

June 3rd 2015 Workshop on

“New Perspectives in Spintronic and Mesoscopic Physics”

Quantum transport phenomena in 3D topological insulator thin films

Ryutaro Yoshimi^A,

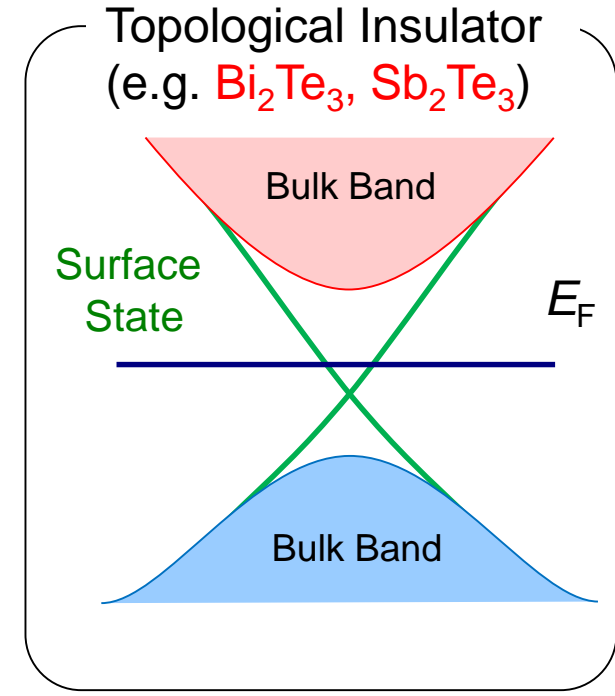
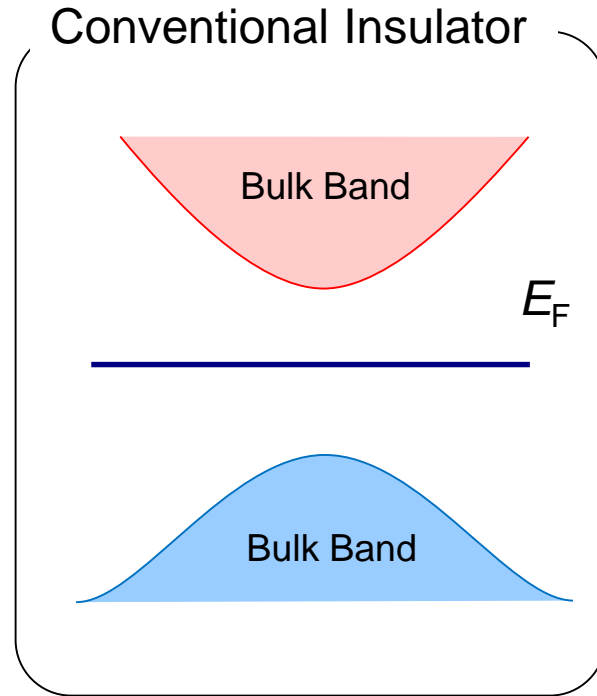
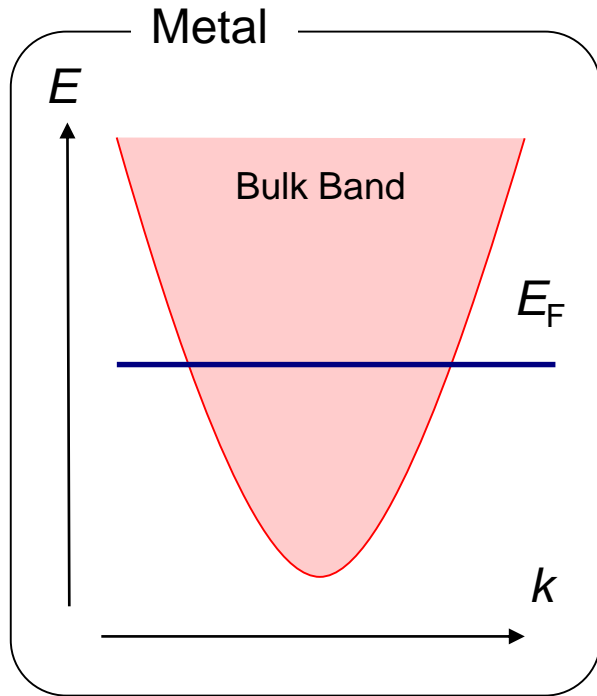
A. Tsukazaki^B, M. Mogi^A, K. Yasuda^A,

Y. Kozuka^A, J. Falson^A, J. G. Checkelsky^D, K. S. Takahashi^C,

N. Nagaosa^{A,C}, M. Kawasaki^{A,C}, and Y. Tokura^{A,C}

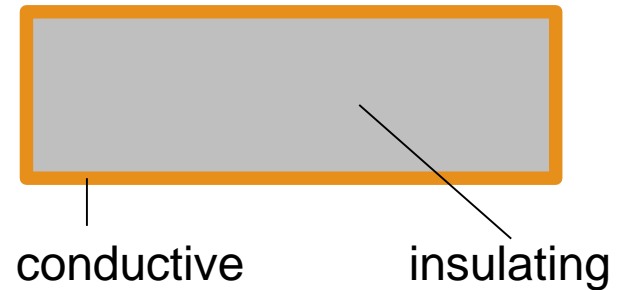
Univ. of Tokyo^A, Tohoku Univ.^B, RIKEN-CEMS^C, MIT^D

Introduction to Topological Insulator

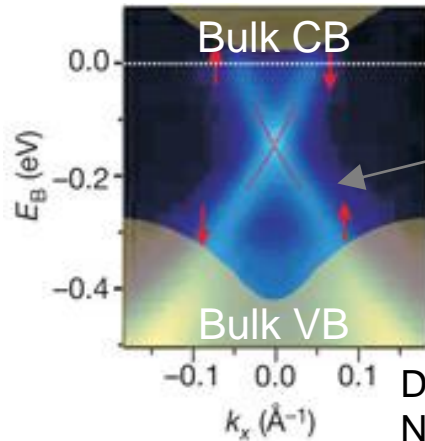


Topological Insulator

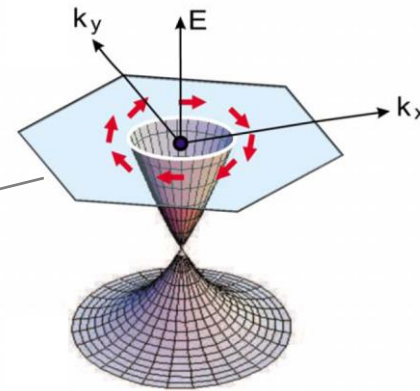
- Bulk (inside the sample) is insulating
- Surface is **conductive**



Momentum space



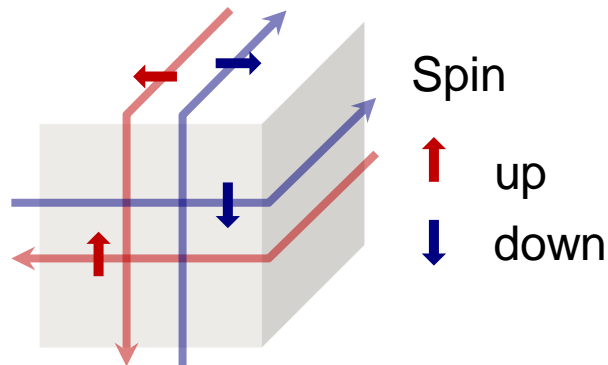
D. Hsieh *et al.*,
Nature (2009)



Two-dimensional surface gapless band

- massless **Dirac** dispersion
- **spin** momentum locking
- **gapless** under time reversal symmetry

Real space



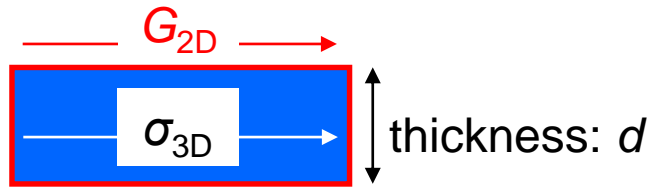
Dissipationless **spin-current**
on the surface



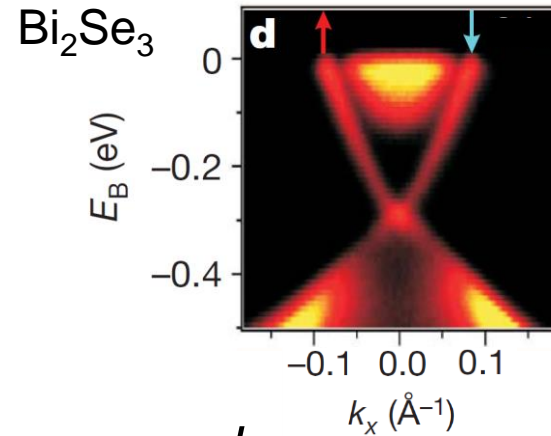
Future application to spintronics device or quantum computation

Bi₂Se₃, Bi₂Te₃, Sb₂Te₃,

- Small band gap (~ 250 meV)
- **Bulk residual carrier** by crystal defects



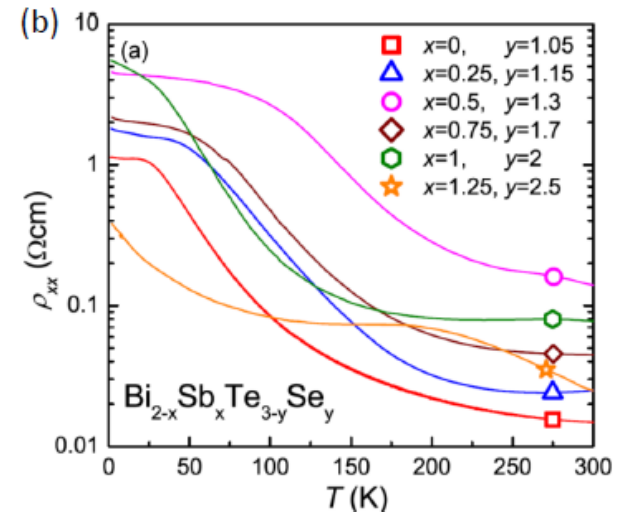
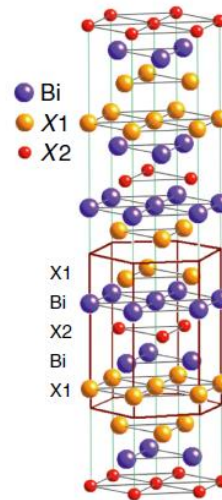
$$G_{\text{total}} = \underbrace{G_{2D}}_{\text{surface}} + \underbrace{\sigma_{3D} \times d}_{\text{bulk conduction}}$$



D. Hsieh *et al.*, Nature (2009)

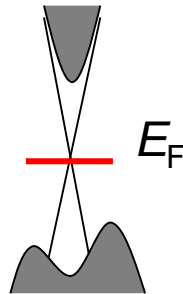
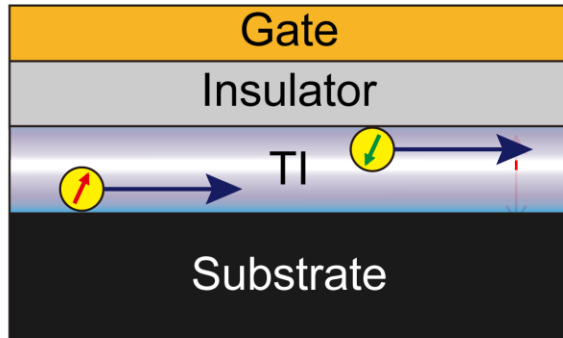
Reduction of bulk conduction

- Exfoliation (reduction in d)
- Find insulating compound (reduction in σ_{3D})



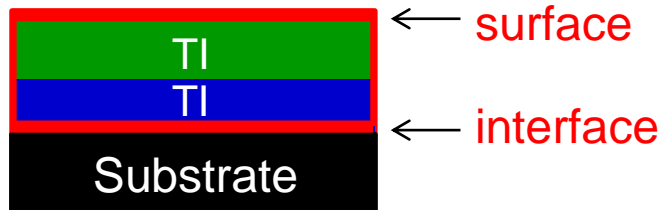
Y. Ando, Review in JPSJ (2013)

Thin film + FET device



- **Surface dominant** transport
- **E_F tuning** in the Dirac state

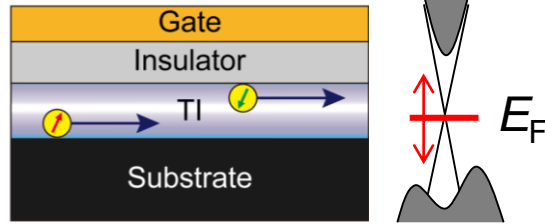
Surface/interface degrees of freedom



- Controlling surface/interface (lift the degeneracy) by **layered structure**

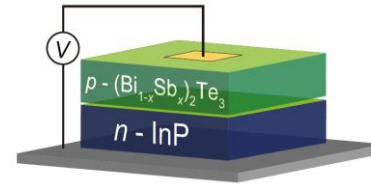
- Thin film device enable us to access the ideal Dirac states of TI without bulk state
- Surface/interface degrees of freedom shows up

Thin film device fabrication



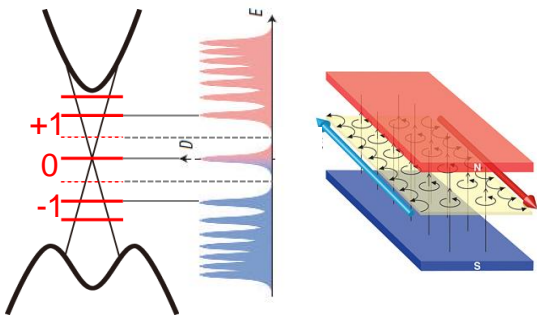
Bulk insulating thin film
FET device for E_F tuning

Tunneling Spectroscopy



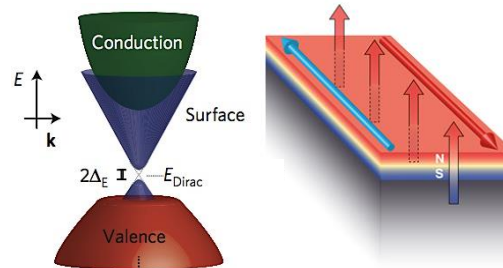
R. Yoshimi *et al.*,
Nature Materials **13** 253 (2014)
Detection of
interface Dirac state

Quantum Hall Effect



R. Yoshimi, *et. al*
Nature Commun. **6**,6627 (2015)

Quantum Anomalous Hall Effect



J. G. Checkelsky, R. Yoshimi *et.al.*
Nature Physics, **10**, 731 (2014)

Quantization **at zero field**

QHE in semi-magnetic TI bilayer

mag



non-mag



InP substrate

R. Yoshimi *et al.*, submitted.

QHE **at higher temperature**

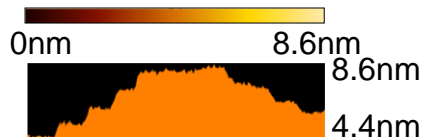
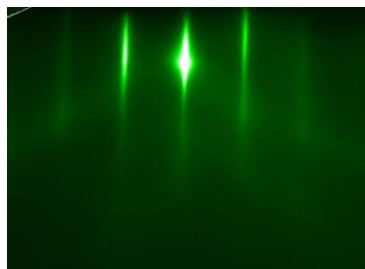
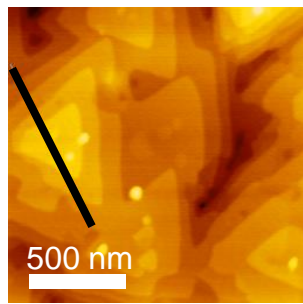
$(\text{Bi}_{1-x}\text{Sb}_x)_2\text{Te}_3$ (BST) thin film

Molecular Beam Epitaxy (MBE)

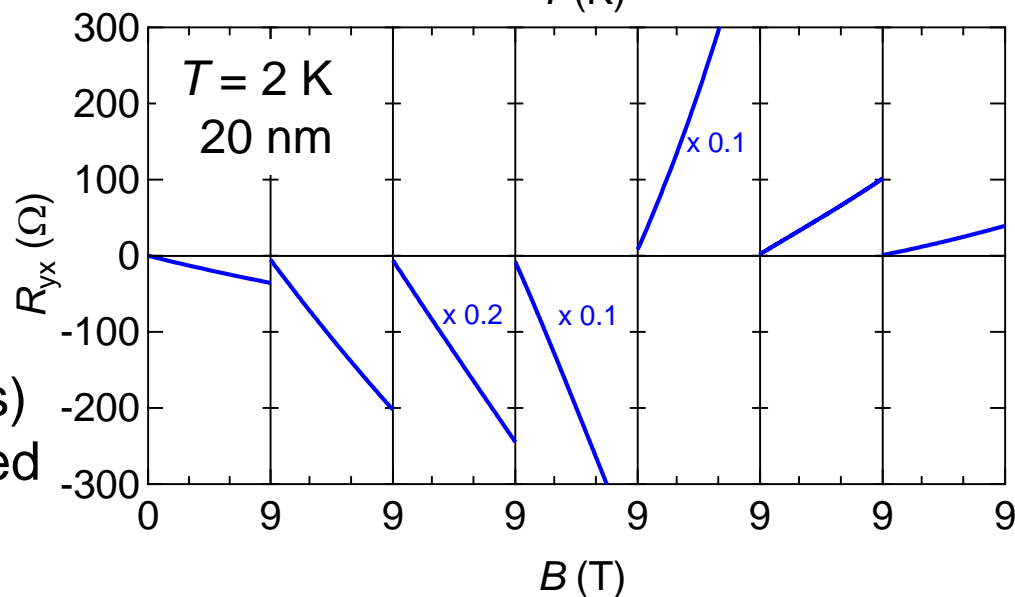
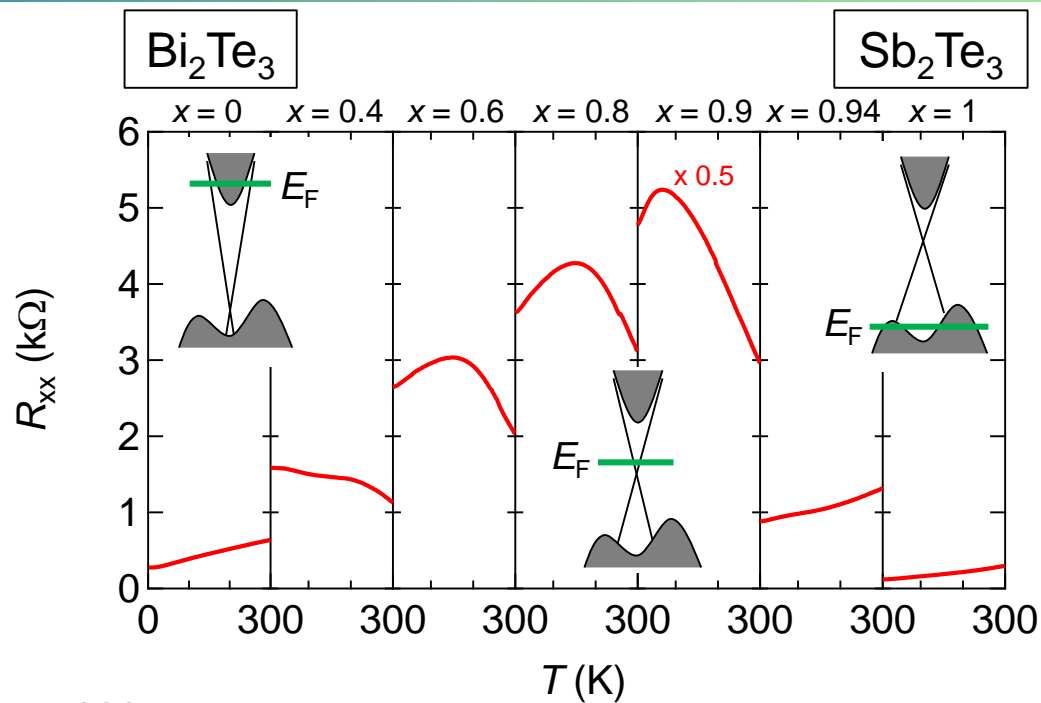


AFM

RHEED

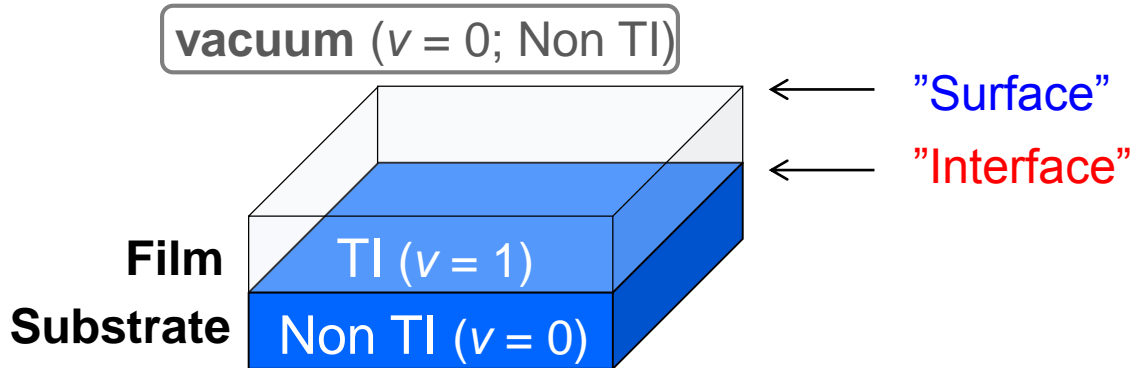


- Substrate: InP (lattice matches)
- Carrier density widely controlled with Sb content

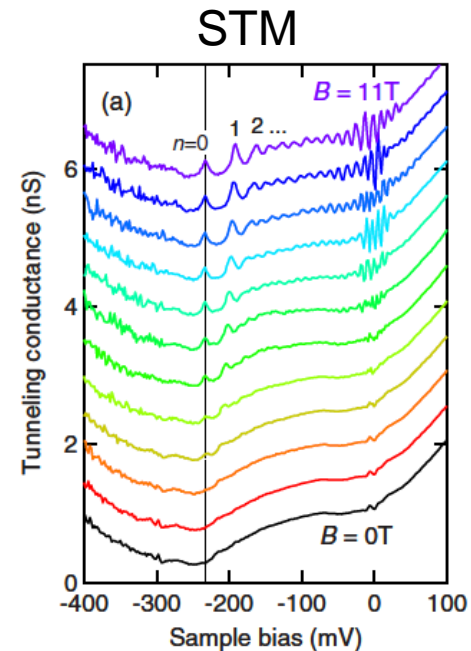
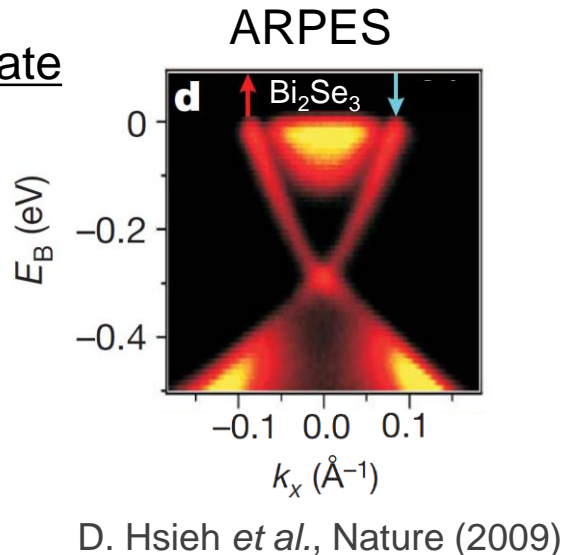


Surface/Interface Dirac states in topological insulator 8

Edge states at the boundary of TI



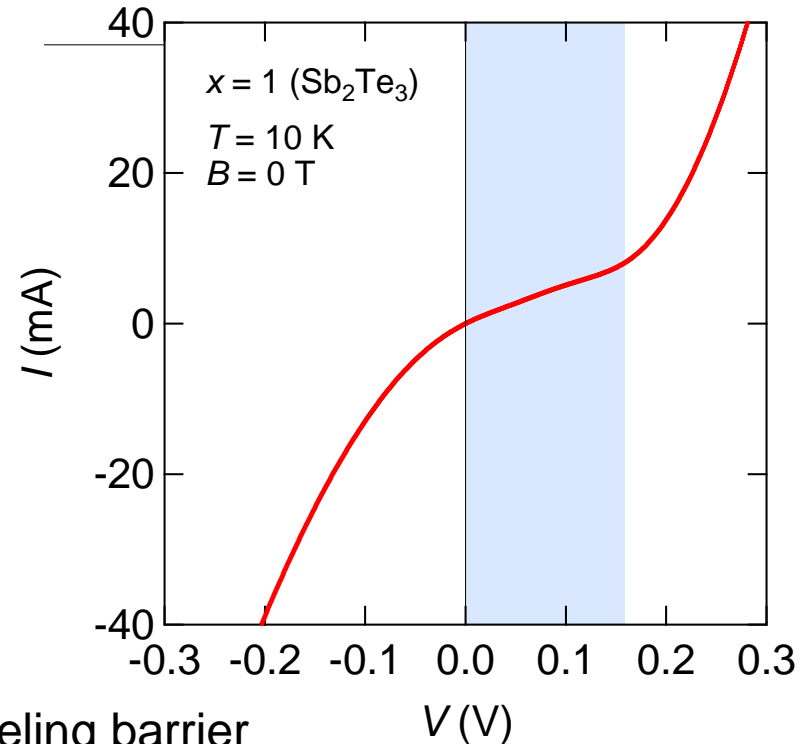
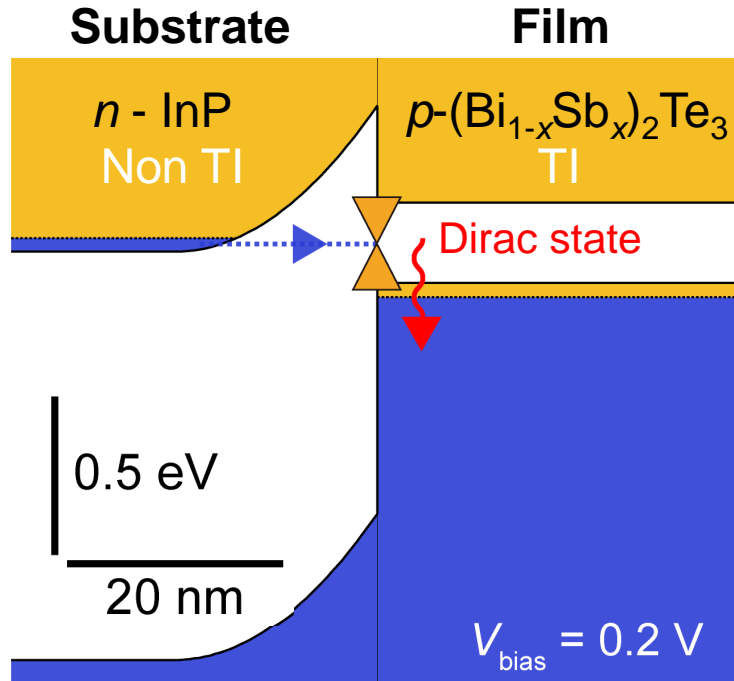
Surface Dirac state



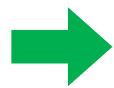
T. Hanaguri *et al.*,
PRB (2010)

- Surface Dirac states are verified by ARPES and STM
- Dirac state at interface is not yet clearly observed

Tunneling spectroscopy on TI/non-TI p - n heterojunction

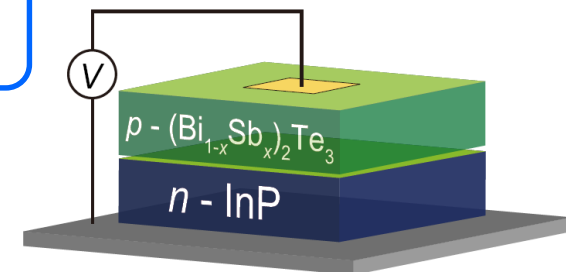


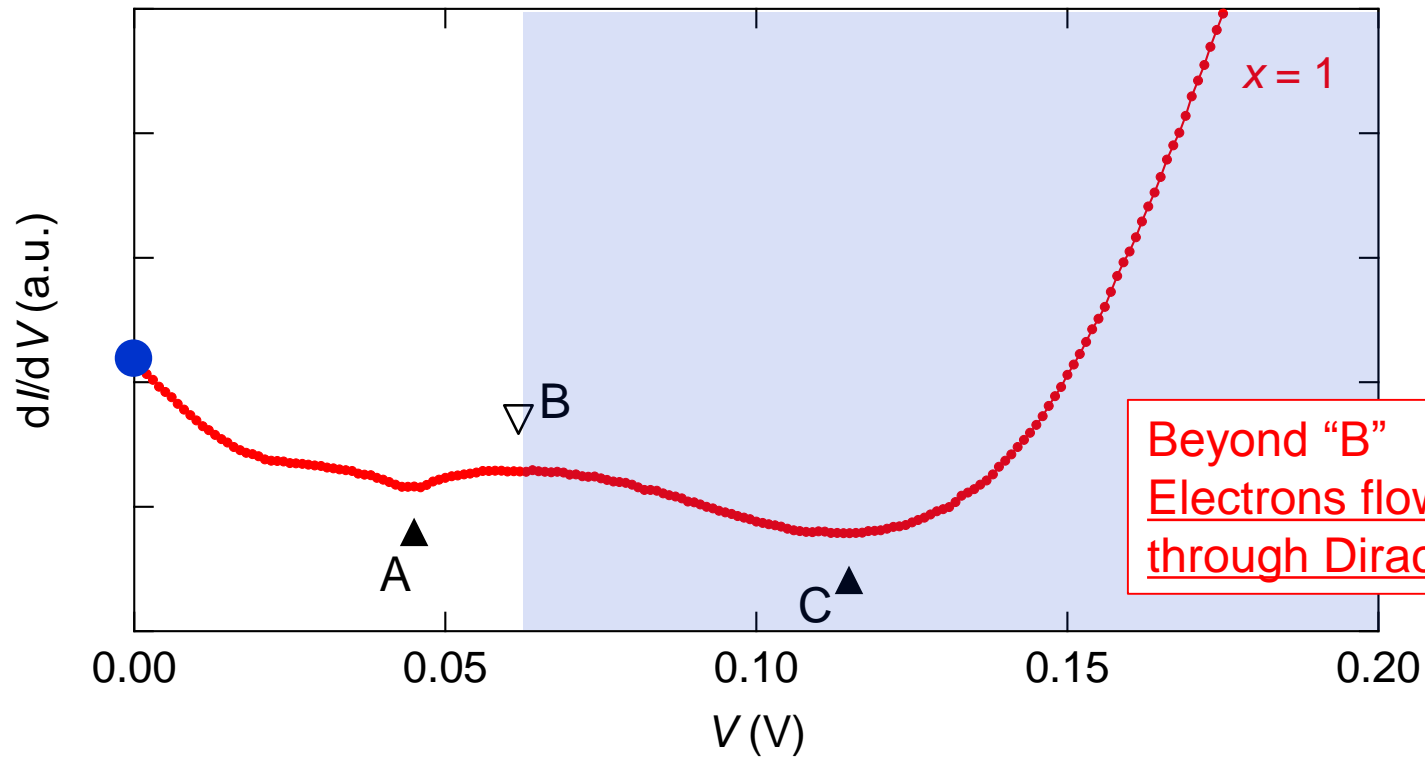
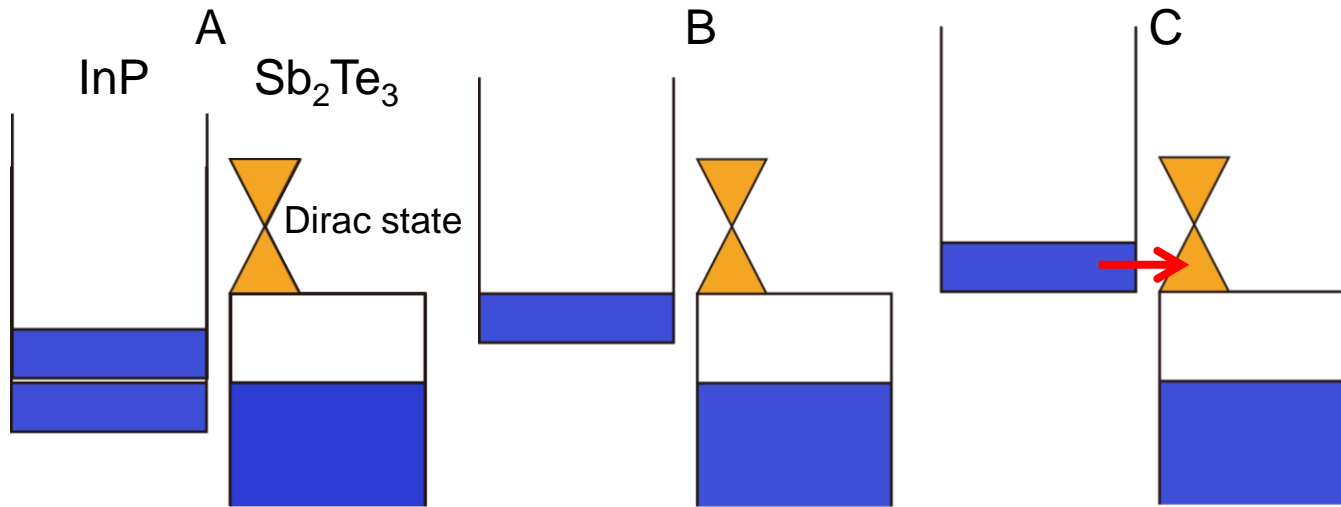
- Thin depletion layer acts as the tunneling barrier
- Tunneling conductance is proportional to the interface density of states



Interface Dirac states and its Landau level formation can be spectroscopically detected

- Typical Esaki-diode character in I - V curve

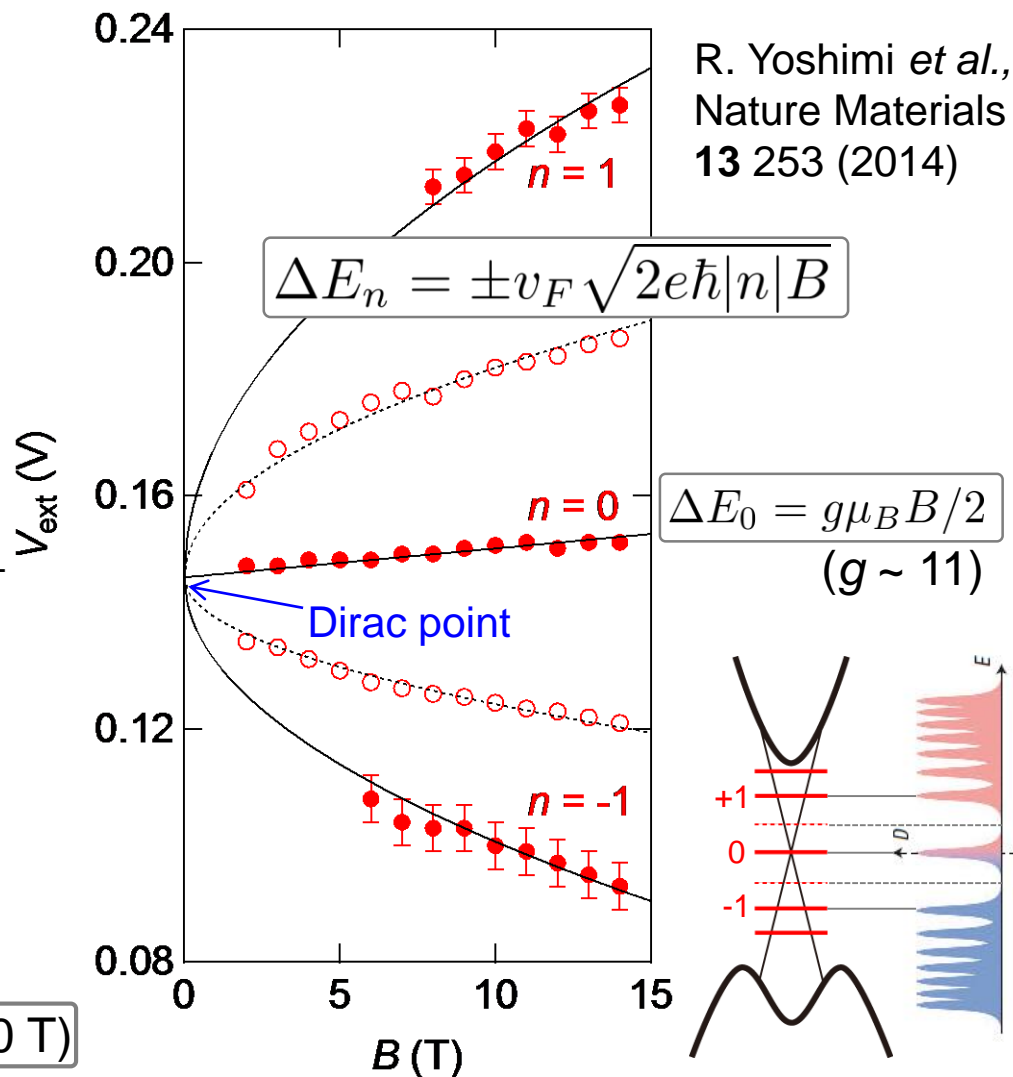
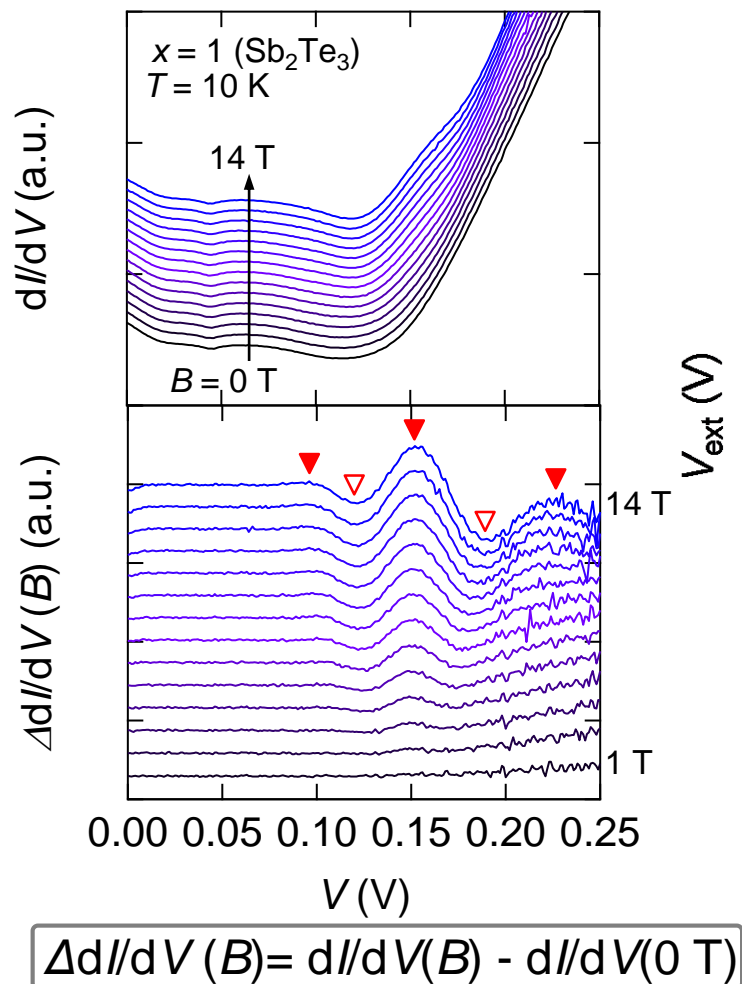




Beyond "B"
Electrons flow
through Dirac state

Landau level formation observed in tunneling spectra¹¹

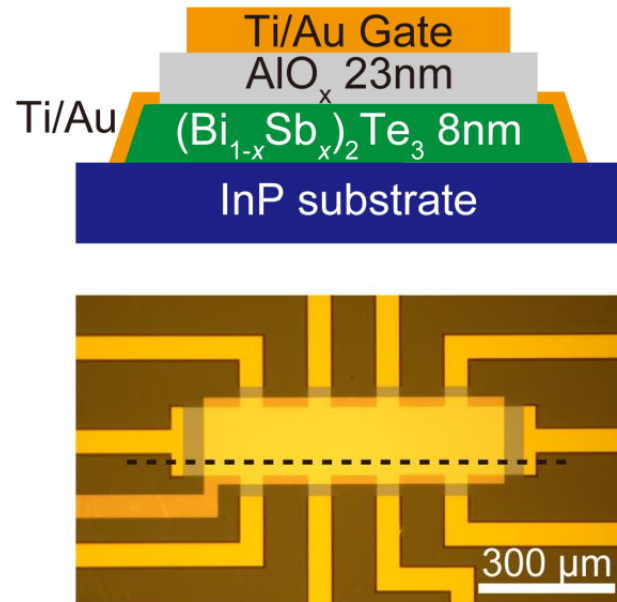
Tunneling conductance dI/dV



R. Yoshimi *et al.*,
Nature Materials
13 253 (2014)

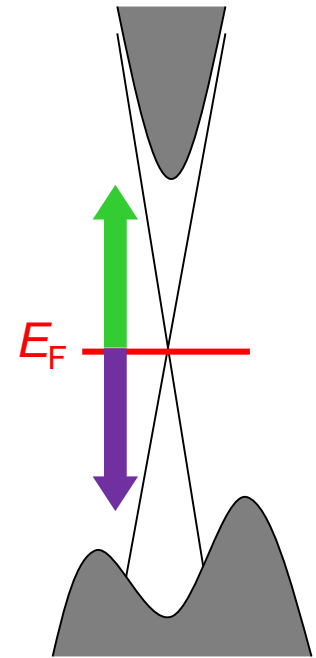
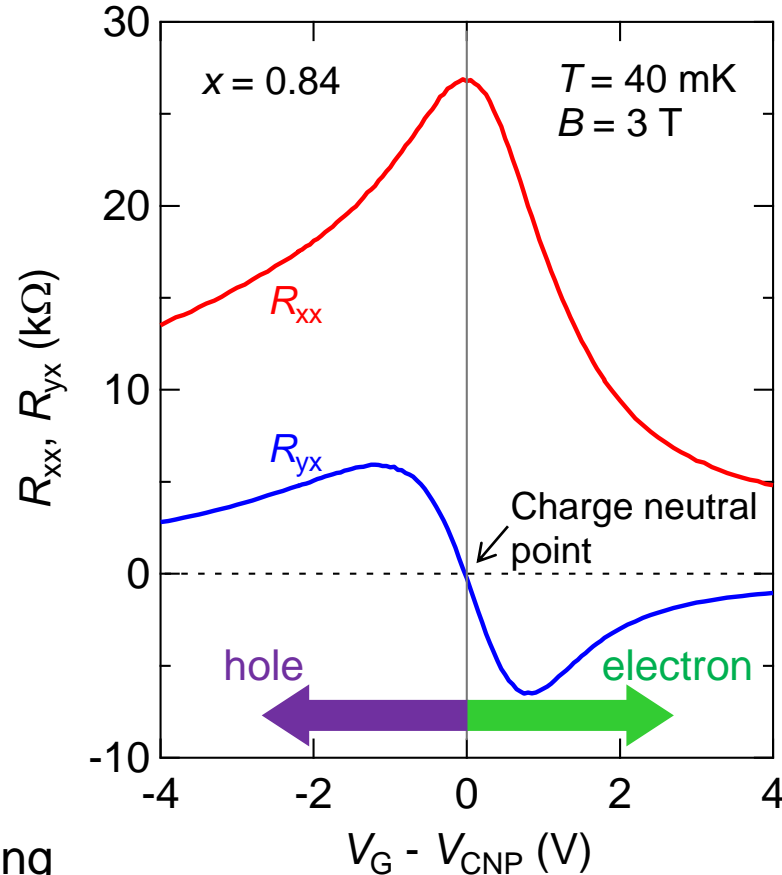
- Oscillatory feature in $\Delta dI/dV$
- Landau levels ($n = 0, \pm 1$) due to **interface Dirac dispersion** ($\propto B^{1/2}$)

FET device



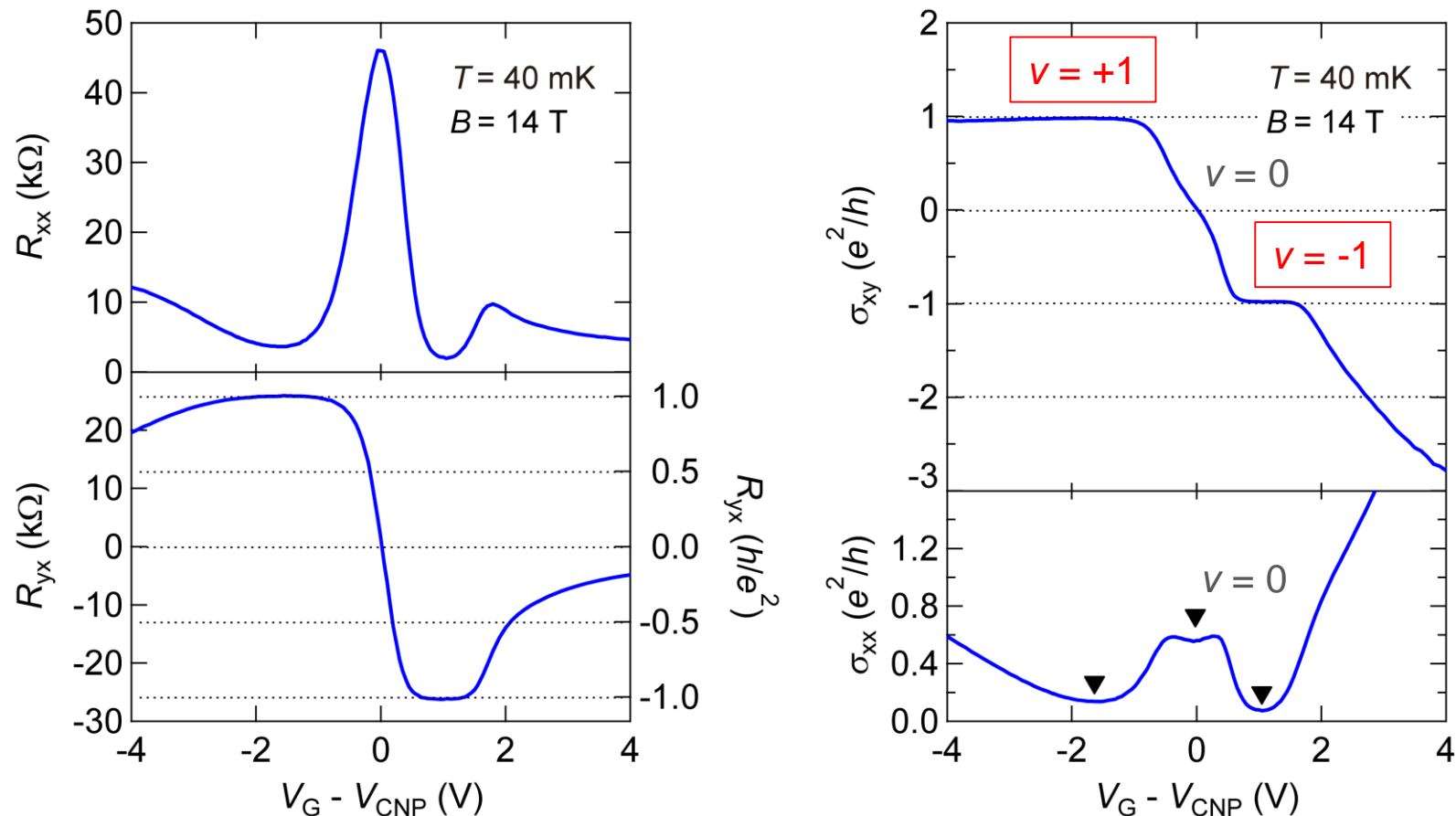
- AlO_x capping before patterning
- Low temperature process ($< 120 \text{ C}^\circ$)
- Gate dielectric by ALD at R.T.

Transport measurement in dilution fridge

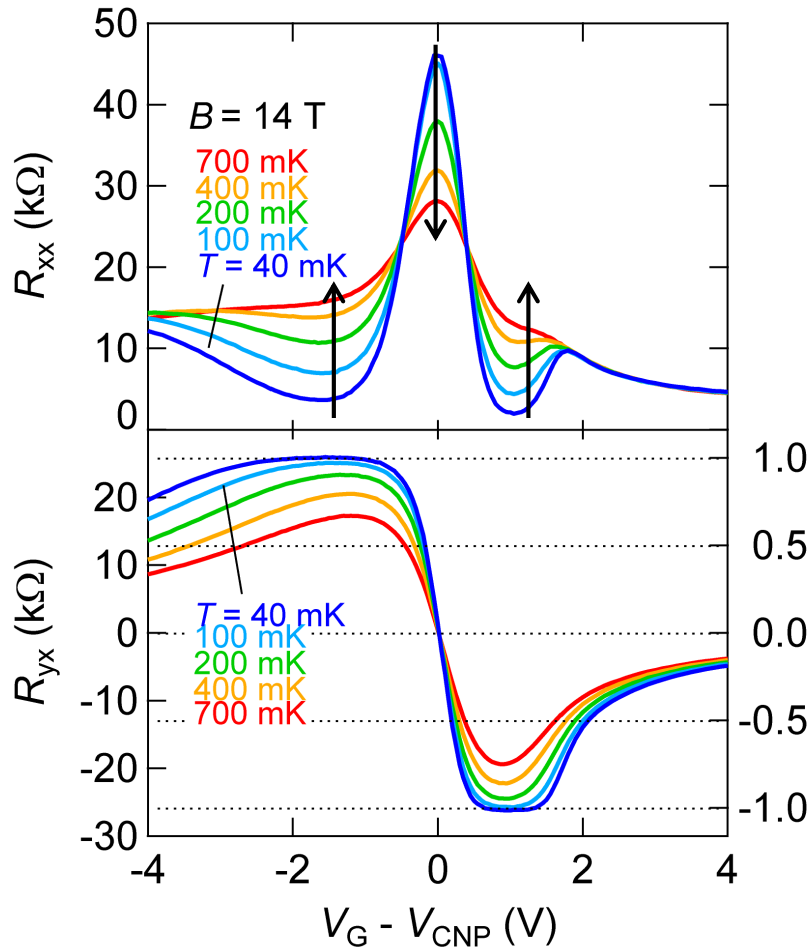


R. Yoshimi, et. al
Nature Commun.
6,6627 (2015)

- Ambipolar conduction by gate tuning
 E_F is modulated across Dirac point
- mobility $\mu \sim 1,500 \text{ cm}^2/\text{Vs}$, $n \sim 4 \times 10^{11} \text{ cm}^{-2}$

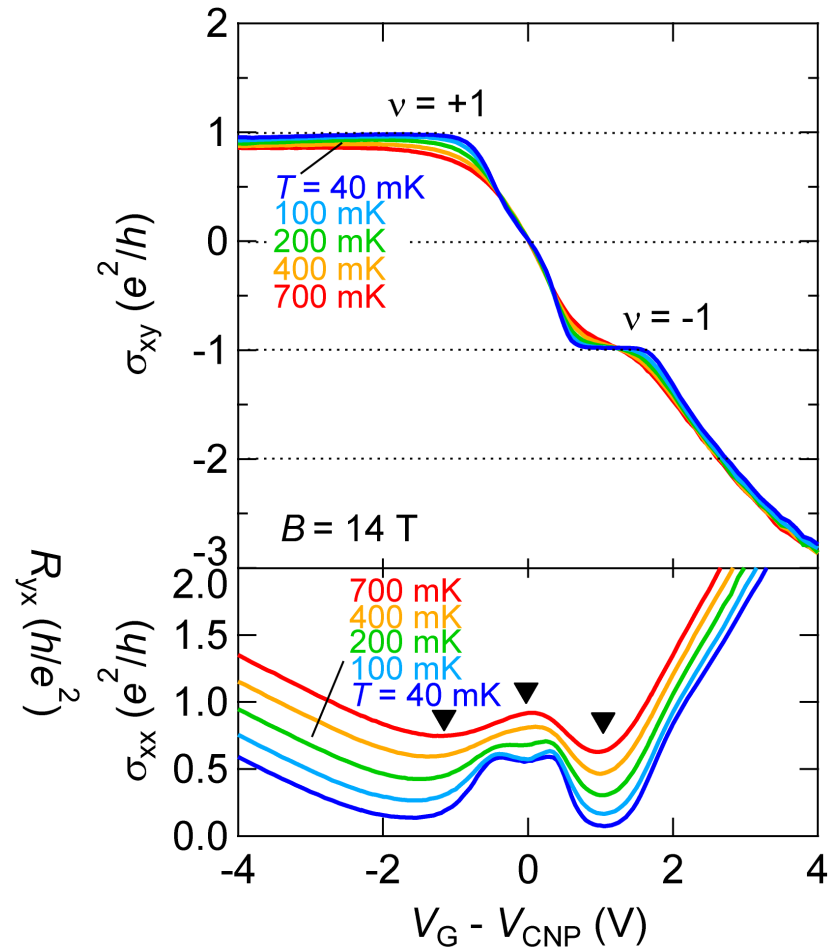


- Quantum Hall plateau $\sigma_{xy} = \pm e^2/h$ ($B = 14$ T)
- 2D nature of Dirac states
- At charge neutral point, $\sigma_{xy} = 0$ ($\nu = 0$) QHE



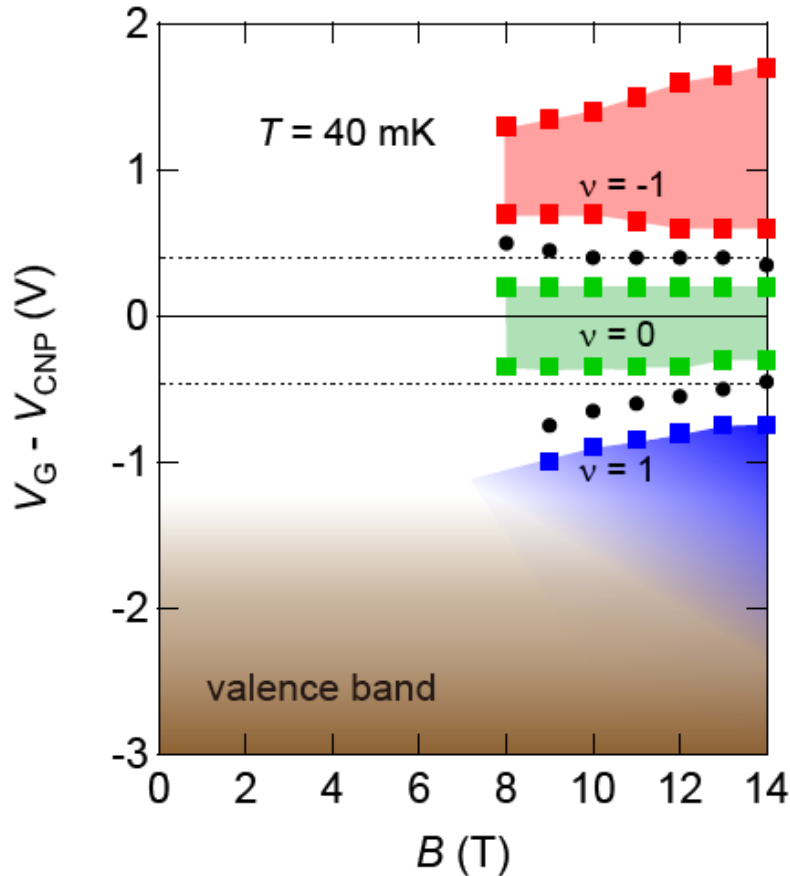
$\nu = \pm 1$: R_{xx} is metallic

$\nu = 0$: R_{xx} is insulating

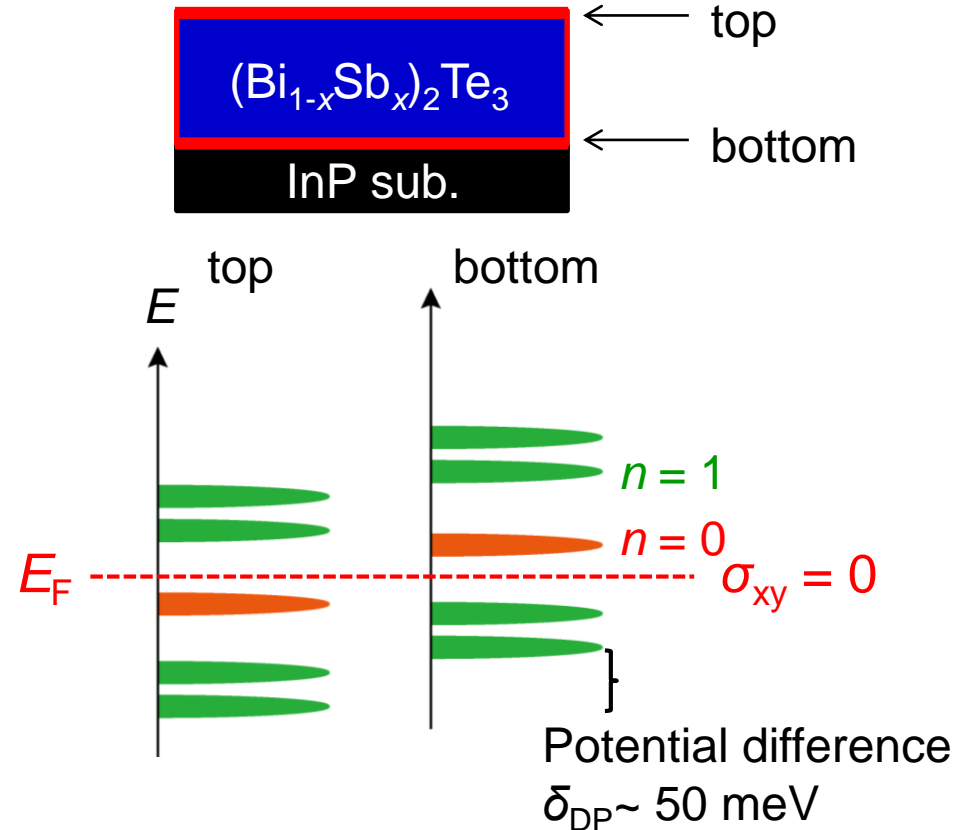


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 Nature Commun.
 6,6627 (2015)

Phase diagram of QH states

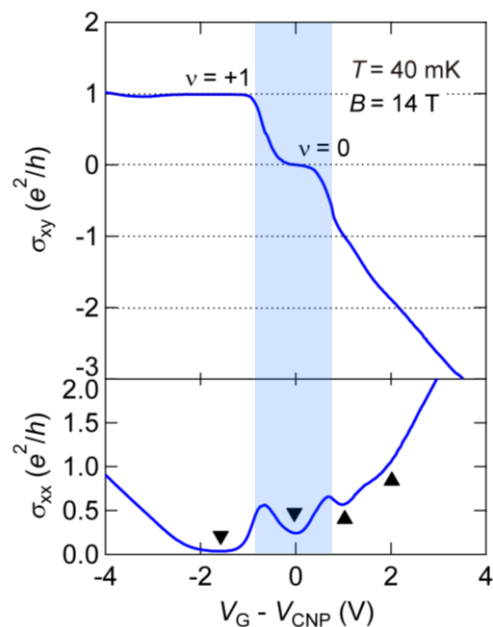
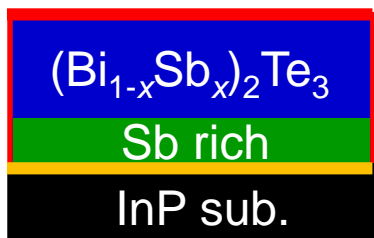
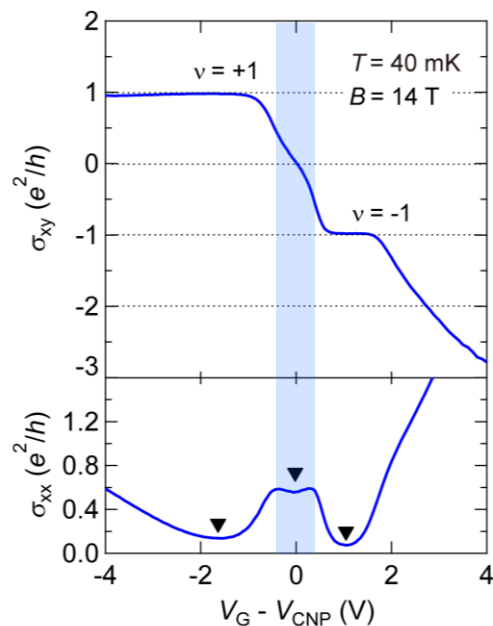
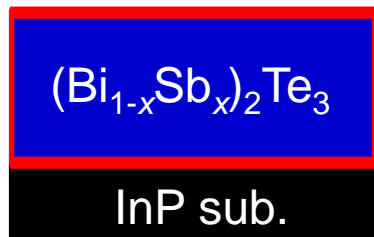


Landau levels at top and bottom surface

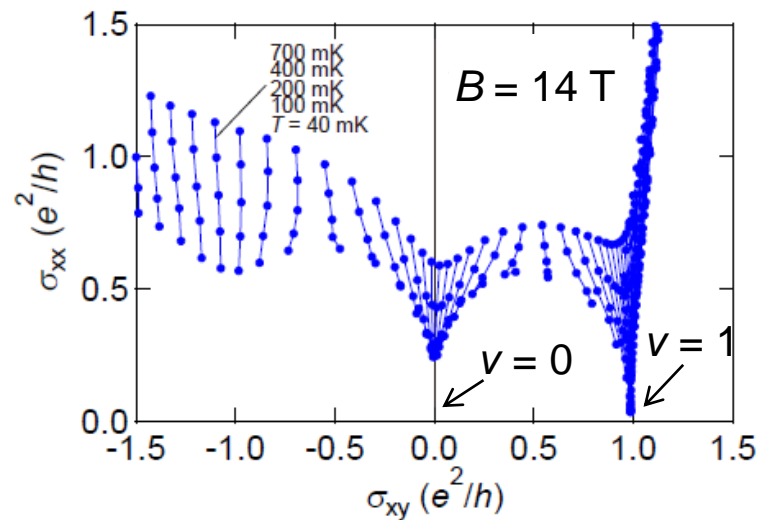
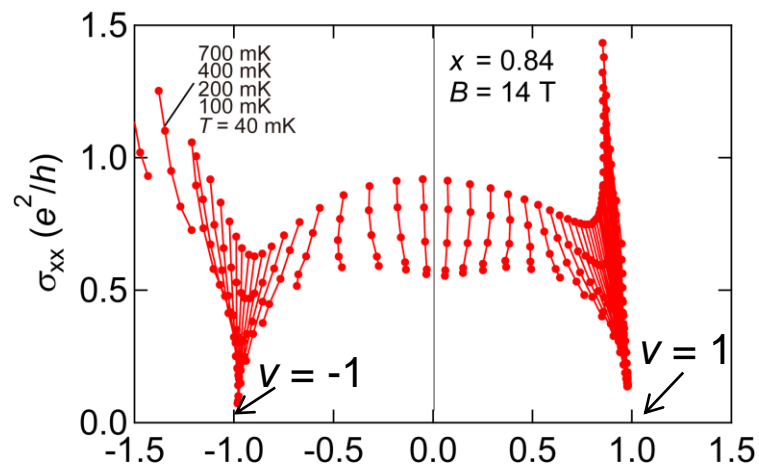


- weak B dependence in σ_{xy} plateau
- Different environment of top and bottom state $\rightarrow \delta_{\text{DP}} \sim 50$ meV

Structure

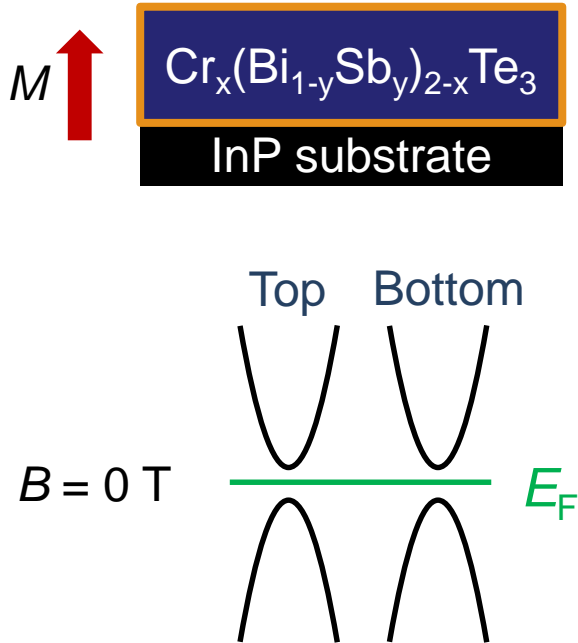


Conductivity mapping



Observation of Quantum Anomalous Hall Effect

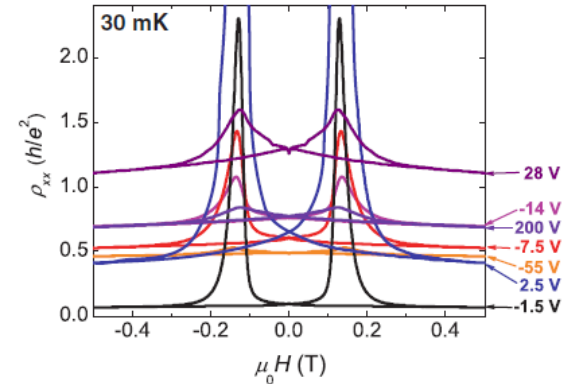
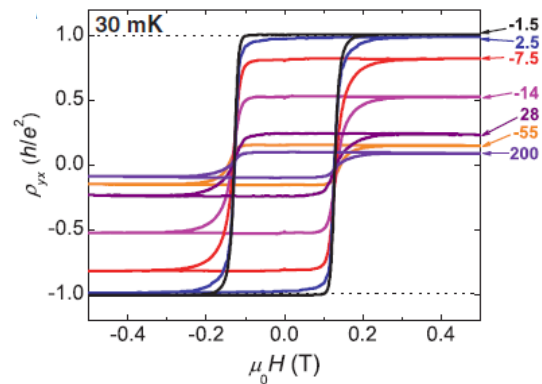
QAHE



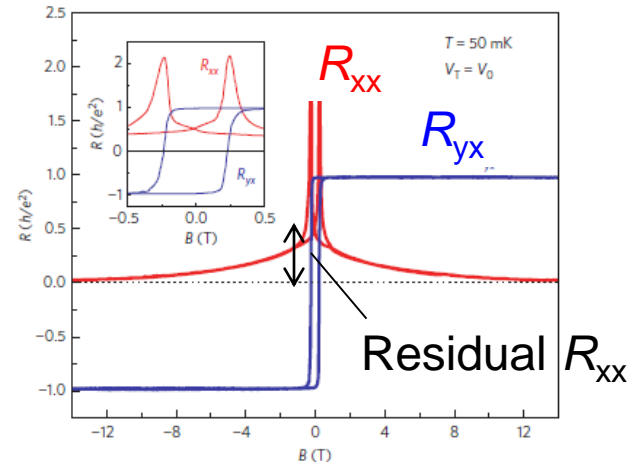
- Spontaneous gap opening by magnetism
- Quantization of **Anomalous Hall term**

$$\begin{cases} R_{yx} = R_{yx}^A = h/e^2 \\ R_{xx} = 0 \end{cases}$$

$\text{Cr}_x(\text{Bi}_{1-y}\text{Sb}_y)_{2-x}\text{Te}_3$



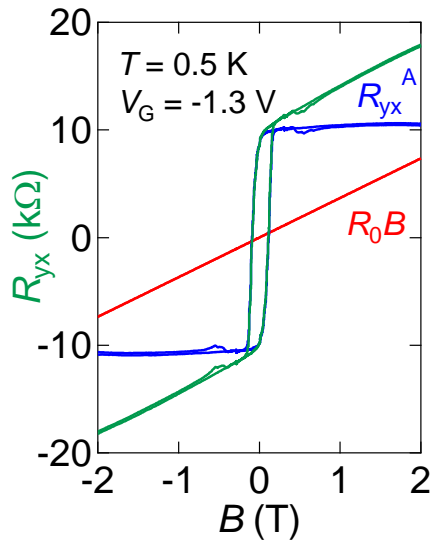
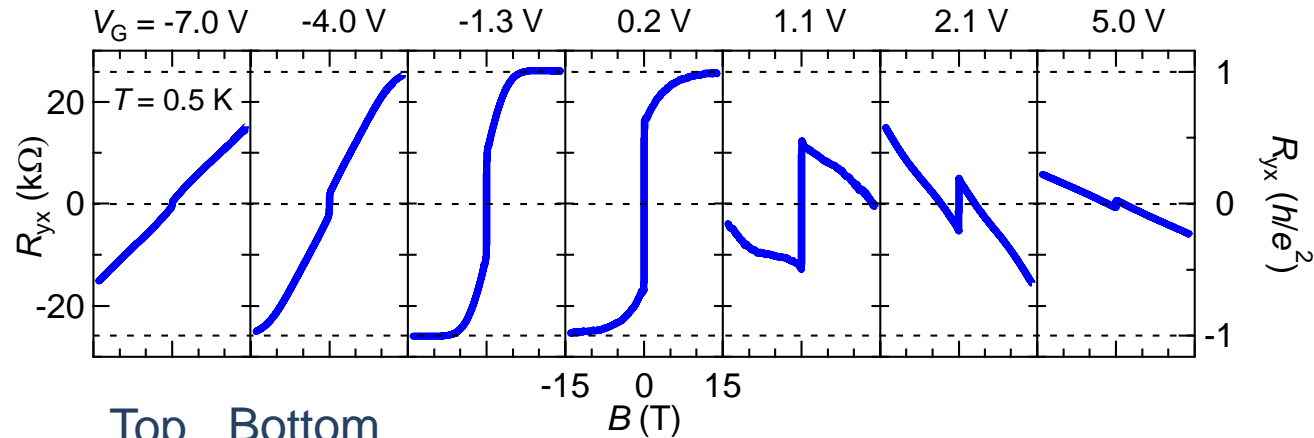
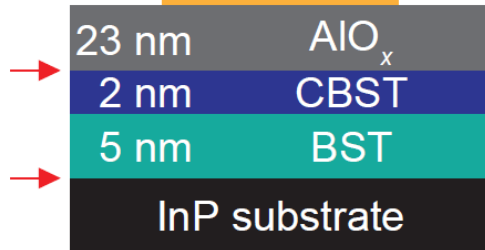
C-. Z. Chang *et al.*, Science **340** 167 (2013).



- QAHE with help of magnetic field (R_{xx} at 0 T $\sim 0.5 h/e^2$)

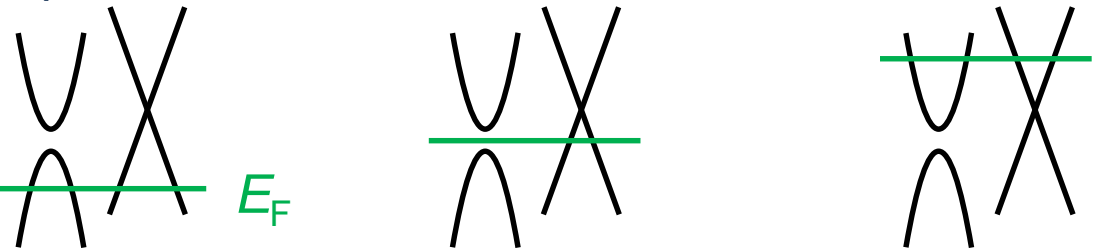
J. G. Checkelsky,
R. Yoshimi *et al.*,
Nature Phys. (2014)

Gate Ti/Au



$B = 0$ T

Top Bottom



- R_{yx} reaches h/e^2 by adding up **OHE** and **AHE**
- Band structure

$$R_{yx} = R_0 B + R_{yx}^A$$

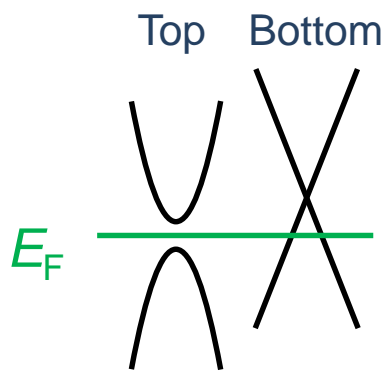
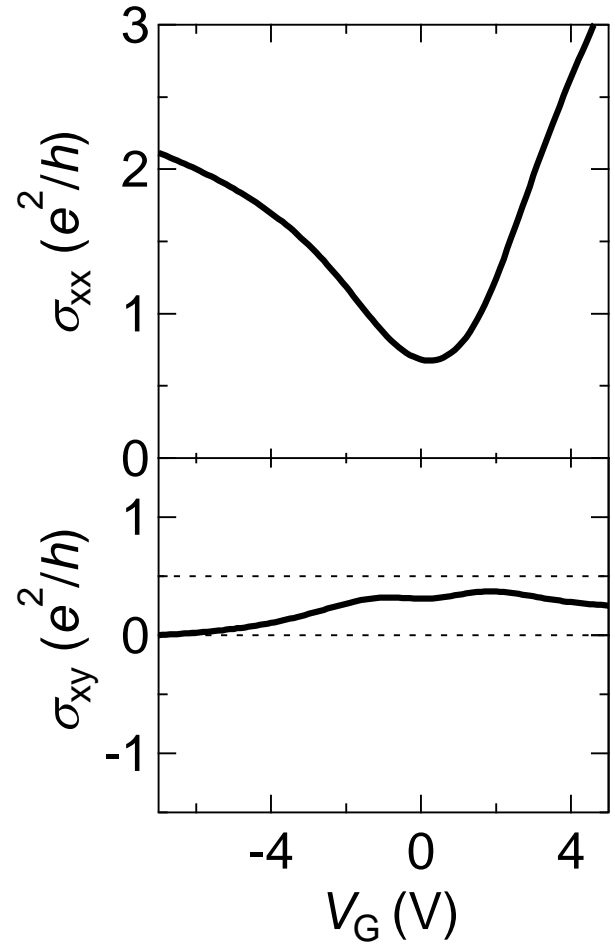
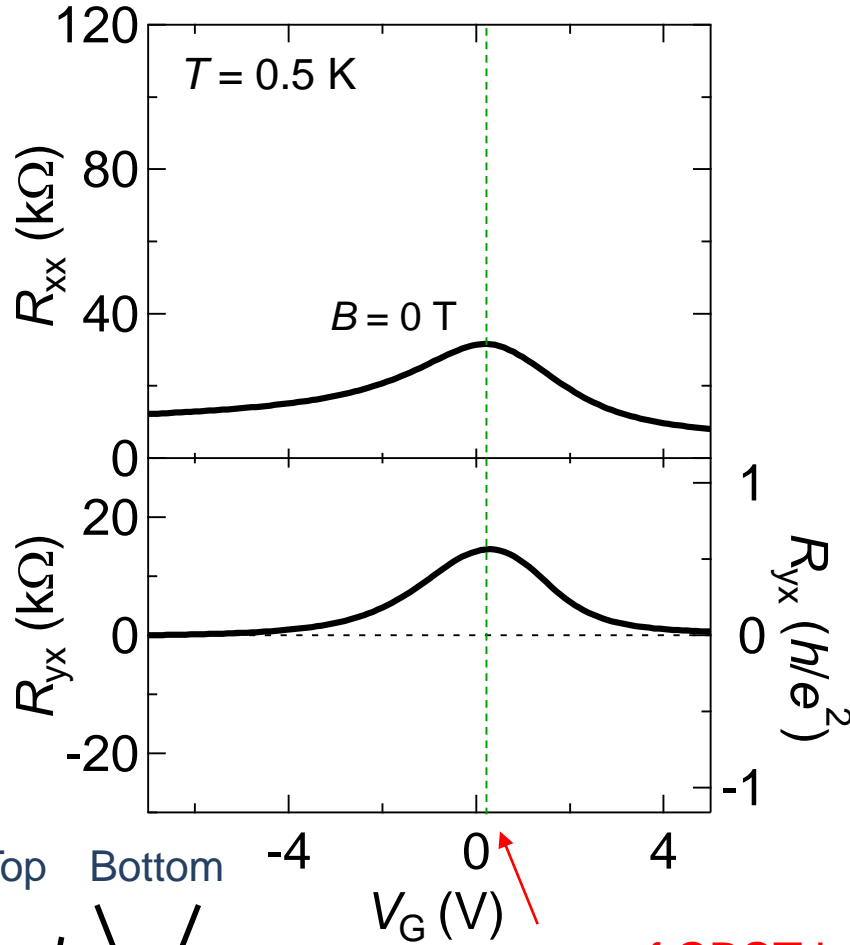
Ordinary Hall term Anomalous Hall term

BST CBST

gapped CBST surface
gapless BST surface

Transport Phenomena at 0T

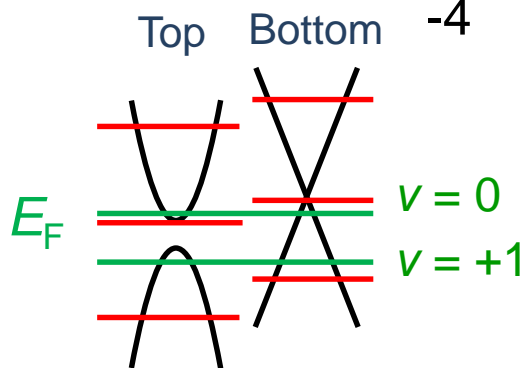
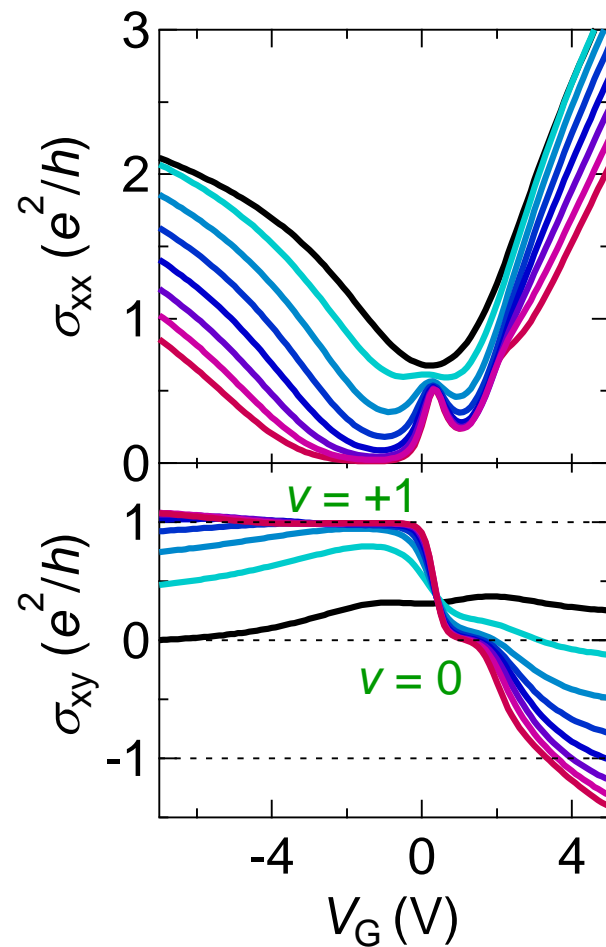
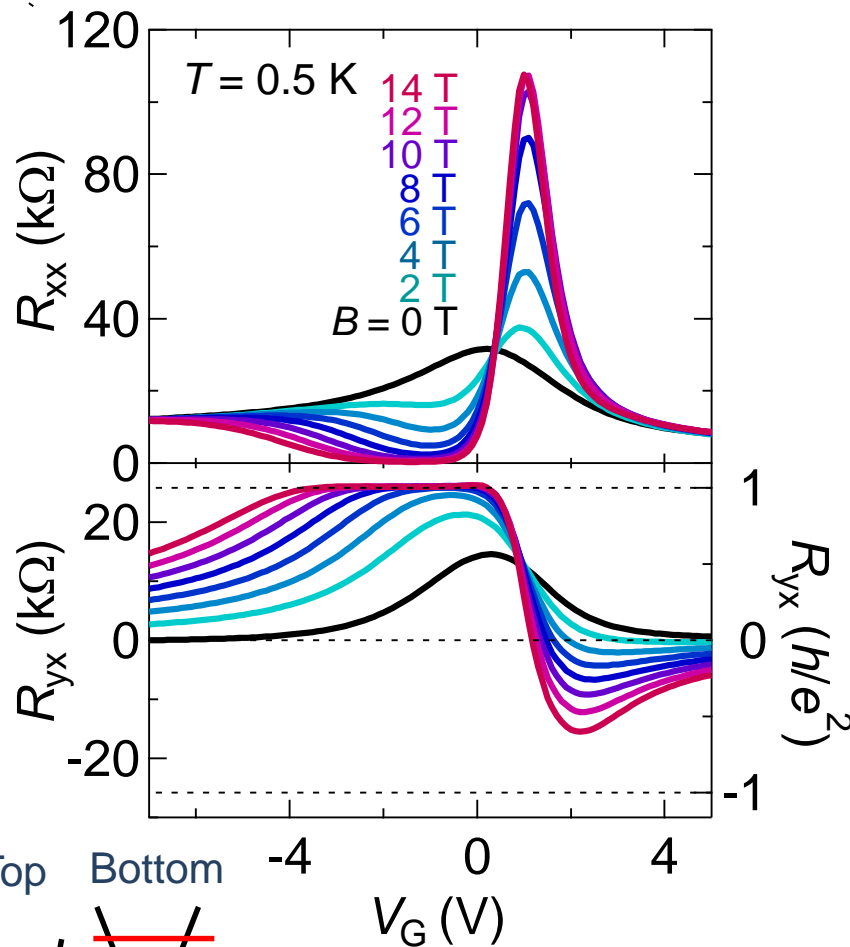
Gate Ti/Au	
23 nm	AlO_x
2 nm	CBST
5 nm	BST
InP substrate	



gap of CBST layer

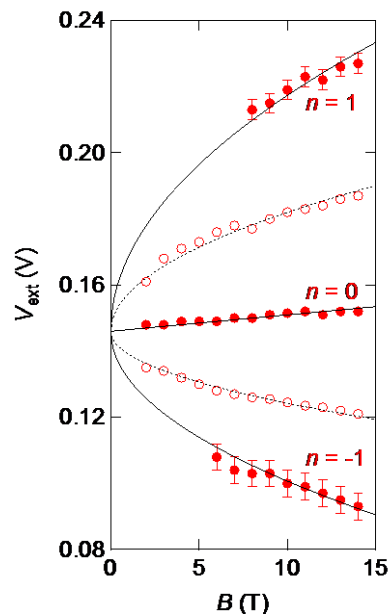
Magnetic field dependence

Gate Ti/Au	
23 nm	AlO _x
2 nm	CBST
5 nm	BST
InP substrate	



• $\nu = +1, 0$ states due to LL formation are observed

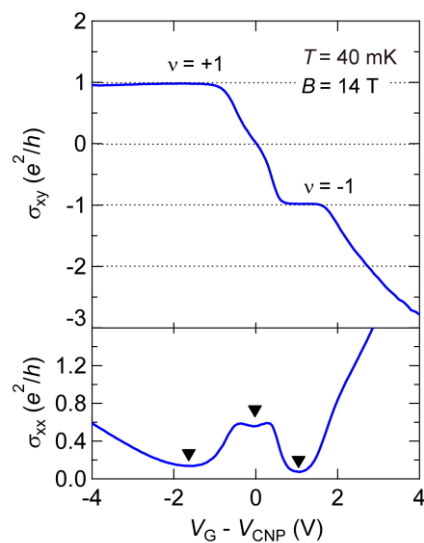
Tunneling Spectroscopy



R. Yoshimi *et al.*,
Nature Materials **13** 253
(2014)

Detection of
interface
Dirac state

QHE in BST

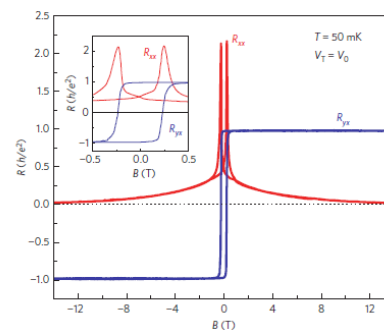


R. Yoshimi, *et al.*
Nature Commun.
6,6627 (2015)

FET device of
BST thin film

QHE with $\nu = \pm 1, 0$

QAHE in CBST

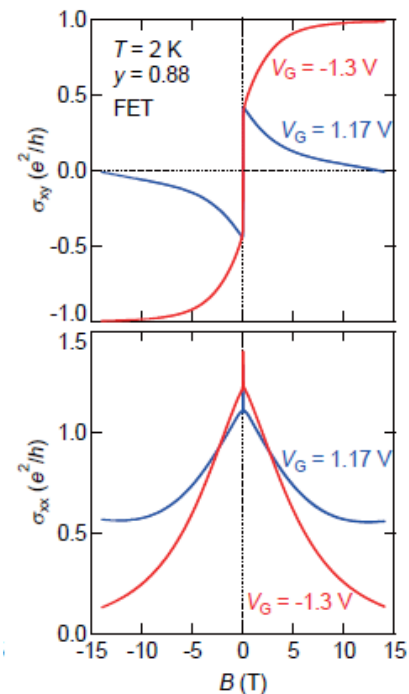


J. G. Checkelsky,
R. Yoshimi *et al.*,
Nature Phys. (2014)

Cr-doped BST

Quantization
at B = 0 T

QHE in bilayer



R. Yoshimi *et al.*,
submitted.

**QHE at higher
temperature**