

OBSERVATION OF SPIN POLARIZATION IN TRANSITION METAL DICHALCOGENIDE

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Monolayer transition-metal dichalcogenides (TMDC) have been attracting great attention as the candidate material for next-generation electronic devices utilizing spin and/or valley degrees of freedom [1-5]. As shown in Fig. 1(a), for example, the monolayer-MoS₂ is formed by the 2-dimensional network of Mo sitting in the center of S triangular prisms, thus possessing the 3-fold symmetry without an inversion center. Because of this peculiar trigonal structure, in contrast to the simple honeycomb graphene, the spontaneous spin and valley coupled band structure (*i.e.* spin/valley polarization) is realized at the Brillouin zone corners. Indeed, the control and detection of the valley polarization are demonstrated by using the photoluminescence circular dichroism in monolayer-MoS₂ [1-4]. Monolayer-MoS₂ hitherto has been mostly obtained by mechanically exfoliating the 2H-polytype bulk crystal, which itself lacks the spin/valley polarization due to its centrosymmetric structure recovered by the hexagonal stacking. The exfoliated monolayer-MoS₂ thus obtained has the limited size of area ($< 1 \mu\text{m}^2$), thus preventing the direct investigations on its spin and electronic structure. Here we focus on another stable phase of polytype, 3R-MoS₂ (space group $R\bar{3}m$) [6-8], which is composed of the so-called *abc* stacking where the 3-fold symmetry of monolayer-MoS₂ remains in bulk, as shown in Fig. 1(b). In 3R-MoS₂, the spin/valley coupling similar to monolayer-MoS₂ should reside in bulk, appearing as the opposite spin-polarizations at \bar{K} - and \bar{K}' -points [Fig. 1(c)].

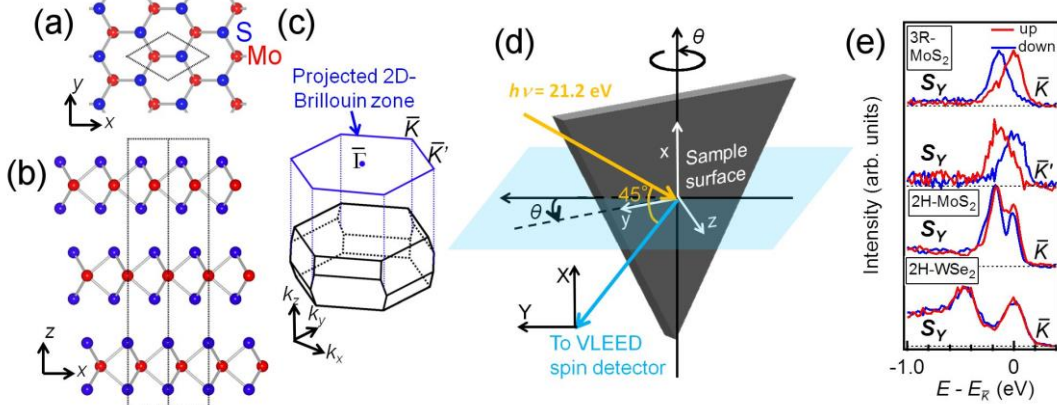


FIGURE 1. (a) Top view of MoS₂ monolayer. (b) Side view of 3R-MoS₂. (c) The first Brillouin zone of 3R-MoS₂. The blue hexagonal plane with $\bar{\Gamma}$, \bar{K} , and \bar{K}' represents the corresponding two-dimensional projected Brillouin zone. (d) Schematic of the SARPES experimental geometry at BL19A in photon factory, KEK[9]. (e) SARPES energy distribution curves at the \bar{K} (\bar{K}') point, with detection of the spin component along Y-direction [as shown in (d)], for 3R-MoS₂, 2H-MoS₂, and 2H-WSe₂.

Spin-resolved angle-resolved photoemission spectroscopy (SARPES) with HeI α light source ($h\nu = 21.2 \text{ eV}$) was performed by using the VLEED (Very Low Energy Electron Diffraction) spin detector at BL19A in photon factory, KEK [9]. The angular resolution was set to $\pm 1^\circ$ and the total energy resolution was set to 100 meV. Samples ($\sim 1 \text{ mm} \times 1 \text{ mm}$) were cleaved in situ near room-temperature and measured at 110 K. Measurement geometry is shown in Fig. 1(d). The VLEED detector can measure the spin polarization along X and Y directions. Note that the sample orientation (x, y, z) is rotated around X-axis to measure the \bar{K} -point.

Figure 1(e) shows the SARPES spectra detecting Y -oriented spin ($I_Y^{\uparrow,\downarrow}$) of 3R-MoS₂, 2H-MoS₂ and 2H-WSe₂, measured at \bar{K} - and \bar{K}' -points. $I_Y^{\uparrow,\downarrow}$ includes the information on spin polarization along z and y , due to the sample orientation. $I_Y^{\uparrow,\downarrow}$ data for 3R-MoS₂ at \bar{K} -point clearly indicate the huge spin polarization, *i.e.* spin up (down) dominant for the upper (lower) band, which gets inverted at \bar{K}' -point. By comparing with the *ab-initio* calculation, the opposite spin-polarizations at \bar{K} - and \bar{K}' - points observed in 3R-MoS₂ offer the evidence of the spontaneous out-of-plane spin polarization which consequently accompanies the \bar{K} - and \bar{K}' -valley polarization, as similarly expected in monolayer-MoS₂ [1-5]. 2H-MoS₂ and 2H-WSe₂, on the other hand, show the nearly equivalent intensities of I_Y^{\uparrow} and I_Y^{\downarrow} , indicating that the spin polarization is absent or very weak. It should be reflecting the centrosymmetric crystal structure of 2H-polytype, which at the same time suggests the spin-splitting observed in 3R-MoS₂ is indeed from bulk, not arising from the inversion breaking at the surface.

To summarize, we investigated the spin- and electronic structure at the Brillouin zone corner for noncentrosymmetric 3R-MoS₂ and centrosymmetric 2H-MoS₂ and 2H-WSe₂ by using spin-resolved ARPES. The observed spin polarization in 3R-MoS₂ is identical to that expected in monolayer-MoS₂ and other TMDC families. The noncentrosymmetric 3R phase should be useful for future device fabrications, that will drive the valley/spin-tronics based on 2 dimensional crystals

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