

phase approximation (RPA). The paramagnon spectral density function, $\rho_p(x, \omega) = \frac{1}{\pi} \Im \chi_\omega^R(x, x)$, is enhanced in the constriction and grows rapidly while $I_0 = I(\pi \hbar v_F)^{-1}$ approaches 1^- (the mean-field Stoner instability point for homogeneous system). The characteristic energy of the paramagnon mode behaves like $\hbar \omega_{para} \equiv \pi \sqrt{1 - I_0} \hbar v_F / L$.

Linear response theory: We examined the conductance in a linear response (Kubo) theory. [7] The DC conductance is obtained using current-current response function.

With a perturbation expansion only including RPA propagator in the lowest order as shown in the Feynman diagrams in Fig.1, we obtained the conductance formula: $G = G_0 \, d\eta(-\frac{\partial f(\eta)}{\partial \eta})(|t_\eta|^2 + \delta T^a(\eta) + \delta T^b(\eta))$, where the first term corresponds to the Landauer's formula and the other terms are including χ^R . The first correction δT^a corresponds to the current vertex correction and vanishes at $T = 0$. The second correction δT^b comes from the self-energy of single-particle Green's function. The imaginary part of the self-energy vanishes, however, there is a non-vanishing *positive* contribution at $T = 0$, which is a static potential renormalization.

The conductance of relatively higher temperatures is described by a model that the electron scattered by the fluctuating potential, $\delta U(x, t)$, which is due to *internal* dynamics (Coulomb interaction). By using the time-dependent perturbation theory [8] for $\delta U(x, t)$ we obtained the stationary current [6] and current noise [9] formula, which explain the experiments well.

Conclusions: We proposed a microscopic theory of electron scattering by spin fluctuations at a constriction in the lowest order perturbation as the origin of the *0.7 anomaly*, which explains successfully the higher temperature behavior. The low temperature and/or non-linear characteristics are interesting, however the lowest-order RPA treatment may not be enough to account for Kondo-like behavior. It would be quite interesting if an electron spin resonance or thermal/shot noise experiment could be done to detect the slow spin fluctuations.

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[References]

- [1] K. J. Thomas, *et al.*, Phys. Rev. Lett. **77**, (1996) 135.
- [2] A. Kristensen, *et al.*, Phys. Rev. **B 62**, (2000) 10950. D. J. Reilly, *et al.*, Phys. Rev. **B 63**, (2001) 121311. S. Nuttinck, *et al.*, Jpn. J. Appl. Phys. **39** (2000) L655.
- [3] S. M. Cronenwett, *et al.*, Phys. Rev. Lett. **88** (2002) 226805.
- [4] Y. Meir, *et al.*, Phys. Rev. Lett. **89**, (2002) 196802.
- [5] C.-K. Wang and K.-F. Berggren, Phys. Rev. **B 54** (1996) 14257. H. Bruus, *et al.*, Physica E **10** (2001) 97. K. Hirose, *et al.*, Phys. Rev. **B 63** (2001) 33315.
- [6] Y. Tokura and A. Khaetskii, Physica E, **12** (2002) 711.
- [7] A. Oguri, J. Phys. Soc. Jpn, **66** (1997) 1427.
- [8] Y. Levinson and P. Wölfle, Phys. Rev. Lett. **83** (1999) 1399.
- [9] Y. Tokura, unpublished.

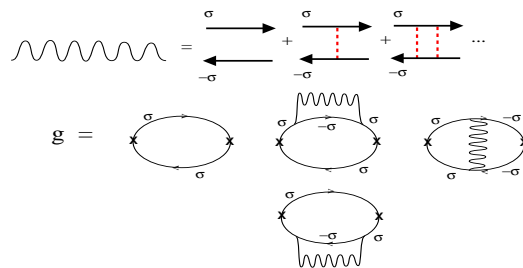


Fig.1 Feynman diagrams in the calculation of the paramagnon propagator (above) and the conductance up to lowest order in the paramagnon propagator. The solid line represents the electron propagator and σ is the electron spin.