

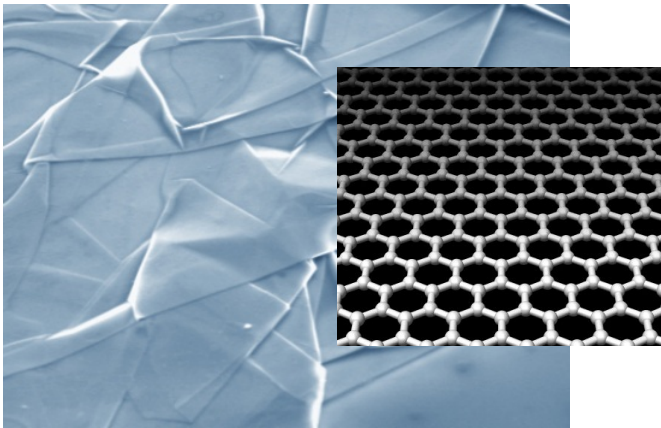
# Electronic transmission through the atomic domain boundary

--- from graphene to transition metal dichalcogenides

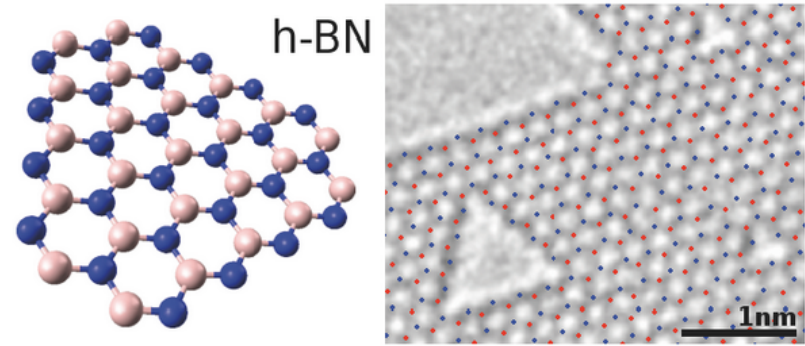
Mikito Koshino (Tohoku University)

# Family of 2D materials

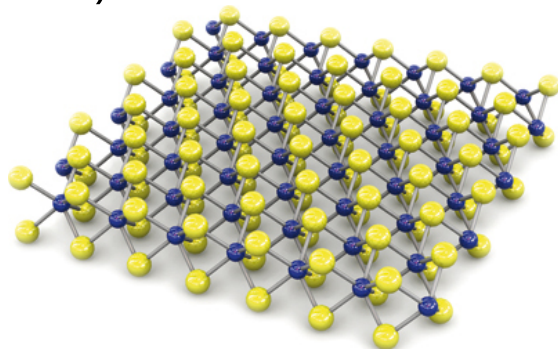
Graphene



hexagonal BN



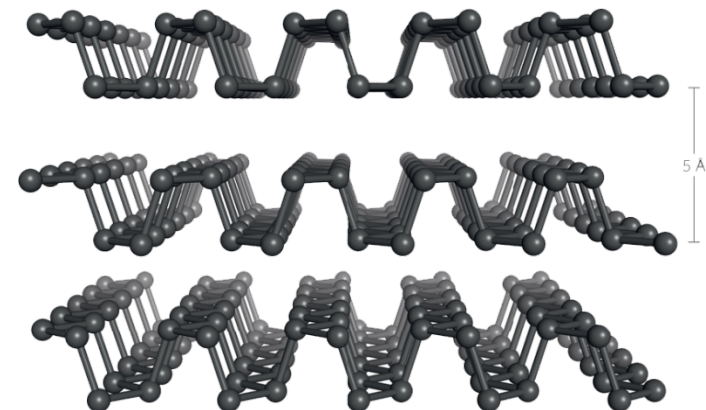
Transition metal dichalcogenides  
(TMD)



$\text{MX}_2$  (M=Mo,W; X=S, Se,Te)

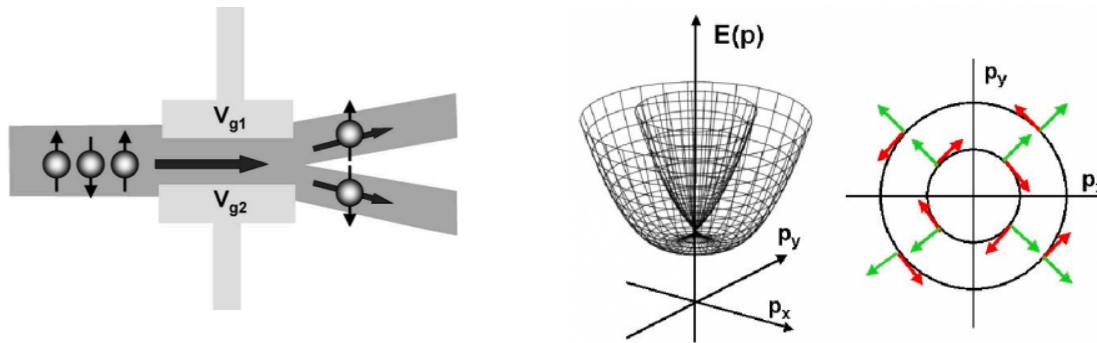


Phosphorene



# Spintronics in 2D materials?

Ex. Spin splitters proposed in conventional 2DEG (Rashba spin-orbit interaction)



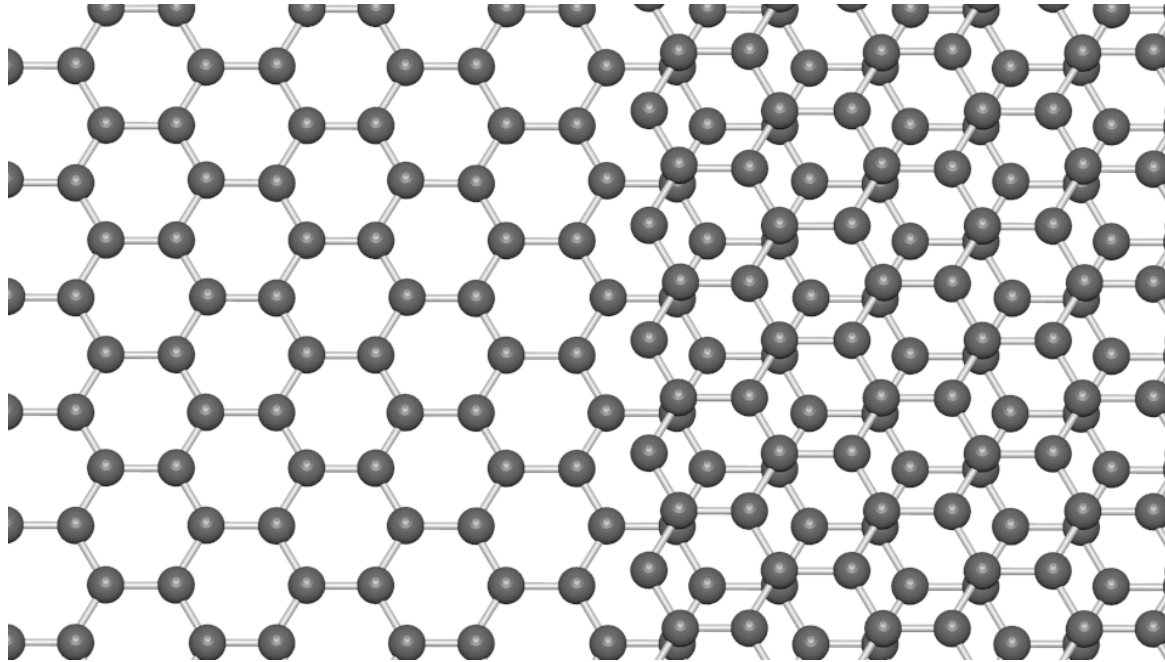
... Not yet realized

Kiselev and Kim, APL 78, 775 (2001).  
Ohe et al, PRB 72, 041308 (2005).  
Yamamoto et al, PRB72, 115321 (2005).

... Can we do this in better way in 2D materials ?

# This talk: **Electron transmission through atomic boundary**

## Graphene

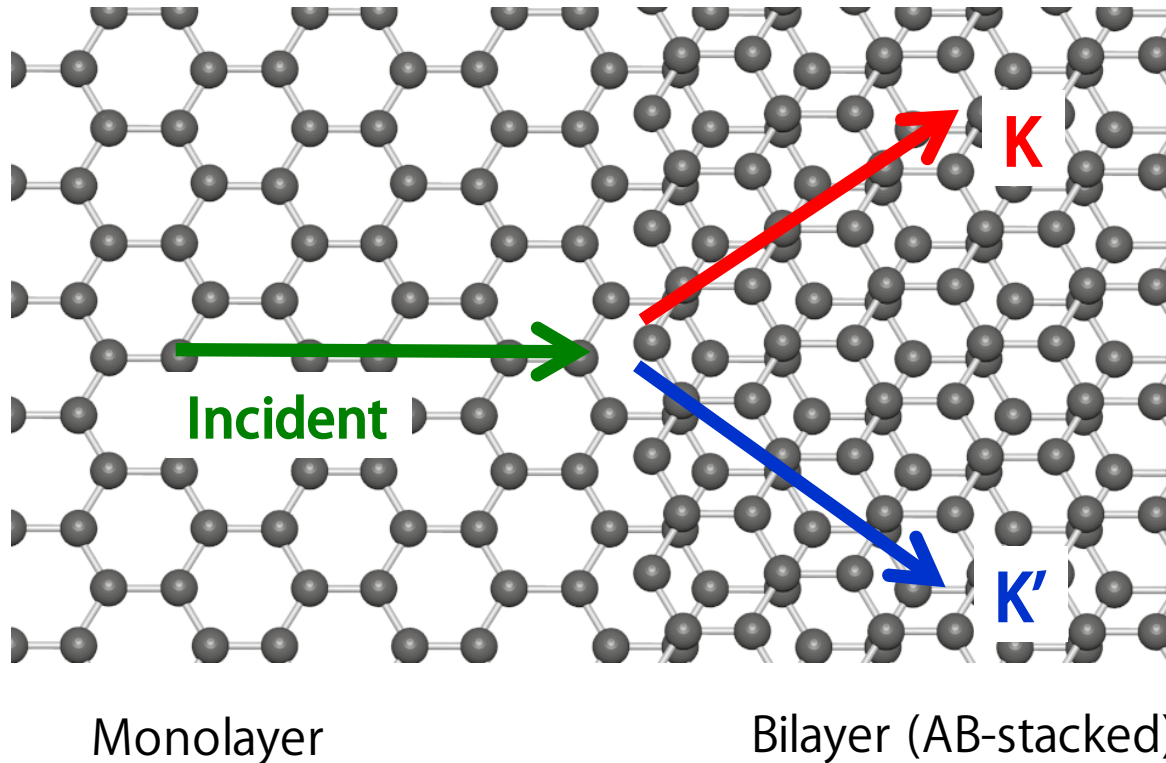


Monolayer

Bilayer (AB-stacked)

# This talk: Electron transmission through atomic boundary

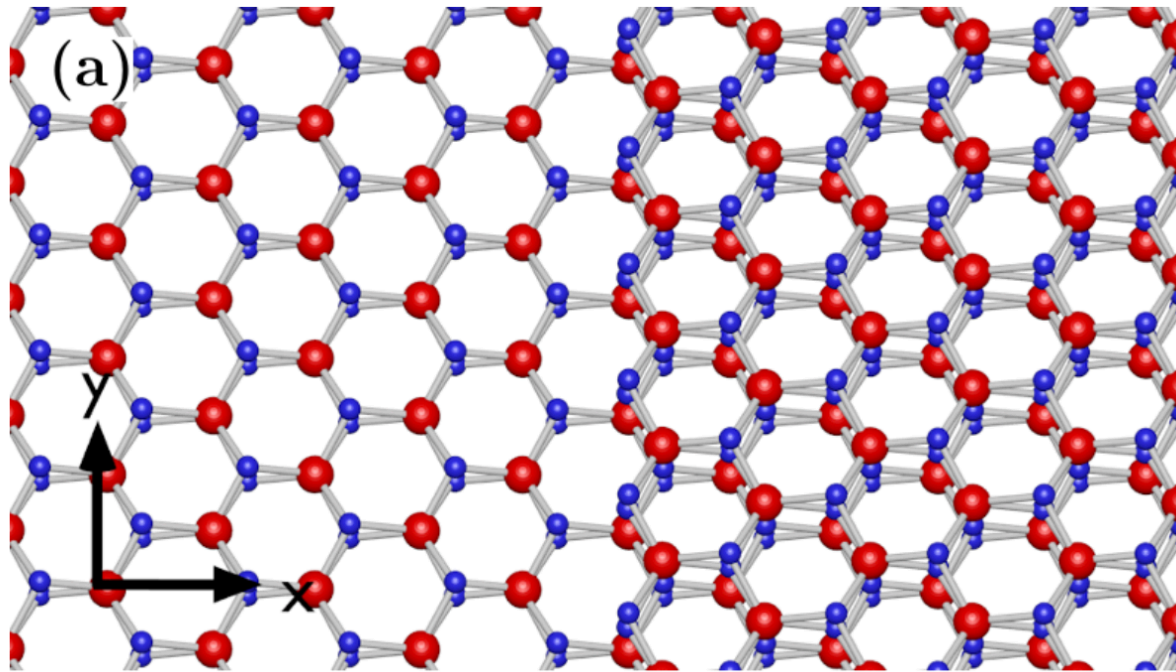
## Graphene



Graphene atomic boundary splits valley pseudo-spins (K,K')

# This talk: Electron transmission through atomic boundary

## Transition metal dichalcogenides (TMD)

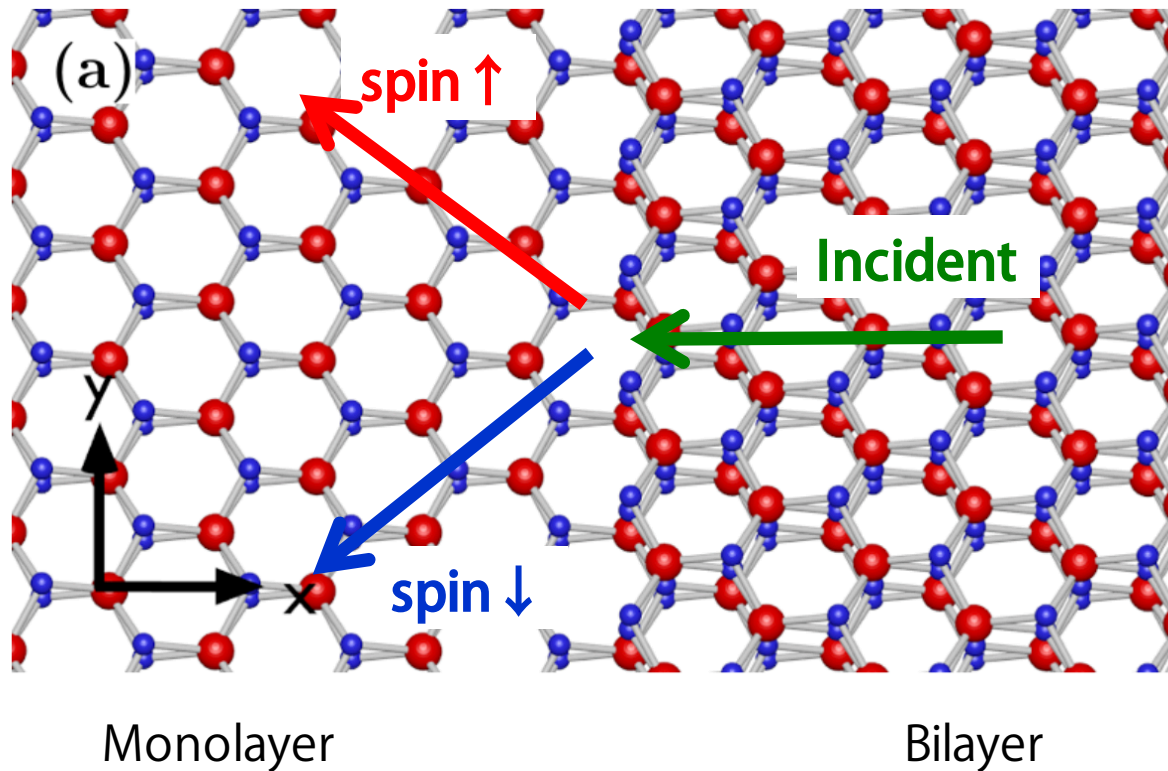


Monolayer

Bilayer

# This talk: Electron transmission through atomic boundary

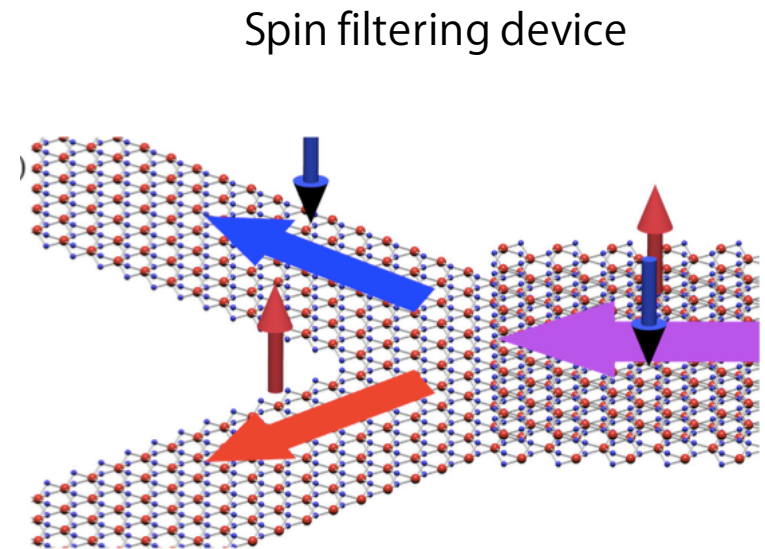
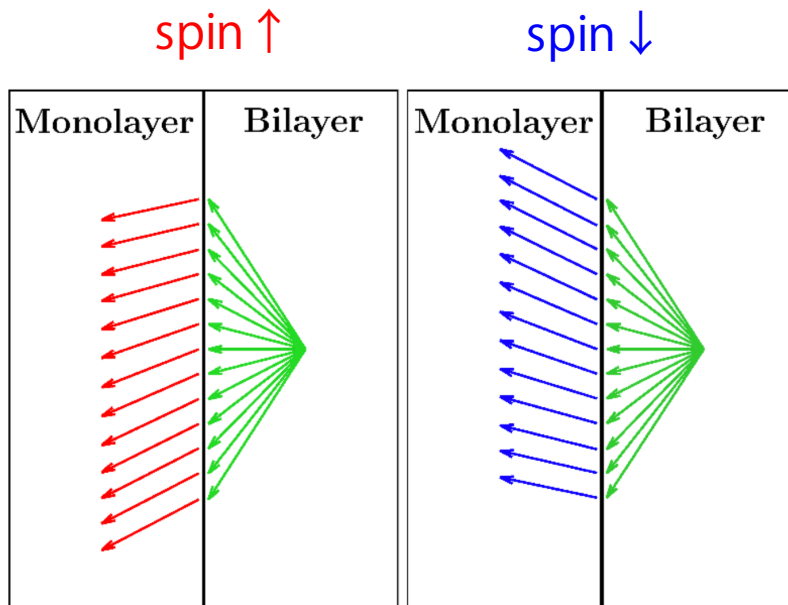
## Transition metal dichalcogenides (TMD)



TMD atomic boundary splits real spins!

# This talk: Electron transmission through atomic boundary

TMD monolayer-bilayer junction works as a **spin splitter**



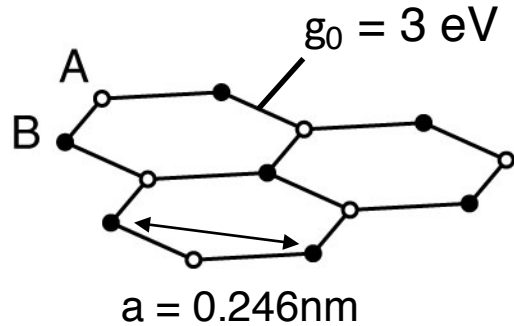
T. Habe and M. Koshino, Phys. Rev. B 91, 201407(R) (2015)



# Graphene's band structure

## Monolayer graphene

McClure, Phys. Rev. 104, 666 (1956).



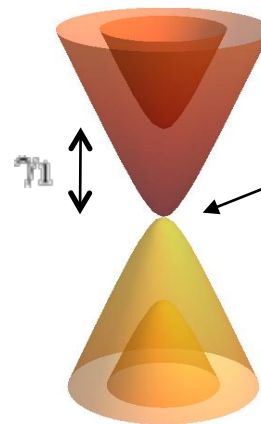
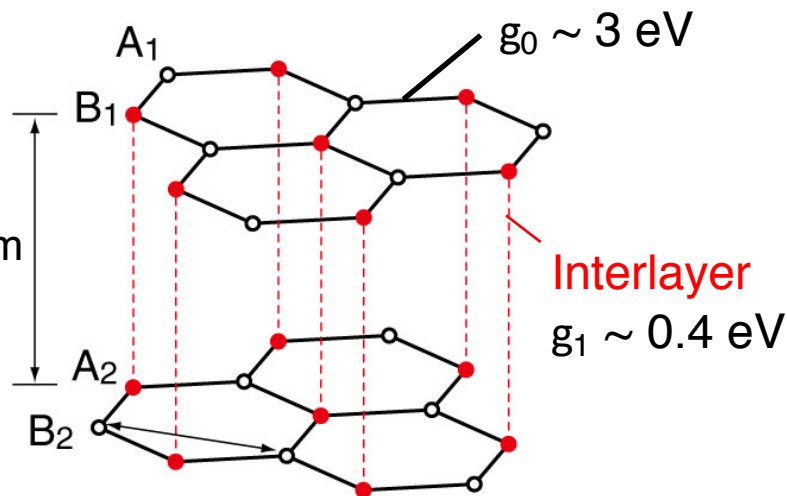
**Linear**

$$v = \frac{\sqrt{3}a}{2\hbar} \gamma_0 \approx 1 \times 10^6 \text{ m/s}$$

.... constant velocity

## Bilayer graphene

McCann and Fal'ko, PRL 96, 086805 (2006)



**Massive**

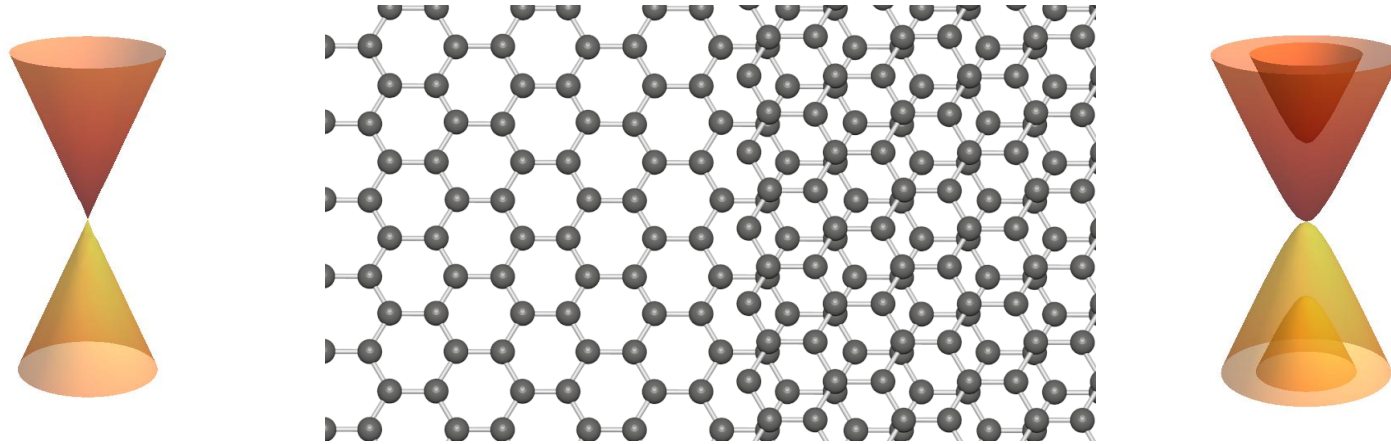
Effective mass:

$$m^* = \frac{\gamma_1}{2v^2} \sim 0.03m_0$$

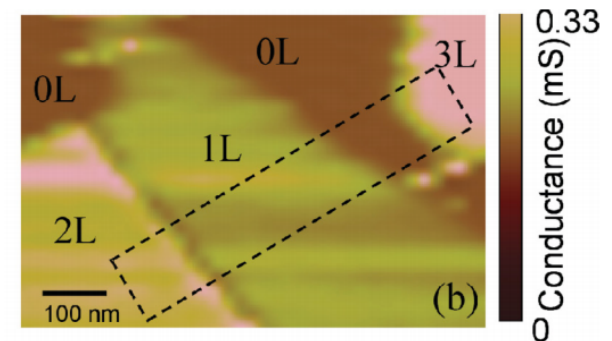
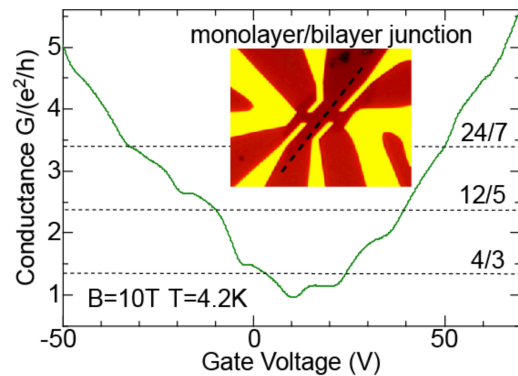
# Graphene monolayer-bilayer junction

## Theoretical studies

Nakanishi, Koshino, Ando, PRB 82, 125428 (2010)  
Koshino, Nakanishi, Ando PRB 82, 205436 (2010)



## Experiments



# Boundary condition in continuum model

**Hamiltonian**

**Monolayer (K-valley)**

$$\mathcal{H}^K = \begin{pmatrix} 0 & vp_- \\ vp_+ & 0 \end{pmatrix}$$

$$p_{\pm} = p_x \pm ip_y$$

**Wavefunction**

$$\mathbf{F}^K(\mathbf{r}) = \begin{pmatrix} F_A^K(\mathbf{r}) \\ F_B^K(\mathbf{r}) \end{pmatrix}$$

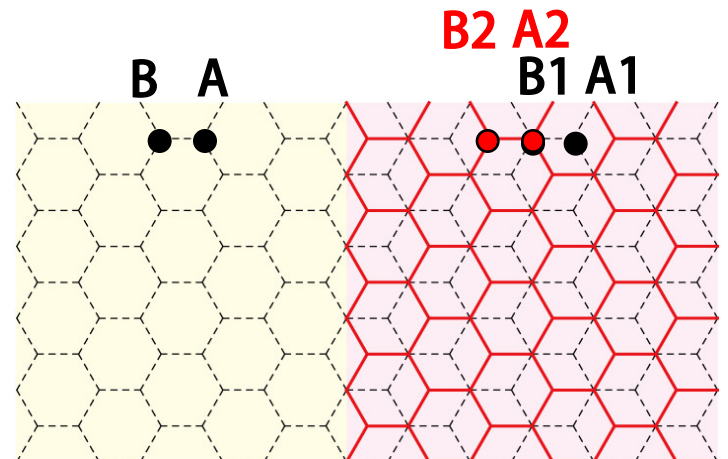
**Bilayer (K-valley)**

$$\mathcal{H}^K = \begin{pmatrix} 0 & vp_- & 0 & 0 \\ vp_+ & 0 & \gamma_1 & 0 \\ 0 & \gamma_1 & 0 & vp_- \\ 0 & 0 & vp_+ & 0 \end{pmatrix}.$$

$$\mathbf{G}^K(\mathbf{r}) = \begin{pmatrix} G_{A1}^K(\mathbf{r}) \\ G_{B1}^K(\mathbf{r}) \\ G_{A2}^K(\mathbf{r}) \\ G_{B2}^K(\mathbf{r}) \end{pmatrix}$$

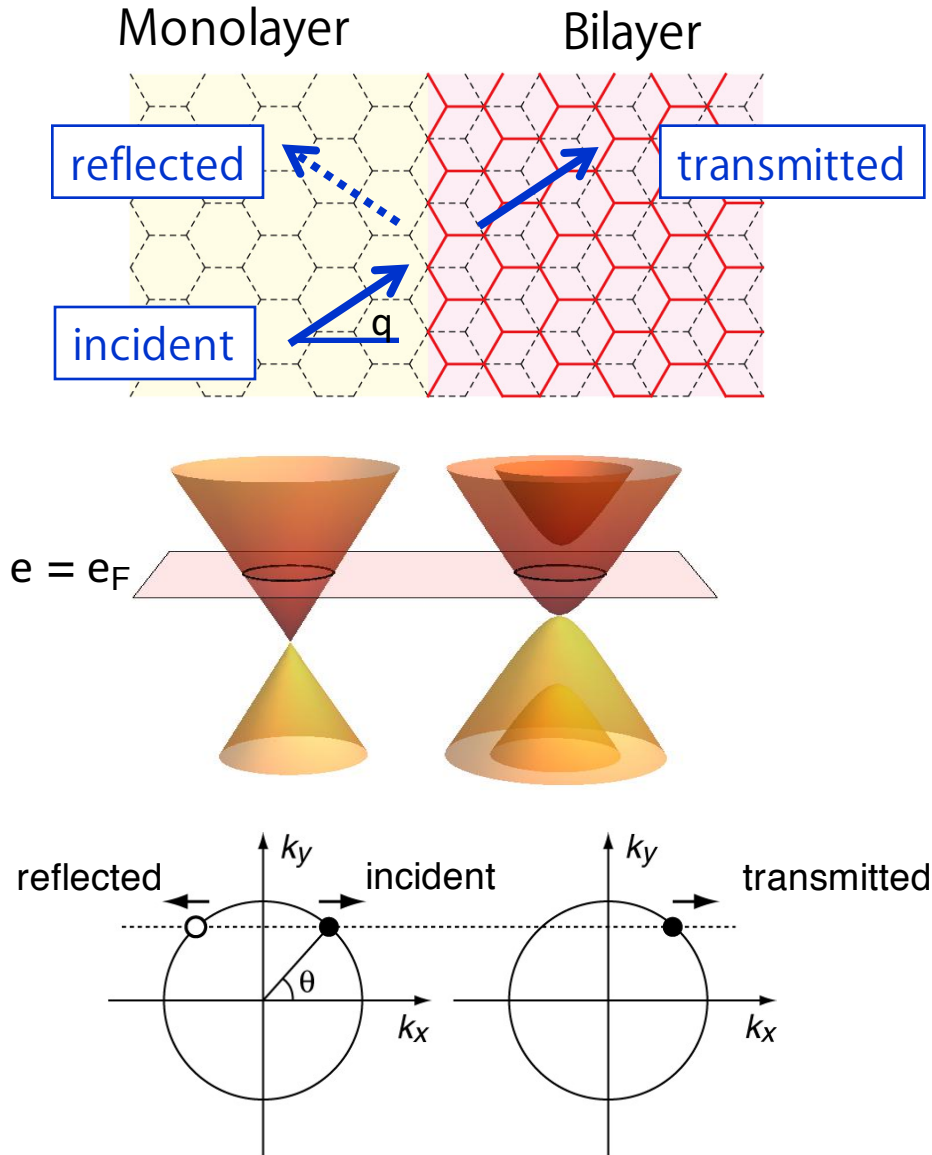
**Boundary condition (Zigzag-1)**

$F_A^K(0, y) = G_{A1}^K(0, y)$	}	<b>1st layer</b>
$F_B^K(0, y) = G_{B1}^K(0, y)$		
$0 = G_{B2}^K(0, y)$	—	<b>2nd layer</b>

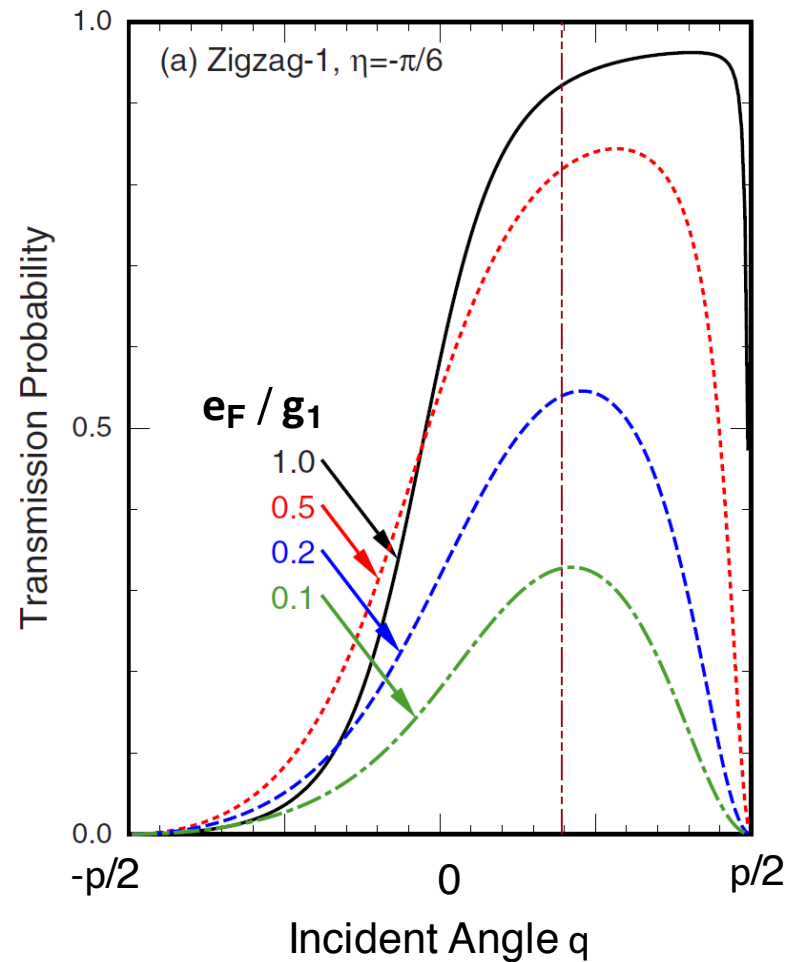


# Transmission through M-B junction

T. Nakanishi, MK, T. Ando  
PRB 82, 125428 (2010)

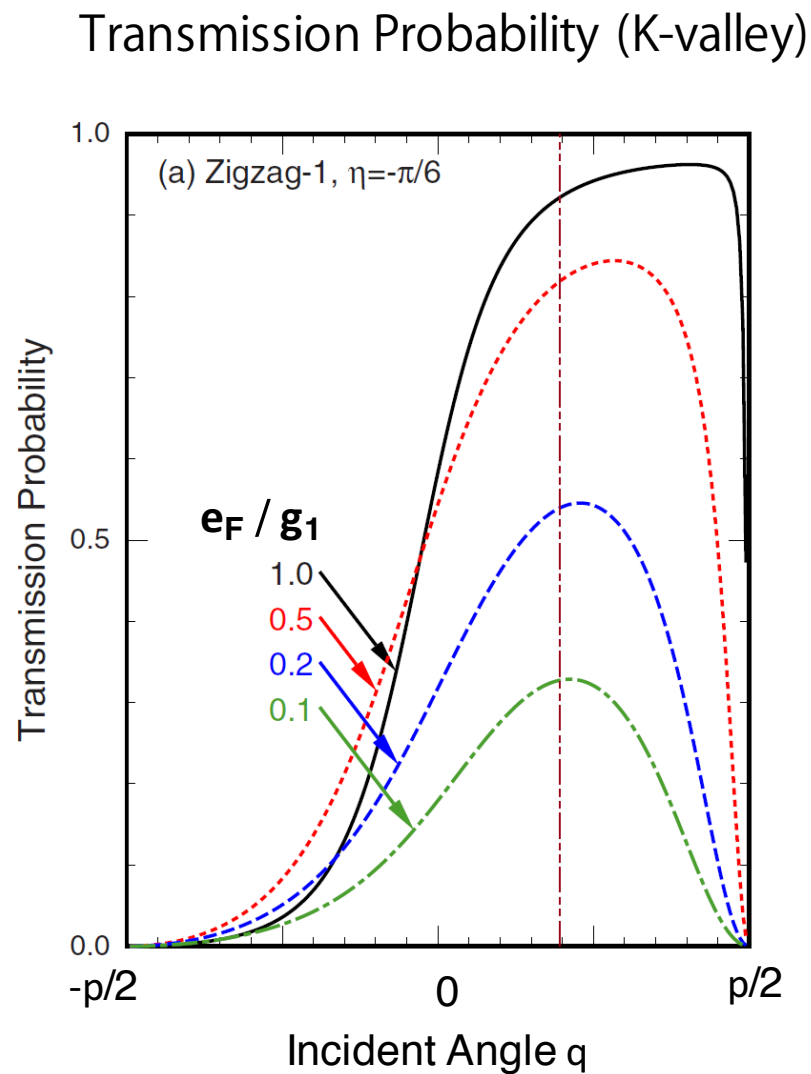
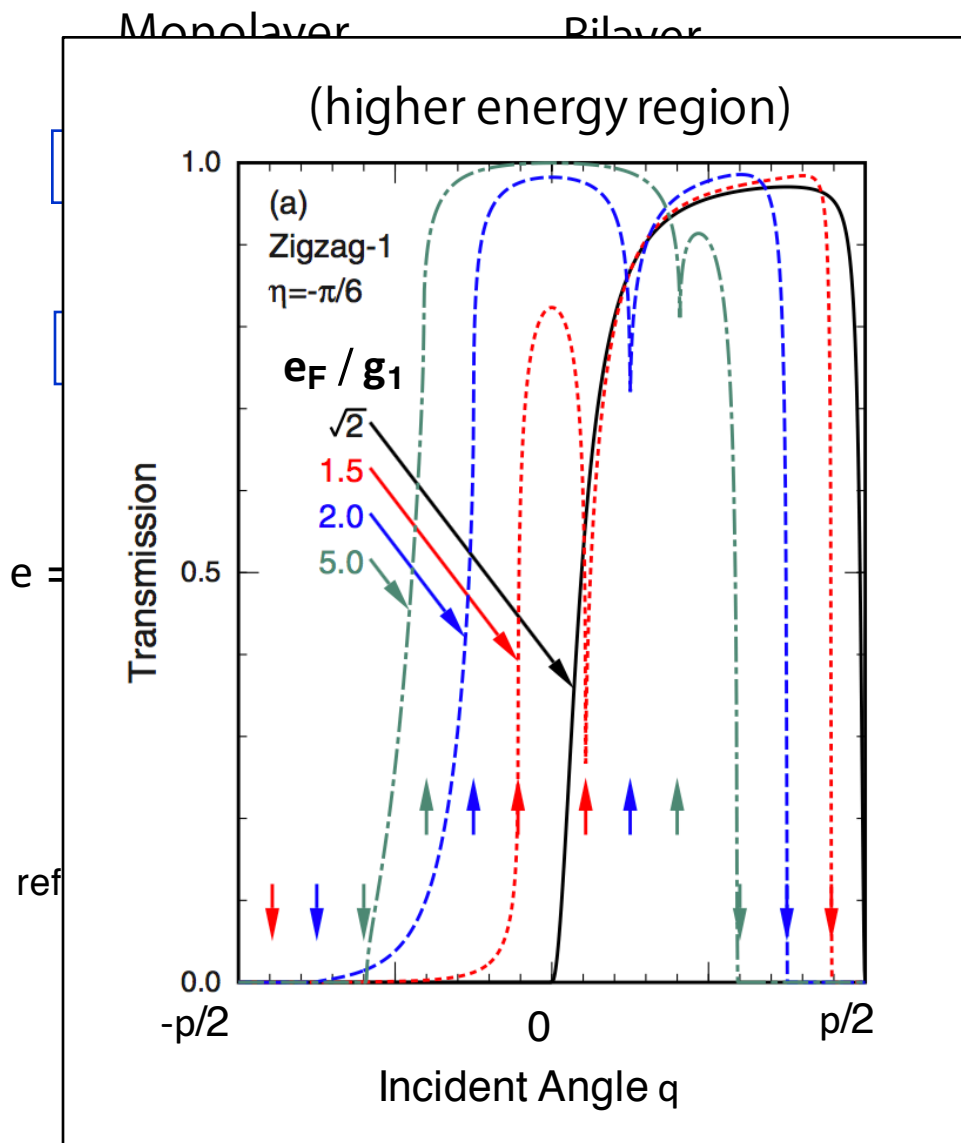


Transmission Probability (K-valley)

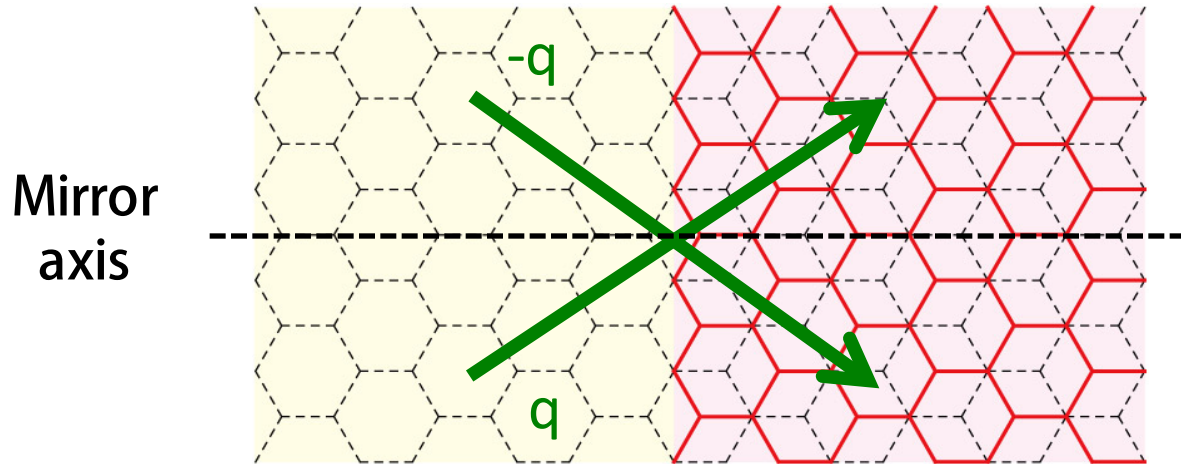


# Transmission through M-B junction

T. Nakanishi, MK, T. Ando  
PRB 82, 125428 (2010)



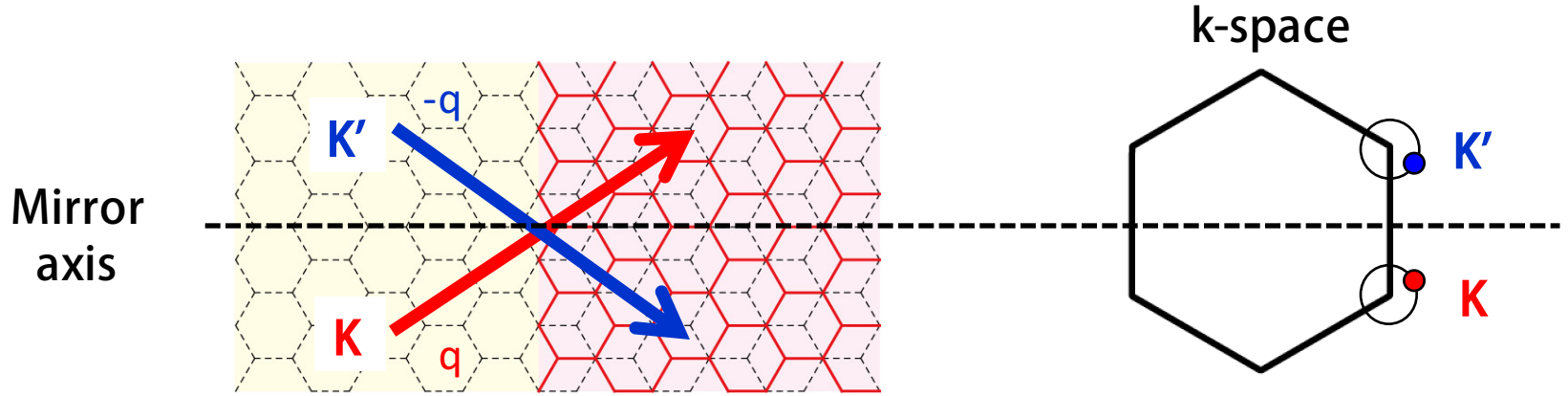
# Why angle-asymmetric?



**Mirror symmetry**

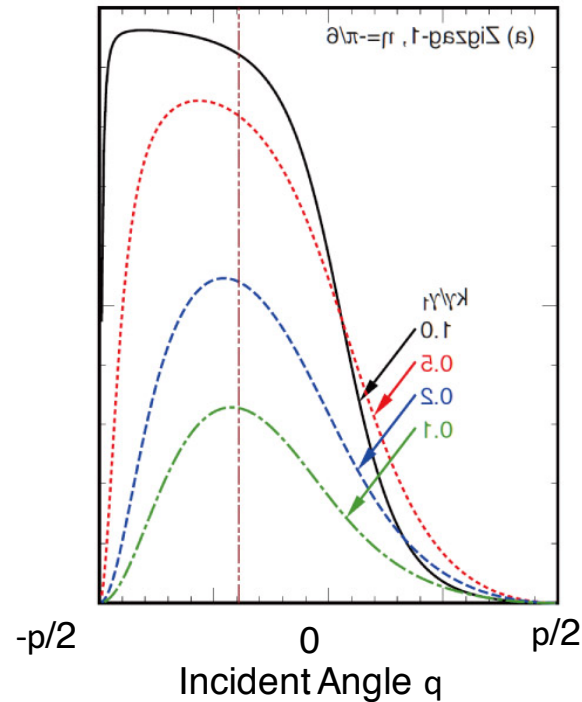
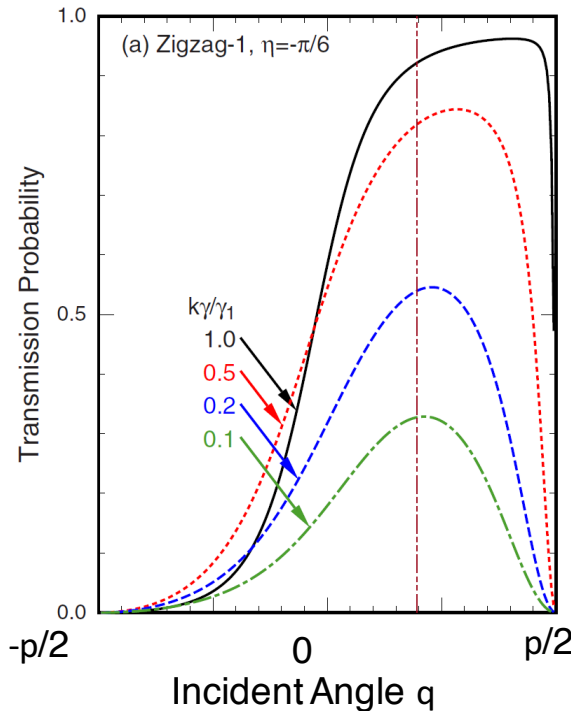
... Transmission should be symmetric with respect to  $q = 0$ ?

# Why angle-asymmetric?



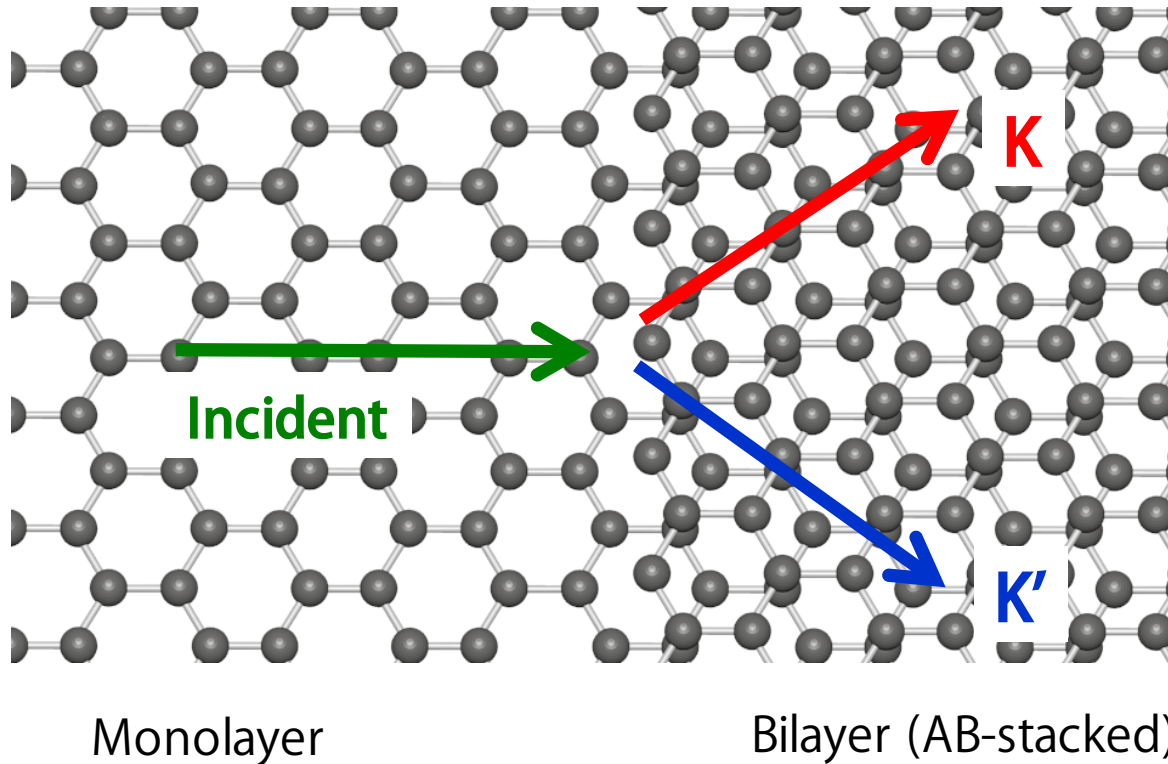
**K-valley**

**K'-valley**



**Asymmetric transmission in K and K'**

# Graphene

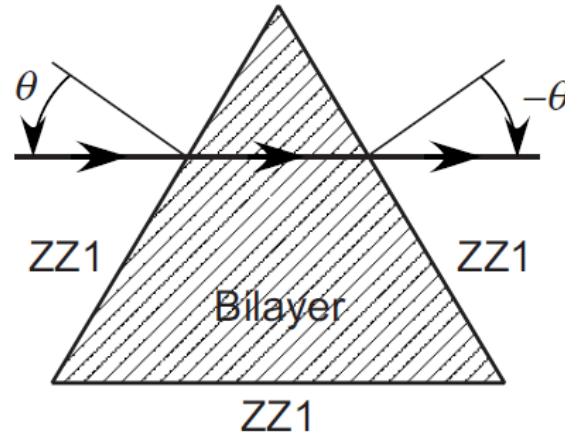


Graphene atomic boundary splits valley pseudo-spins (K,K')

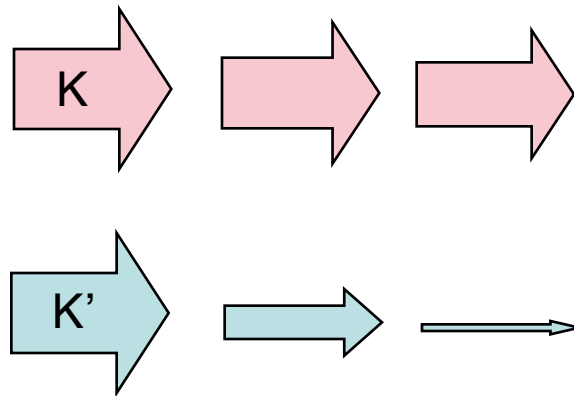


# Valley-polarizing mechanism

T. Nakanishi, MK, T. Ando  
PRB 82, 125428 (2010)



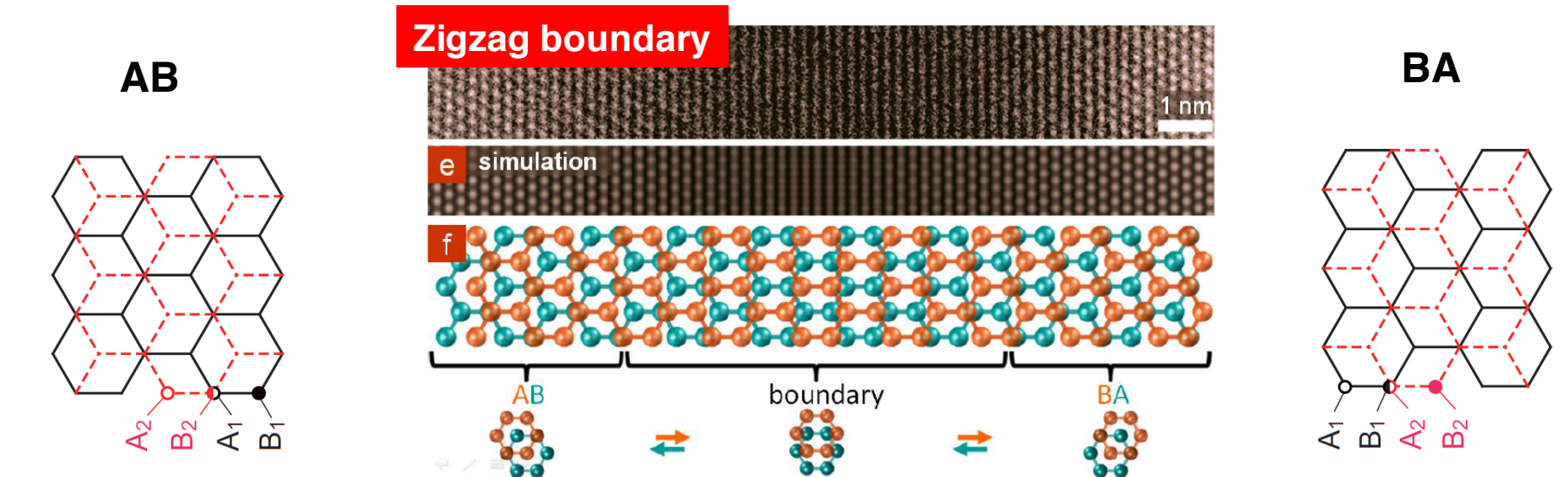
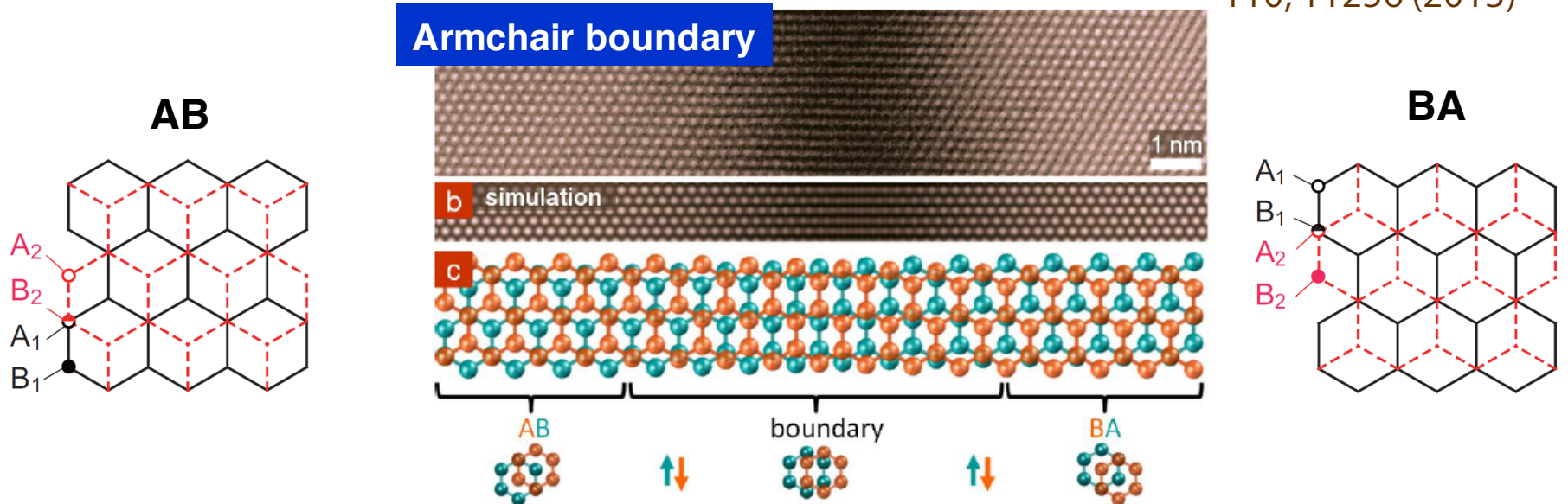
Monolayer



Valley polarization

# AB-BA domain in bilayer graphene

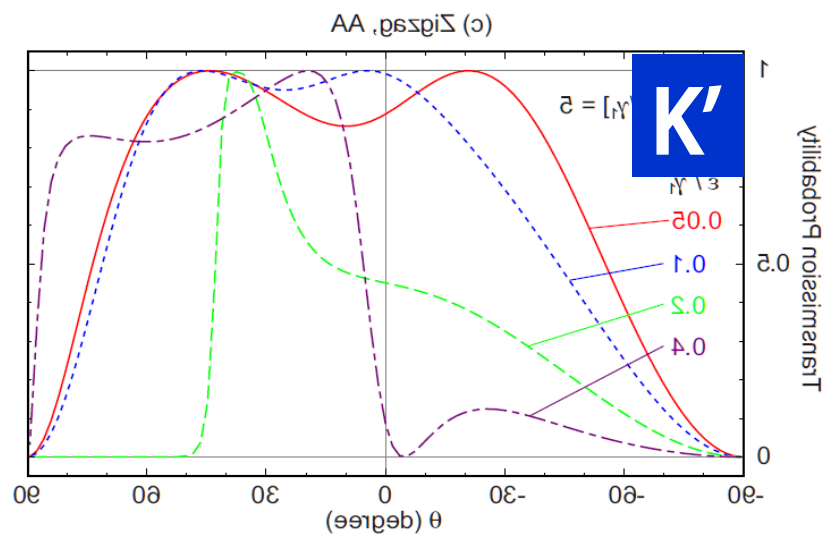
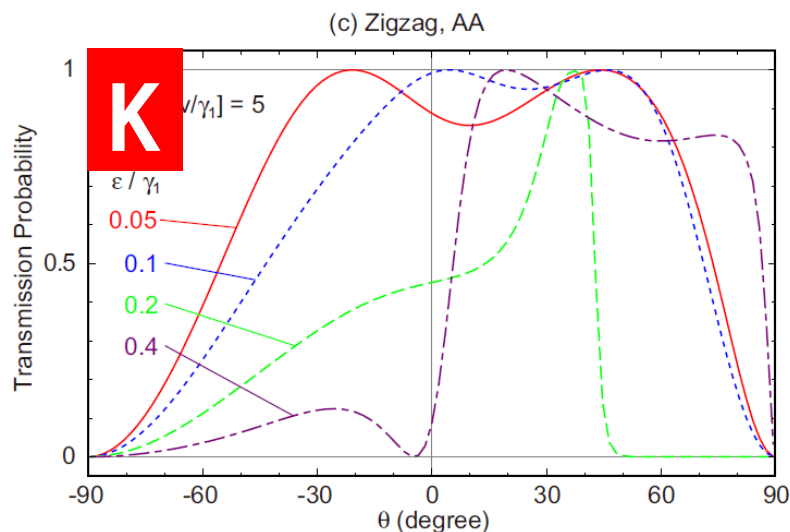
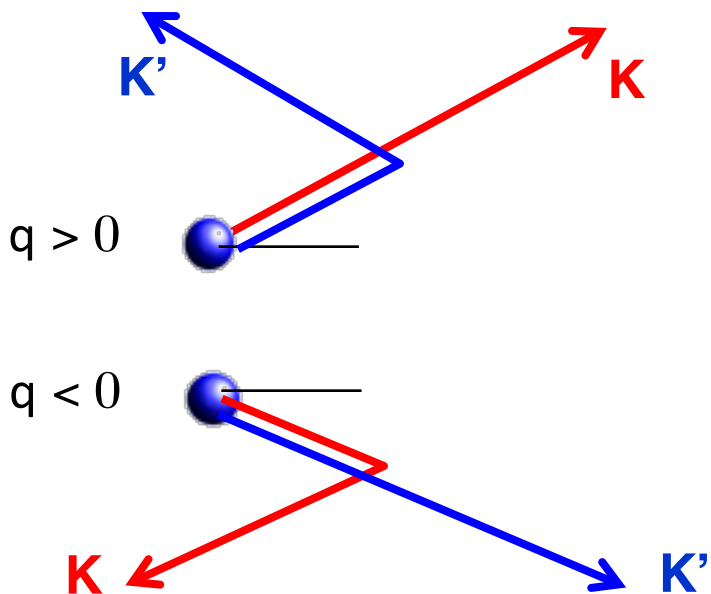
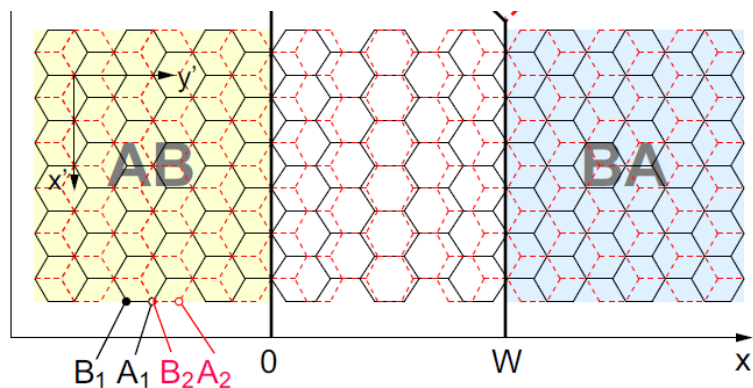
J. S. Alden et al,  
PNAS  
110, 11256 (2013)



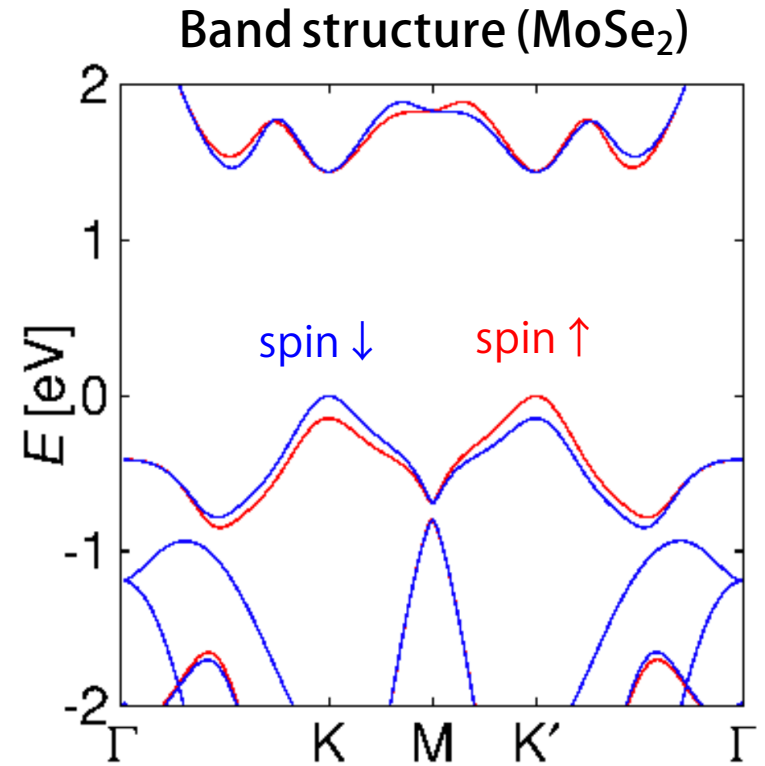
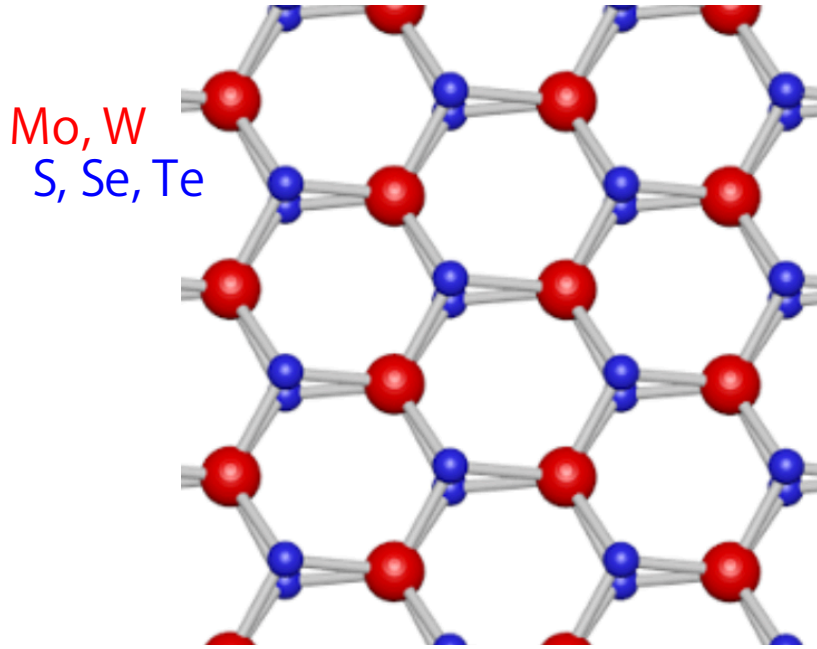
# Valley polarization

M. Koshino,  
Phys. Rev. B 88, 115409 (2013)

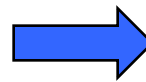
Reflection symmetry:  $K$  and  $K'$  have opposite angle dependence



# TMD (Transition metal dichalcogenides)



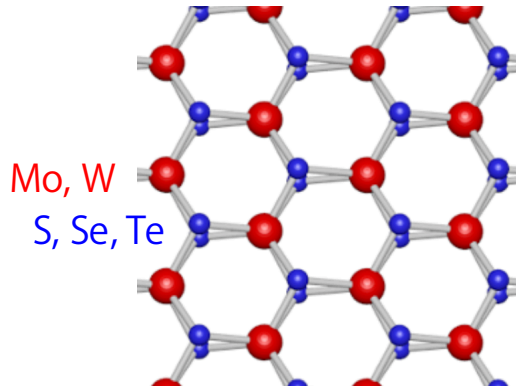
Inversion symmetry  
breaking  
+  
Spin-orbit coupling



Spin split  
(opposite in K and K')

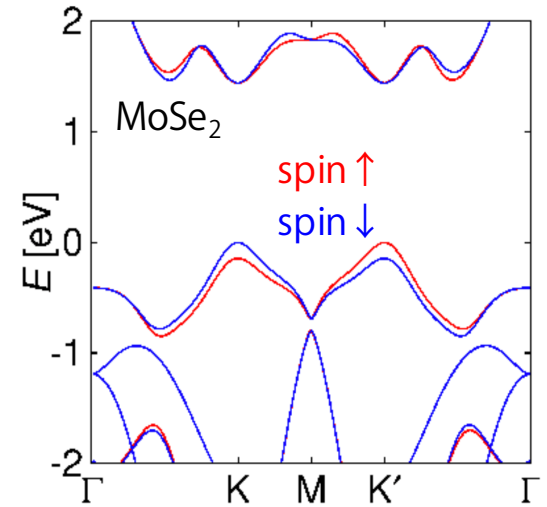
# TMD (Transition metal dichalcogenides)

## Monolayer

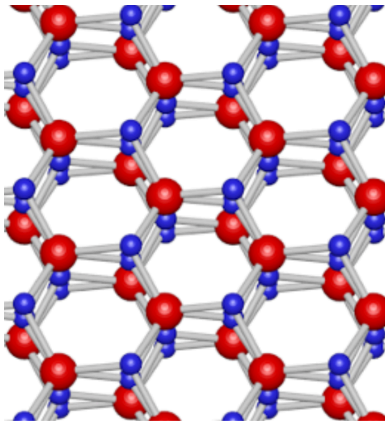


Inversion symmetry  
breaking  
+  
Spin-orbit coupling

↓  
Spin split

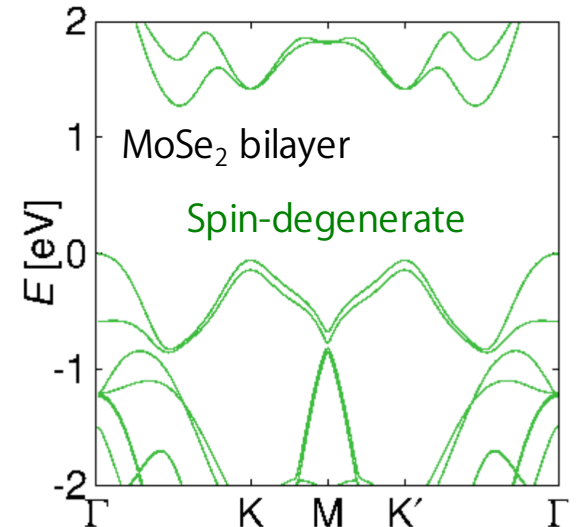


## Bilayer (2H phase)



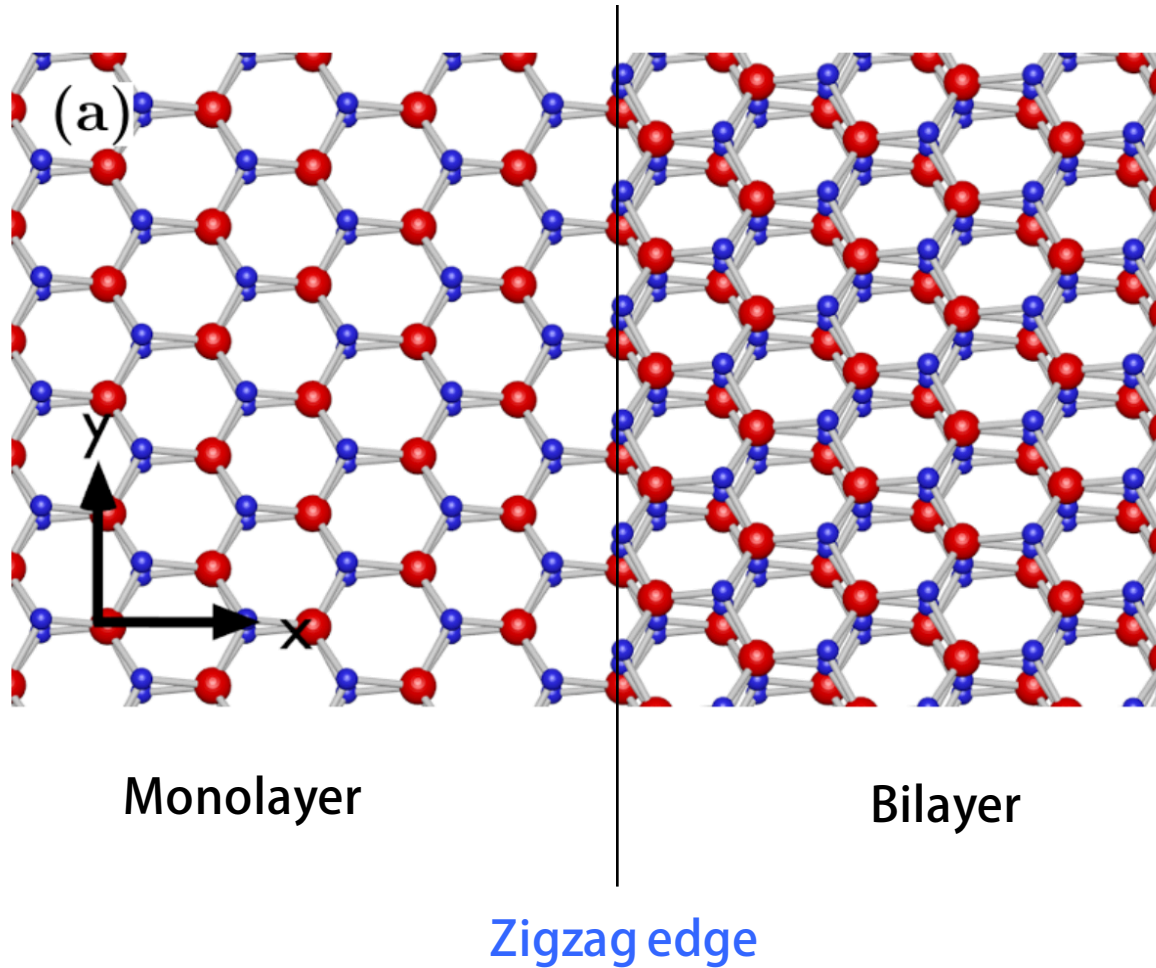
Inversion symmetry  
recovers

↓  
Spin degenerate



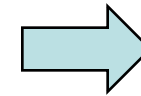
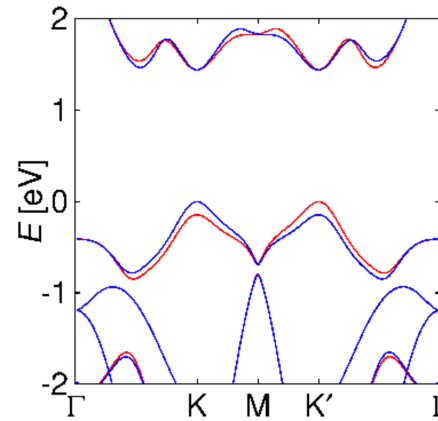
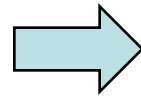
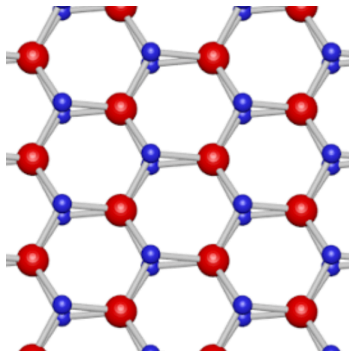
# TMD atomic junction

T. Habe and M. Koshino,  
Phys. Rev. B 91, 201407(R) (2015)

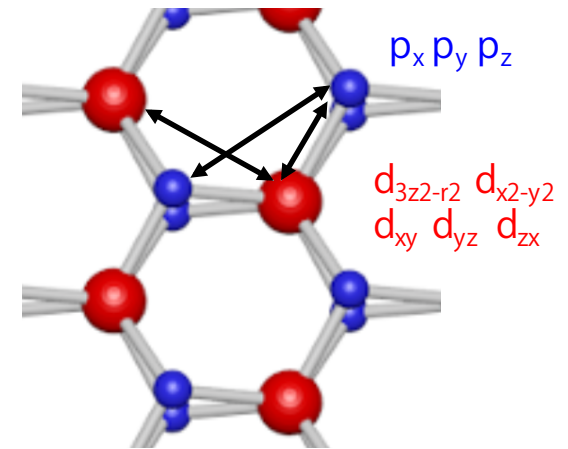


# Theoretical Method

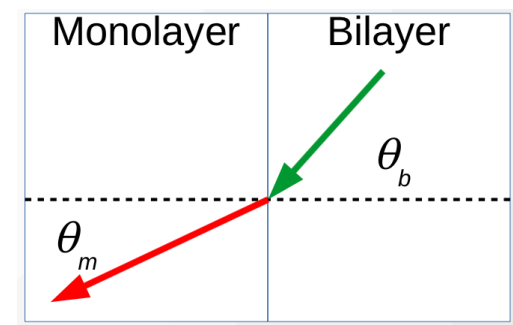
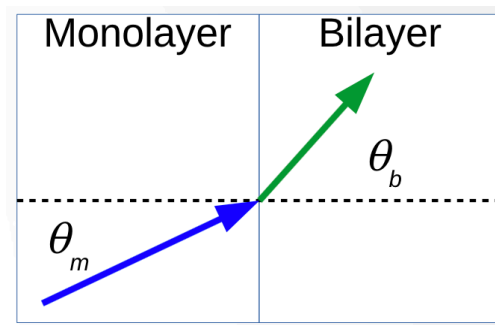
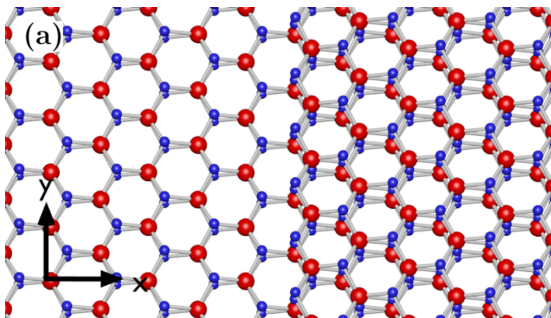
i) **DFT** band calculation  
(Quantum Espresso)



ii) Create **tight-binding model**  
(Wannier 90)



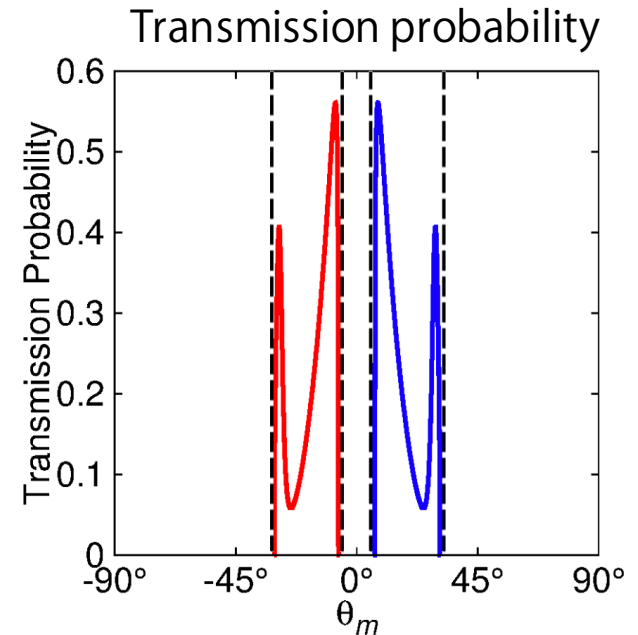
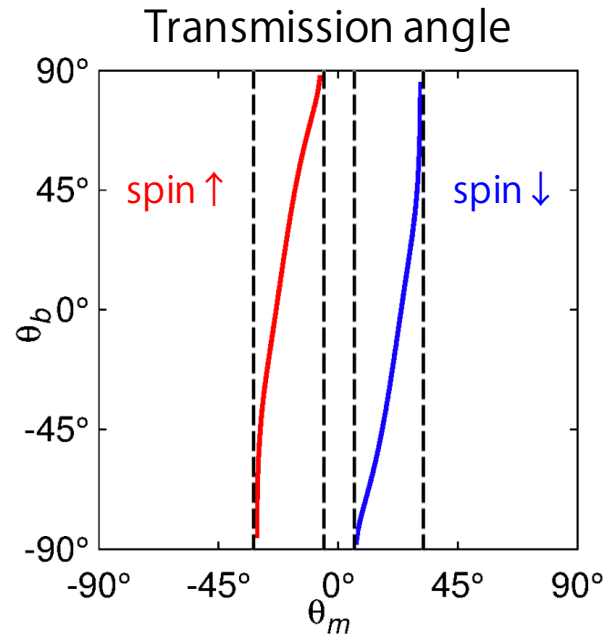
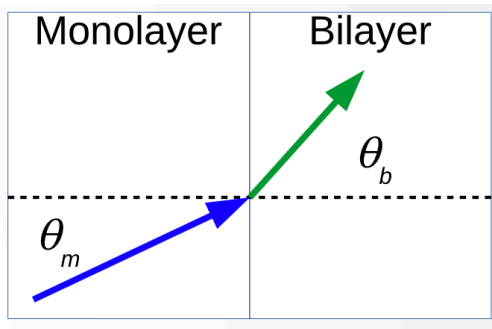
iii) Calculate **the transmission probability**



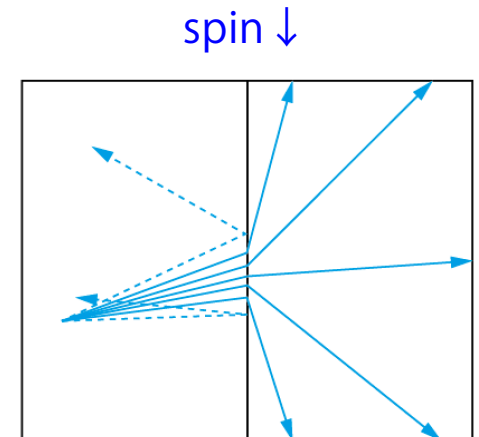
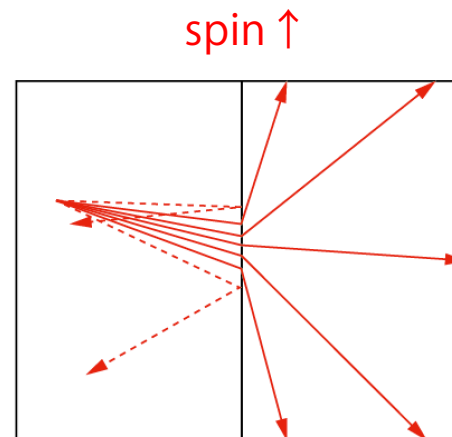
# Electron transmission: MoTe<sub>2</sub>

[hole-doped;  
 $n = -7.02 \times 10^{13} \text{ cm}^{-2}$ ]

From monolayer side



Transmission is highly angle-selective

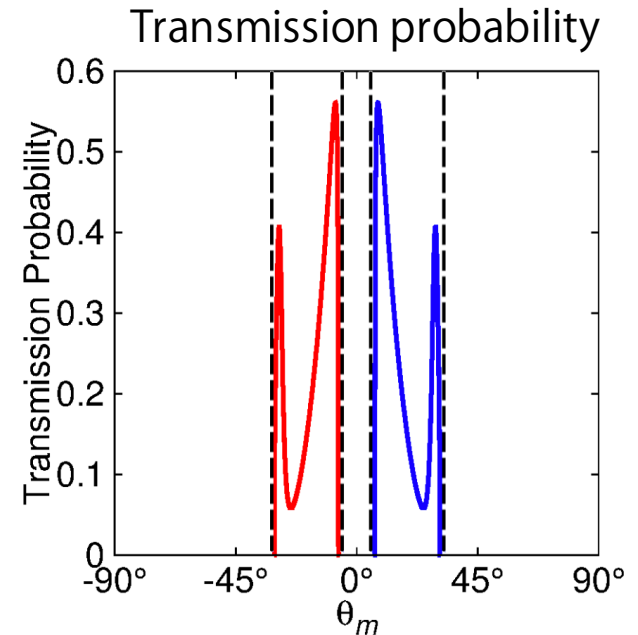
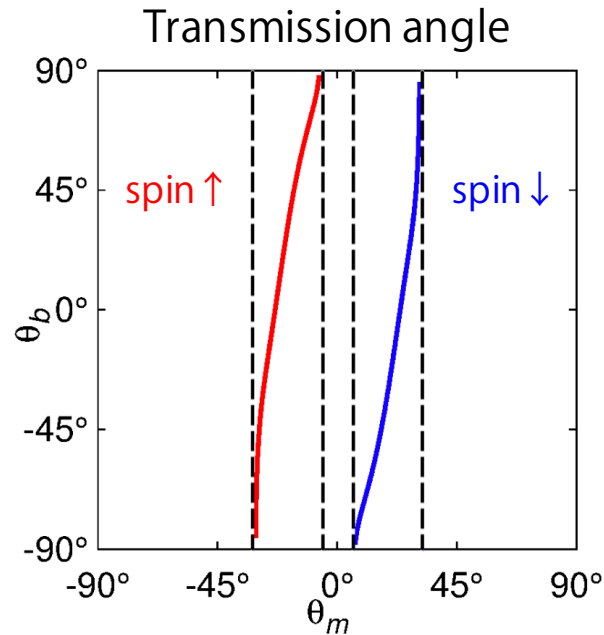
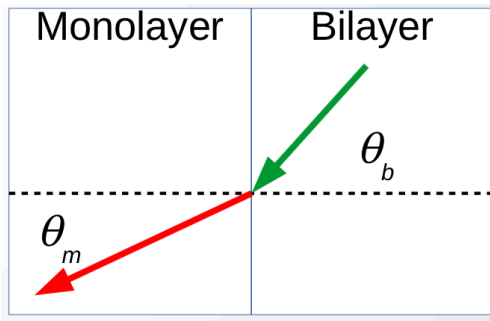




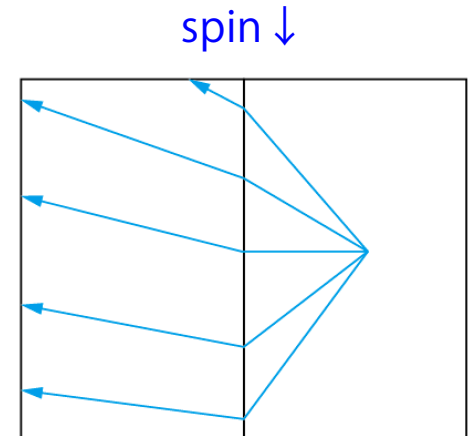
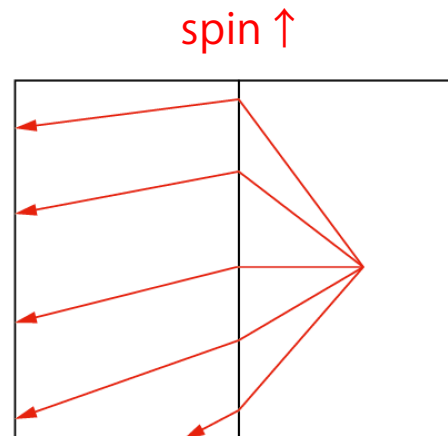
# Electron transmission: MoTe<sub>2</sub>

[hole-doped;  
 $n = -7.02 \times 10^{13} \text{ cm}^{-2}$ ]

From bilayer side

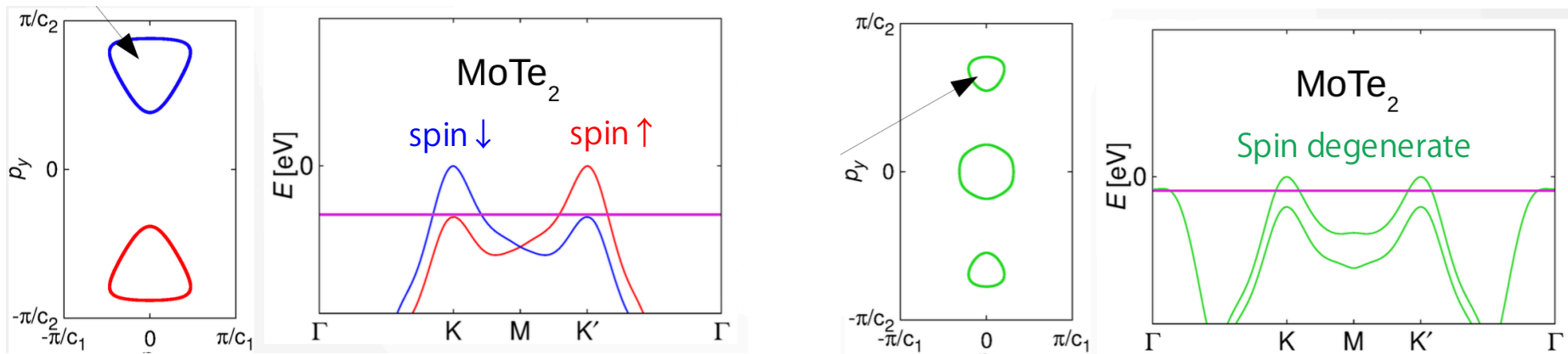


Collimated to different directions depending on spin!



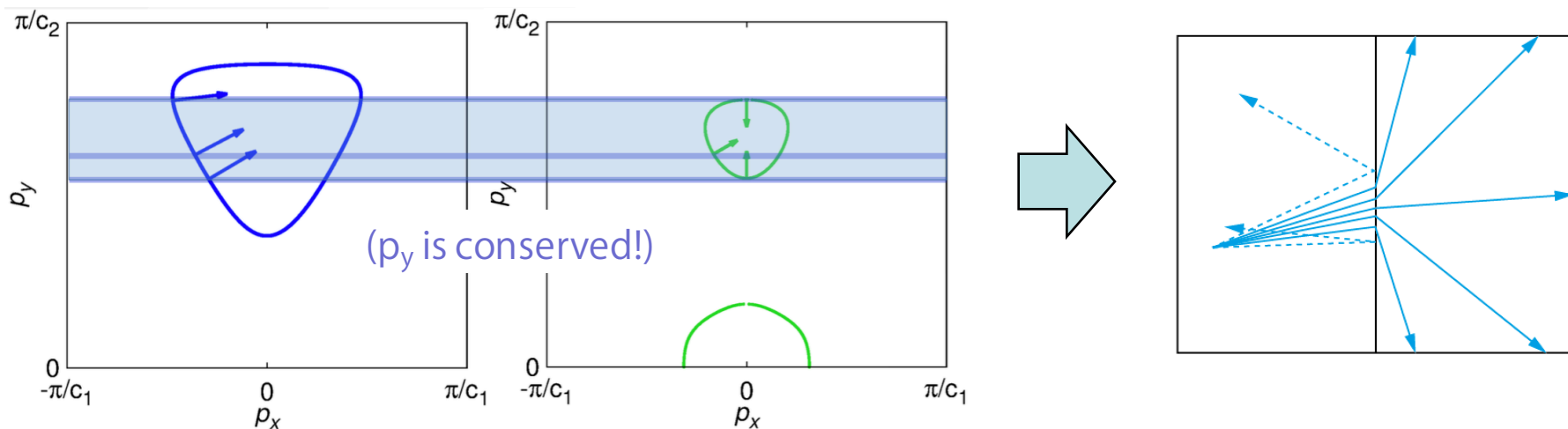
# Mechanism of spin-dependent transmission

Monolayer ← (equal carrier density) → Bilayer



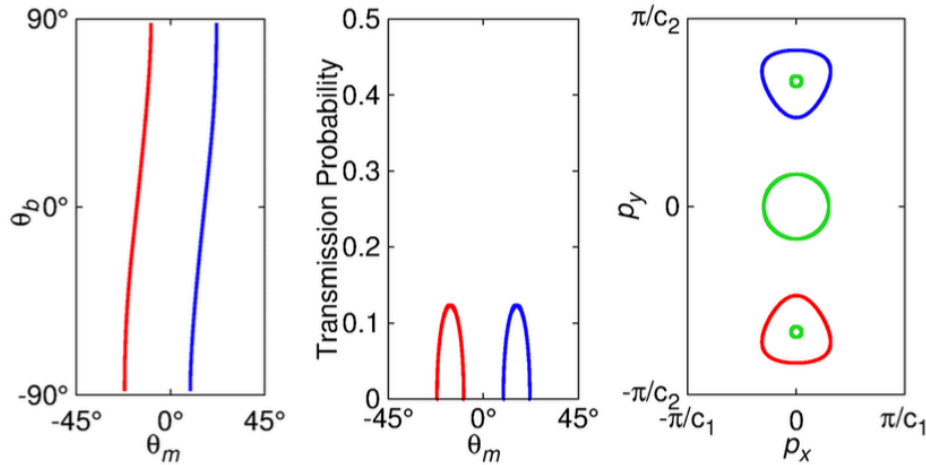
... Larger Fermi circle in monolayer than in bilayer

Wave number matching

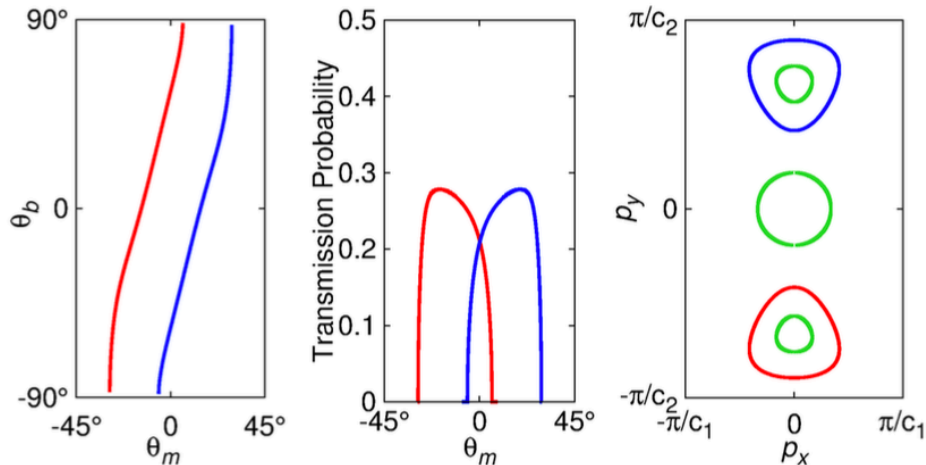


# Other TMDs

(a) MoSe<sub>2</sub> ( $n = 4.80 \times 10^{13} [\text{cm}^{-2}]$ )



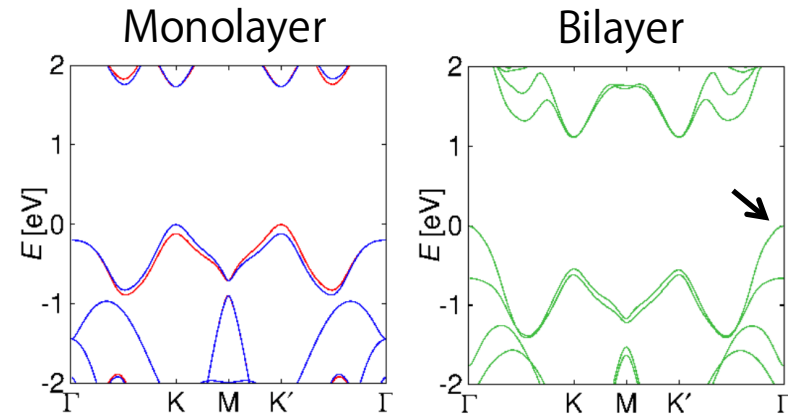
(b) WSe<sub>2</sub> ( $n = 8.85 \times 10^{13} [\text{cm}^{-2}]$ )



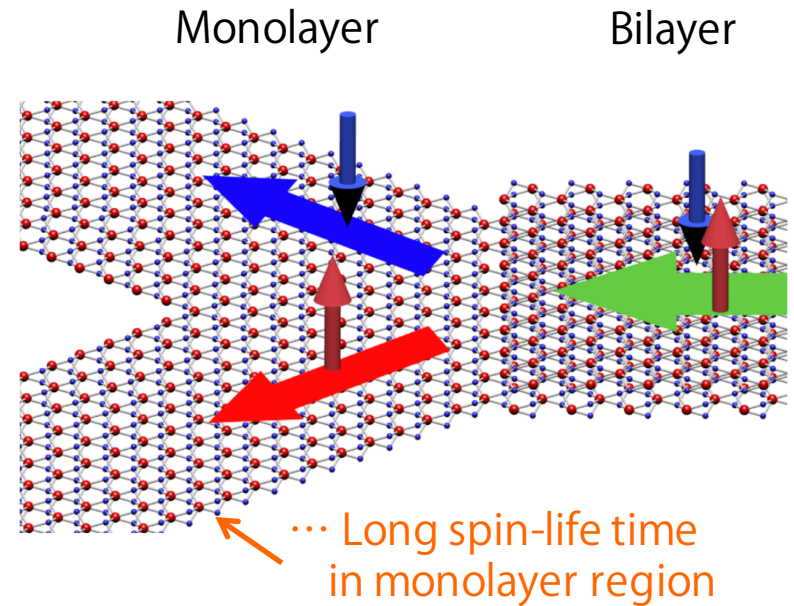
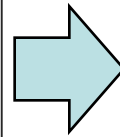
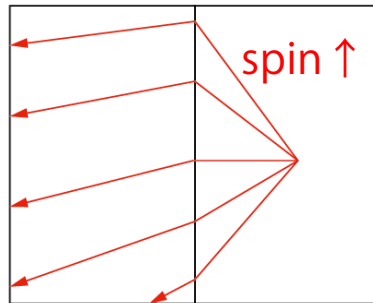
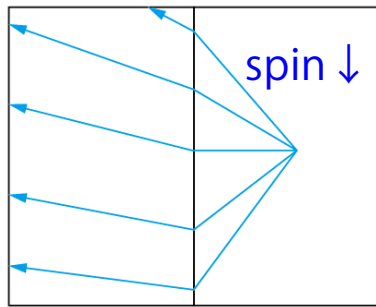
... Similar effects

... MoS<sub>2</sub>?

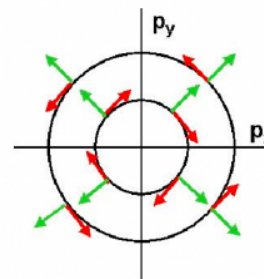
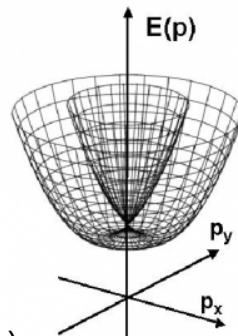
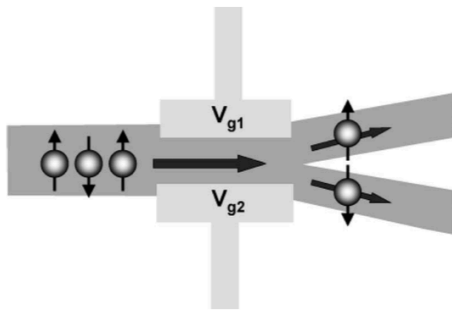
**NO transmission** at the boundary  
due to too high G point



# Spin splitter



Cf. spin splitters proposed in conventional 2DEG (Rashba spin-orbit interaction)



... Spin is not a conserved quantity under impurity scattering

Kiselev and Kim, APL 78, 775 (2001).

Ohe et al, PRB 72, 041308 (2005).

Yamamoto et al, PRB72, 115321 (2005).

# Summary

- Atomic boundary in 2D material causes **flavor-dependent electron transmission**
- Graphene mono-bi junction splits **valley pseudospins**
- **TMD** ( $\text{MoSe}_2$ ,  $\text{MoTe}_2$ ,  $\text{WSe}_2$ ,  $\text{WTe}_2$ ) mono-bi junction splits **real spins**
- Possible application to spin-filtering devices

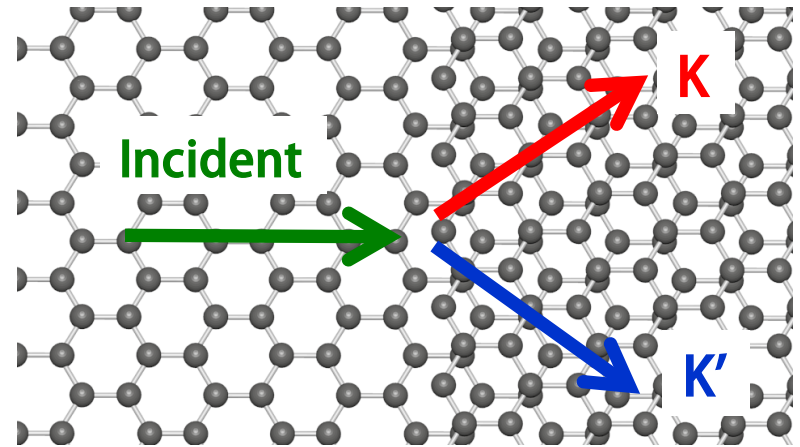
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## Acknowledgements

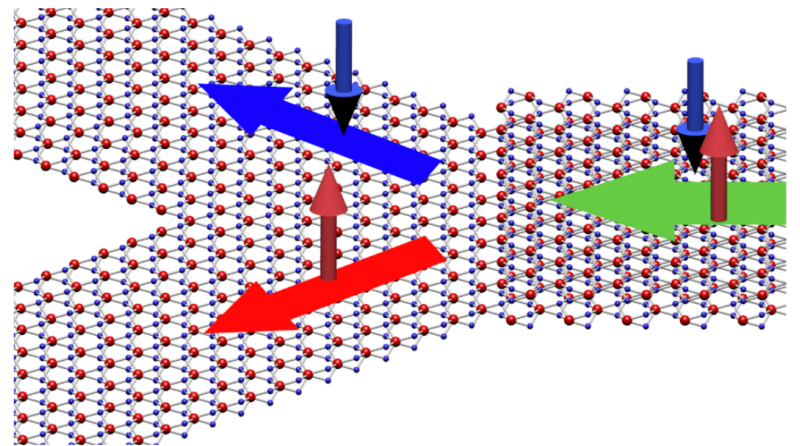
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# Dependence on edge configuration

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