip} \right) \sqrt{a^3 n} \right], \quad \frac{n_0}{n} = 1 - \left(\frac{8}{3\sqrt{\pi}} + c_{ip} \right) \sqrt{a^3 n}.$$

where a , n , and m are the s -wave scattering length, particle density, and particle mass, respectively.

Second, the lifetime of the one-particle excitation is also affected, as clarified by our calculation of the one-particle spectral function. It is shown that each excitation should have a finite lifetime proportional to the s -wave scattering length a , instead of a^2 for the normal state. Thus, the 1PR diagrams are predicted to change the nature of the one-particle excitation of BEC substantially from the Bogoliubov[2] mode with an infinite lifetime into a "bubbling" mode with a considerable decay rate.

[1] T. D. Lee *et al.*, Phys. Rev. **106** (1957) 1135. [2] T. Kita, PRB **80** (2009) 214502. [3] N. N. Bogoliubov, J. Phys. (USSR) **9** (1947) 23.

Conserving-Gapless theory

This theory is summarized below in two steps.

1. Construct the Luttinger-Ward functional Φ by including diagrams characteristic of condensed Bose system. (numerical weights of diagrams characteristic of BEC are left undetermined.)
2. Determine these weights by using Goldstone's theorem.

Self-energies

By carrying out functional differentiation for Φ , we obtain self-energies.

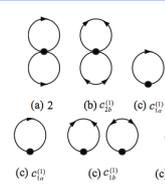
$$\Sigma(1,2) = -2\beta \frac{\delta\Phi}{\delta G(2,1)} \quad \Delta(1,2) = -2\beta \frac{\delta\Phi}{\delta F(2,1)} \quad 1 \equiv (\mathbf{r}_1, t_1)$$

1PR self-energy diagrams should be present.

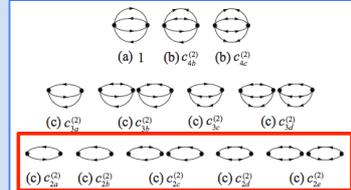


Diagrammatic structures of Φ

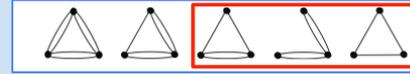
First-order



Second-order



Third-order (without arrows)



New class of Feynman diagrams

Dilute Bose gas system

Hamiltonian

$$H = \sum_{\mathbf{p}} (\epsilon_{\mathbf{p}} - \mu) c_{\mathbf{p}}^{\dagger} c_{\mathbf{p}} + \frac{g}{2} \sum_{\mathbf{p}_1 \mathbf{p}_2 \mathbf{q}} c_{\mathbf{p}_1 + \mathbf{q}}^{\dagger} c_{\mathbf{p}_2 - \mathbf{q}}^{\dagger} c_{\mathbf{p}_2} c_{\mathbf{p}_1}$$

First-order analysis

First-order self-energies

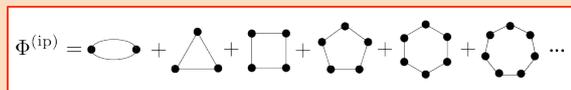


$$\begin{aligned} \Sigma^{(1)} &= 2gn \\ \Delta^{(1)} &= g \left(n - \sum_{\vec{p}} F_{\vec{p}}^{(1)} \right) \\ n &= n_0 - \sum_{\vec{p}} G_{\vec{p}}^{(1)} \\ \mu^{(1)} &= \Sigma^{(1)} - \Delta^{(1)} \end{aligned}$$

$$\begin{aligned} \frac{n_0}{n} &= 1 - \frac{8}{3\sqrt{\pi}} \sqrt{a^3 n} \\ \frac{E}{N} &= 4\pi a n \left(1 + \frac{128}{15\sqrt{\pi}} \sqrt{a^3 n} \right) \end{aligned}$$

Reproduce the results by Lee *et al.*

Structure of functional $\Phi^{(ip)}$ giving 1PR self-energies



→ These polygonal diagrams provide dominant contribution.

Self-energies

$$\begin{aligned} \Sigma(\vec{p}) &= \Sigma^{(1)} + \Delta_{im}(\vec{p}) \quad \Delta(\vec{p}) = \Delta^{(1)} + \Delta_{im}(\vec{p}) \\ \Delta_{im}(\vec{p}) &\equiv \sum_{n=2}^{\infty} \Delta_{im}^{(n)}(\vec{p}) \\ &= 2(gn_0)^2 (G(\vec{p}) + G(-\vec{p}) - 2F(\vec{p})) \\ &\quad + \frac{5}{4} (2gn_0)^3 (G(\vec{p}) + G(-\vec{p}) - 2F(\vec{p}))^2 + \dots \end{aligned}$$

→ These self-energies add an extra constant c_{ip} to results by Lee *et al.*, and give the one-particle excitation a finite lifetime proportional to a !

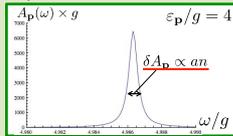
Applications of some approximations

1. one-particle spectral function [4]

$$A_{\mathbf{p}}(\omega) = -2 \text{Im} G_{11}(\vec{p})|_{i\epsilon \rightarrow \omega + i\eta}$$

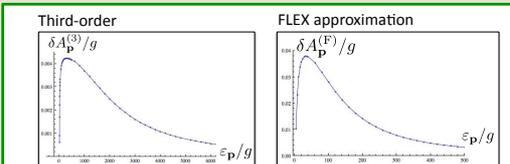
- It has a sharp peak in the one-particle excitation.
- The width of the peak $\delta A_{\mathbf{p}}$ corresponds to the lifetime of one-particle excitation.

*Spectral function of third-order



The width proportional to a !

Momentum dependence of $\delta A_{\mathbf{p}}$



→ lifetime of one-particle excitation changes with the own momentum!

2. Estimates of c_{ip} [5]

$$c_{ip} = \frac{4\sqrt{2}}{gn\pi^{3/2}} \int_0^{\infty} d\epsilon_p \epsilon_p^{1/2} \int_0^{\infty} d\epsilon_{\ell} \frac{(\epsilon_{\ell}^2 - \epsilon_p^2) \Delta_{im}(\vec{p})}{[\epsilon_{\ell}^2 + \epsilon_p(\epsilon_p + 2\Delta_{im}(\vec{p}))][\epsilon_{\ell}^2 + \epsilon_p(\epsilon_p + 2\Delta^{(1)})]}$$

Physical quantities containing the c_{ip}

$$\begin{aligned} \frac{E}{N} &= \frac{2\pi\hbar^2 a n}{m} \left[1 + \frac{16}{5} \left(\frac{8}{3\sqrt{\pi}} + c_{ip} \right) \sqrt{a^3 n} \right], \\ \frac{n_0}{n} &= 1 - \left(\frac{8}{3\sqrt{\pi}} + c_{ip} \right) \sqrt{a^3 n}. \end{aligned}$$

→ Third-order: $c_{ip} \approx 0.412$ FLEX: $c_{ip} \approx 0.563$

[4] K. T and T. Kita: J. Phys. Soc. Jpn. **83** (2014) 033001.

[5] K. T and T. Kita: J. Phys. Soc. Jpn. **82** (2013) 063001.