

Many-Body Perturbation Theory:

(2) Bethe-Salpeter Equation and Optical Excitations

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- Excited states: electrons, holes and excitons
- Equation of motion
- Electron-hole correlation
- Examples

States of a many-electron system

N Electrons

$$|N, 0\rangle$$

1. Ground state $|N, 0\rangle$:

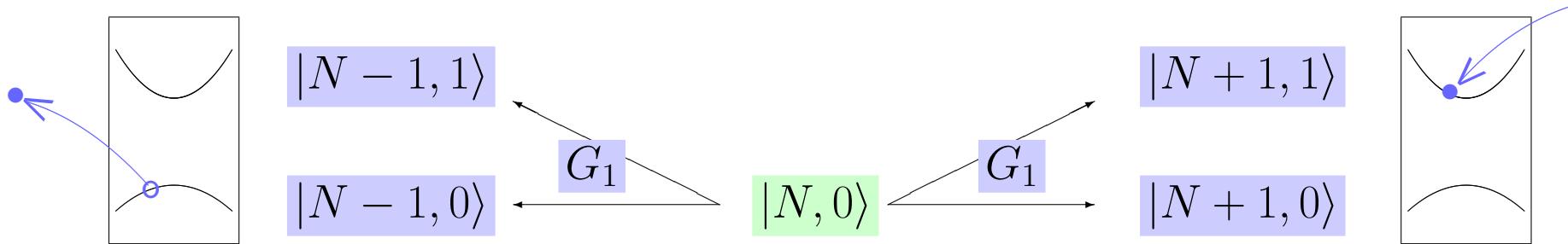
Density-functional theory \implies Geometry

States of a many-electron system

$N - 1$ Electrons

N Electrons

$N + 1$ Electrons



G_1 = Single-particle Green function

1. Ground state $|N, 0\rangle$:

Density-functional theory \Rightarrow Geometry

Many-body perturbation theory:

2. $|N, 0\rangle \rightarrow |N \pm 1, m\rangle$:

$$G_1(\mathbf{x}t, \mathbf{x}'t') = -i \langle N, 0 | T (\psi(\mathbf{x}, t)\psi^+(\mathbf{x}', t')) | N, 0 \rangle$$

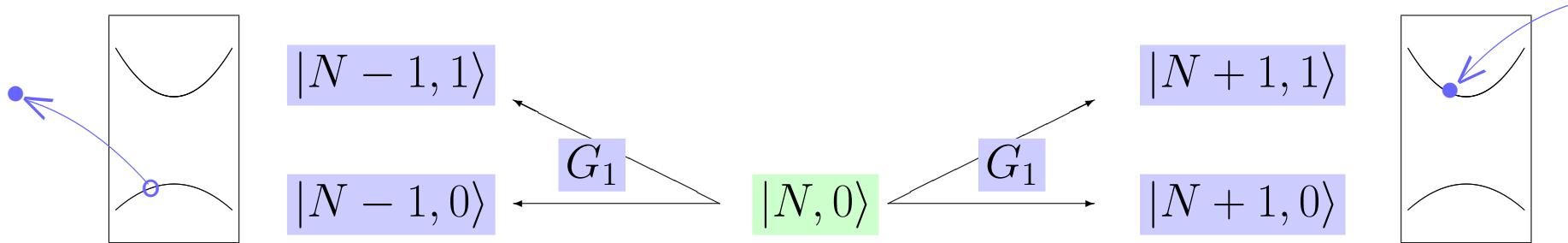
(Propagation of electron or hole between (\mathbf{x}, t) and (\mathbf{x}', t'))

States of a many-electron system

$N - 1$ Electrons

N Electrons

$N + 1$ Electrons



G_1 = Single-particle Green function

1. Ground state $|N, 0\rangle$:

Density-functional theory \Rightarrow Geometry

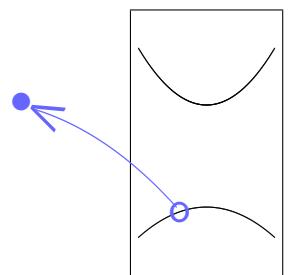
Many-body perturbation theory:

2. $|N, 0\rangle \rightarrow |N \pm 1, m\rangle$:

Dyson equation for G_1 \Rightarrow Band structure

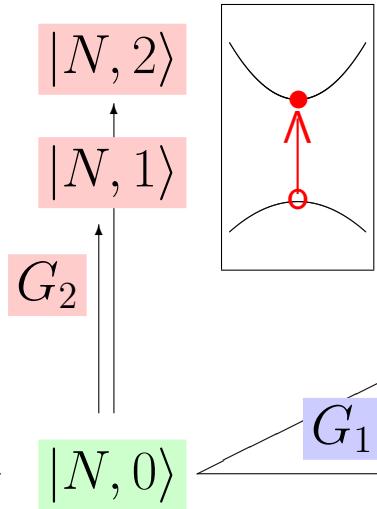
States of a many-electron system

$N - 1$ Electrons

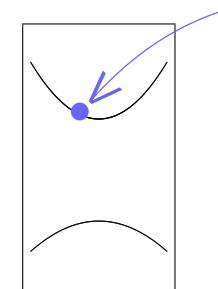


$|N - 1, 1\rangle$
 $|N - 1, 0\rangle$

N Electrons



$N + 1$ Electrons



G_1 = Single-particle Green function

G_2 = Two-particle Green function

1. Ground state $|N, 0\rangle$:

Density-functional theory \Rightarrow Geometry

Many-body perturbation theory:

2. $|N, 0\rangle \rightarrow |N \pm 1, m\rangle$:

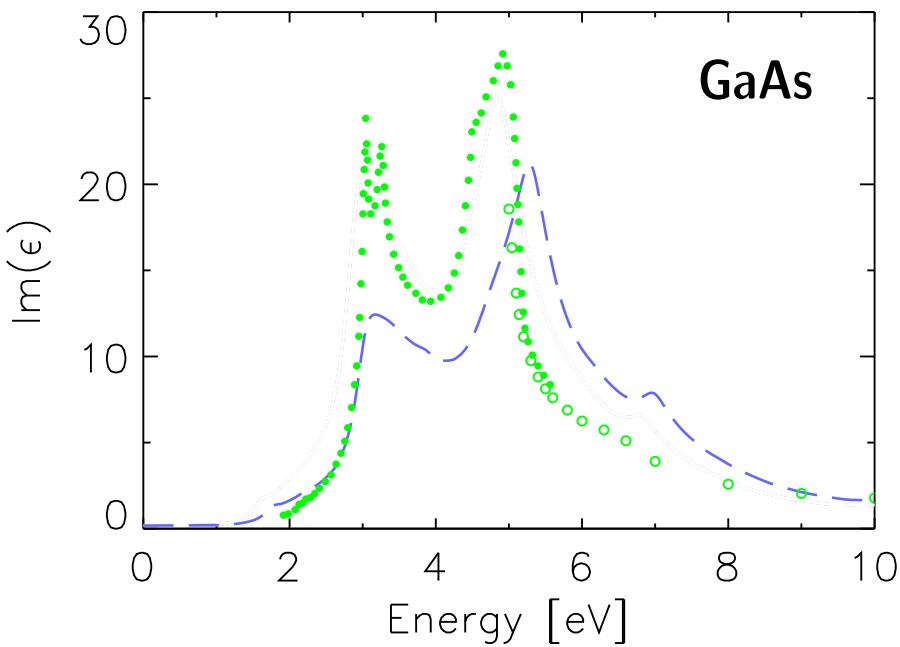
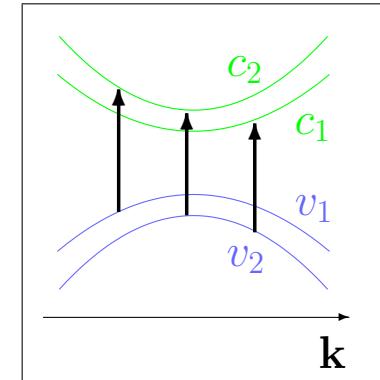
Dyson equation for G_1 \Rightarrow Band structure

3. $|N, 0\rangle \rightarrow |N, m\rangle$:

G_1 + Bethe-Salpeter Eq. for G_2 \Rightarrow El.-hole pairs

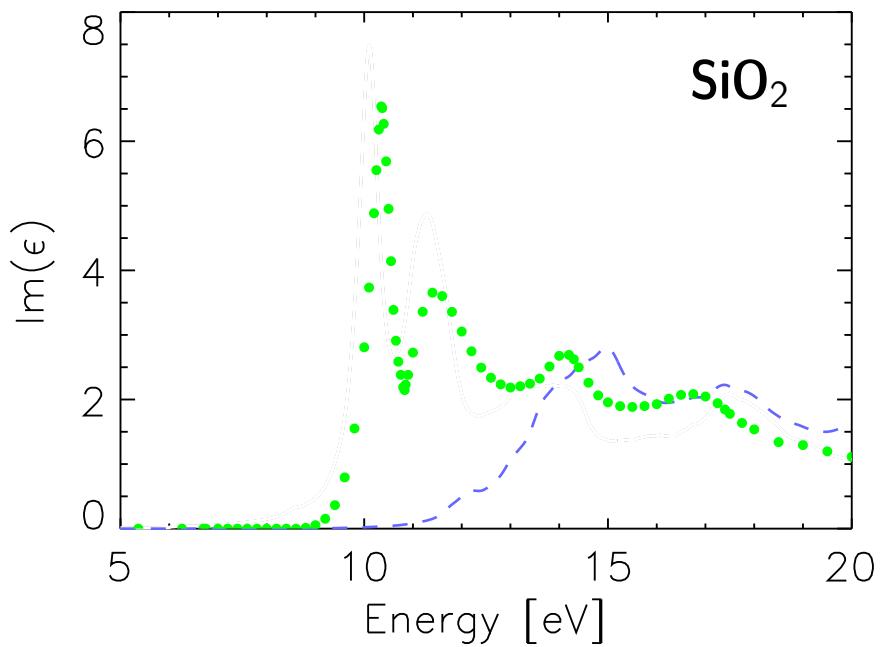
Bulk crystals: Optical absorption spectrum

- - - Free interband transitions: $\epsilon_2(\omega) \sim \sum_{vck} |M_{vck}|^2 \delta(\omega - (\epsilon_{c\mathbf{k}}^{\text{QP}} - \epsilon_{v\mathbf{k}}^{\text{QP}}))$



● ○ ○ Exp.:

D.E. Aspnes and A.A. Sturge, PRB 27, 985 (1983);
 P. Lautenschlager et al., PRB 35, 9174 (1987);
 H.R. Philipp, Sol.St.Comm. 4, 73 (1966).



[M. Rohlffing and S.G. Louie, PRL 81, 2312 (1998);
 E. Chang, M.R. and S.G. Louie, PRL 85, 2613 (2000).]

The two-particle Green function G_2

[G. Strinati, PRB 29, 5718 (1984); Rivista del Nuovo Cimento 11, 1 (1988).]

$$G_2(12; 1'2') = -\langle N, 0 | T(\hat{\psi}(1)\hat{\psi}(2)\hat{\psi}^\dagger(2')\hat{\psi}^\dagger(1')) | N, 0 \rangle$$

Two-particle correlation function: $L(12, 1'2') := -G_2(12, 1'2') + G_1(11') \cdot G_1(22')$

Equation of motion (Bethe-Salpeter equation: BSE):

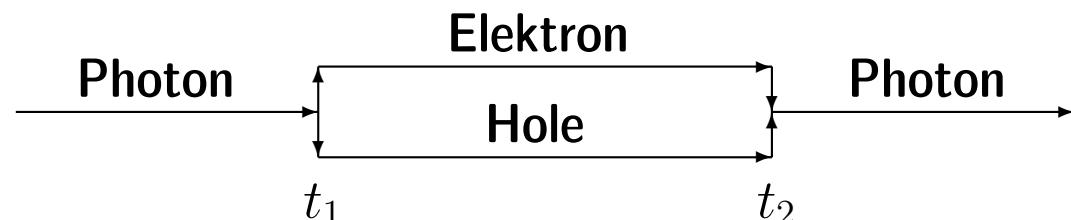
$$L(12; 1'2') = L_0(12; 1'2') + \int d(3456) L_0(14; 1'3) K(35; 46) L(62; 52')$$

$L_0(12; 1'2')$ $\hat{=}$ independent particles $K(35; 46)$: Interaction kernel \implies Coupling

Depending on the order of the times t_1, t_2, t'_1, t'_2 :

Propagation of two electrons / two holes / electron+hole

Here: $t_1=t'_1$ and $t_2=t'_2$ \implies Simultaneous creation of an electron/hole PAIR:



Explicit formulation of the Bethe-Salpeter equation

- Calculate G_1 from the (GW) QP states

$$\Rightarrow L_0(12, 1'2'; \omega) = i \sum_{v,c} \left[\frac{\psi_c(\mathbf{x}_1)\psi_v^*(\mathbf{x}'_1)\psi_v(\mathbf{x}_2)\psi_c^*(\mathbf{x}'_2)}{\omega - (\epsilon_c - \epsilon_v)} - \frac{\psi_v(\mathbf{x}_1)\psi_c^*(\mathbf{x}'_1)\psi_c(\mathbf{x}_2)\psi_v^*(\mathbf{x}'_2)}{\omega + (\epsilon_c - \epsilon_v)} \right]$$

v = Sum over occupied states c = Sum over empty states

- Similar expansion: $L(12, 1'2'; \omega) = i \sum_S \left[\frac{\chi_S(\mathbf{x}_1, \mathbf{x}'_1)\chi_S^*(\mathbf{x}'_2, \mathbf{x}_2)}{\omega - \Omega_S} - \frac{\chi_S(\mathbf{x}_2, \mathbf{x}'_2)\chi_S^*(\mathbf{x}'_1, \mathbf{x}_1)}{\omega + \Omega_S} \right]$
(→ has a structure similar to the QP expansion of G_1)

- Particle-hole amplitudes $\chi_S(\mathbf{x}, \mathbf{x}') = -\langle N, 0 | \psi^\dagger(\mathbf{x}') \psi(\mathbf{x}) | N, S \rangle$
 $= \sum_v^{\text{occ}} \sum_c^{\text{empty}} A_{vc}^S \psi_c(\mathbf{x}) \psi_v^*(\mathbf{x}') + B_{vc}^S \psi_v(\mathbf{x}) \psi_c^*(\mathbf{x}')$

- Expansion of χ_S in the QP states (v, c) $\hat{=}$ "Second Quantization":

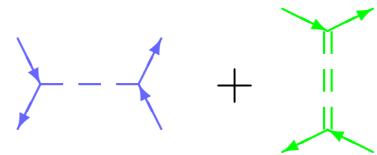
Double transformation from the continuous 3D space into a discrete basis

The electron-hole interaction K^{eh}

$$K^{eh}(12,34) = \frac{\delta[V_{\text{Coul}}(1)\delta(13) + \Sigma(13)]}{\delta G_1(42)} \quad [1 \hat{=} (\mathbf{r}_1, \sigma_1, t_1) \text{ etc.}]$$

Assume: $\Sigma = iG_1W$ and $G_1 \frac{\delta W}{\delta G_1} \approx 0$

$$\implies K^{eh}(12,34) = -i\delta(13)\delta(2^-4)v(14) + i\delta(14)\delta(23)W(1^+3)$$



$$\langle vck | K^{eh} | v'c'k' \rangle = \iint dxdx' \psi_{ck+Q}^*(x) \psi_{v k}(x) \mathbf{v}(\mathbf{r}, \mathbf{r}') \psi_{c'k'+Q}(x') \psi_{v'k'}^*(x')$$

$$-\frac{i}{2\pi} \int d\omega e^{-i\omega 0^+} \iint dxdx' \psi_{ck+Q}^*(x) \psi_{c'k'+Q}(x) \mathbf{W}(\mathbf{r}, \mathbf{r}', \omega) \psi_{v k}(x') \psi_{v'k'}^*(x') \\ \times \left[\frac{1}{\Omega_S - \omega - (\varepsilon_{c'k'+Q}^{\text{QP}} - \varepsilon_{vk}^{\text{QP}}) + i0^+} + \frac{1}{\Omega_S + \omega - (\varepsilon_{ck+Q}^{\text{QP}} - \varepsilon_{v'k'}^{\text{QP}}) + i0^+} \right]$$

$$= \text{Repulsive exchange term } (=:K^{eh,x}) + \text{Attractive direct term } (=:K^{eh,d})$$

Bethe-Salpeter equation: (standard) eigenvalue problem

- Equation of motion for $L \implies$ Eigenvalue problem

$$(E_c - E_v) A_{vc}^S + \sum_{v'c'} K_{vc,v'c'}^{AA}(\Omega_S) A_{v'c'}^S + \sum_{v'c'} K_{vc,v'c'}^{AB}(\Omega_S) B_{v'c'}^S = \Omega_S A_{vc}^S$$

$$\sum_{v'c'} K_{vc,v'c'}^{BA}(\Omega_S) A_{v'c'}^S + (E_c - E_v) B_{vc}^S + \sum_{v'c'} K_{vc,v'c'}^{BB}(\Omega_S) B_{v'c'}^S = -\Omega_S B_{vc}^S$$

- The electron-hole interaction matrix elements are given by

$$K_{vc,v'c'}^{AA}(\Omega_S) = i \int d(3456) \psi_v(\mathbf{x}_4) \psi_c^*(\mathbf{x}_3) K(35, 46; \Omega_S) \psi_{v'}^*(\mathbf{x}_5) \psi_{c'}(\mathbf{x}_6)$$

$$K_{vc,v'c'}^{AB}(\Omega_S) = i \int d(3456) \psi_v(\mathbf{x}_4) \psi_c^*(\mathbf{x}_3) K(35, 46; \Omega_S) \psi_{v'}^*(\mathbf{x}_6) \psi_{c'}(\mathbf{x}_5)$$

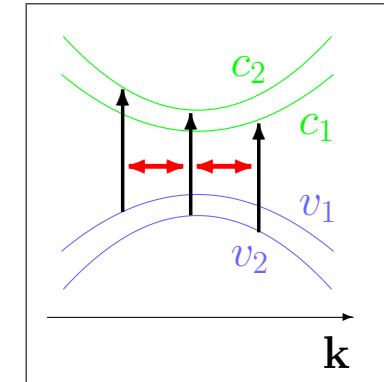
...

- "Tamm-Dankoff approx.": The off-diagonal blocks K^{AB} and K^{BA} are small \implies Neglect them:

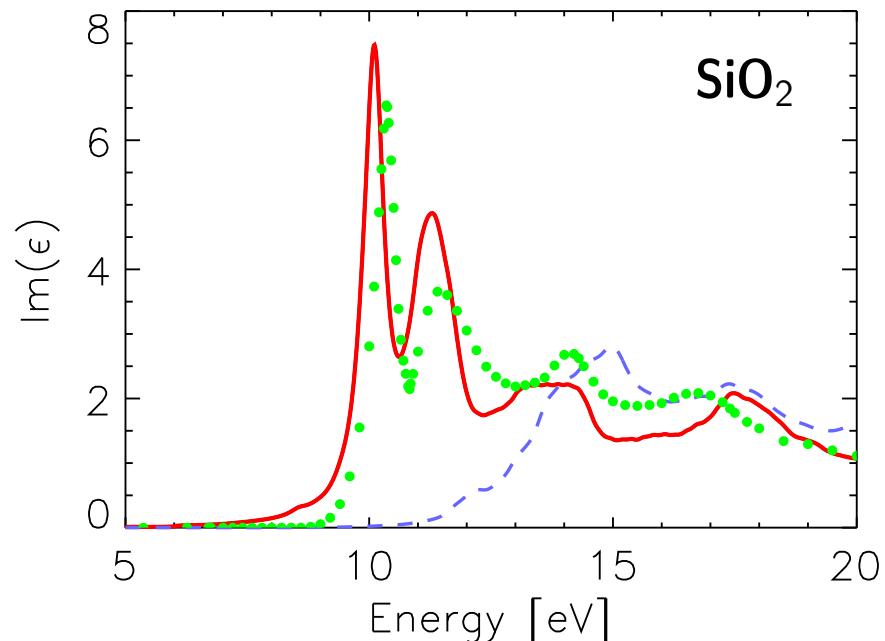
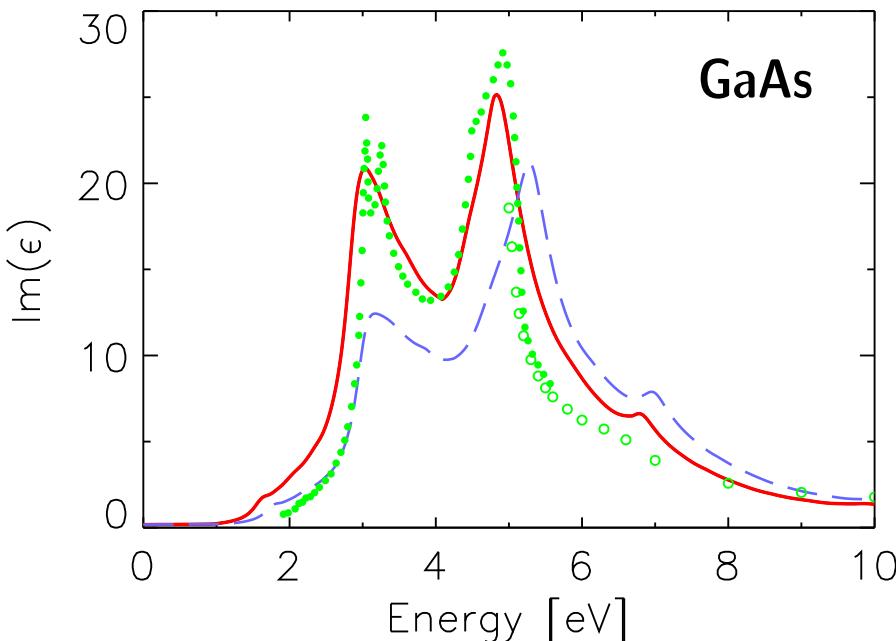
$$(E_c - E_v) A_{vc}^S + \sum_{v'c'} K_{vc,v'c'}^{AA}(\Omega_S) A_{v'c'}^S = \Omega_S A_{vc}^S , \quad \text{i.e.} \quad |N, S\rangle = \sum_{v,c} A_{vc}^S \hat{a}_v^\dagger \hat{b}_c^\dagger |N, 0\rangle$$

Bulk crystals: Optical absorption spectrum

- - - Free interband transitions: $\epsilon_2(\omega) \sim \sum_{vck} |M_{vck}|^2 \delta(\omega - (\epsilon_{ck}^{QP} - \epsilon_{vk}^{QP}))$
- With electron-hole interaction: coupled excitations $|S\rangle$



$$\epsilon_2(\omega) \sim \sum_S |M_S|^2 \delta(\omega - \Omega_S)$$



● ○ ○ Exp.:

D.E. Aspnes and A.A. Sturge, PRB 27, 985 (1983);
 P. Lautenschlager et al., PRB 35, 9174 (1987);
 H.R. Philipp, Sol.St.Comm. 4, 73 (1966).

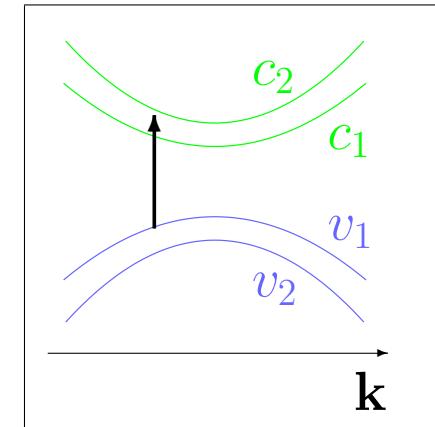
[M. Röhlifing and S.G. Louie, PRL 81, 2312 (1998);
 E. Chang, M.R. and S.G. Louie, PRL 85, 2613 (2000).]

Effect of the electron-hole interaction

- Without interaction:

$$(\varepsilon_{c\mathbf{k}+\mathbf{Q}}^{\text{QP}} - \varepsilon_{v\mathbf{k}}^{\text{QP}}) A_{vc\mathbf{k}}^S = \Omega_S A_{vc\mathbf{k}}^S$$

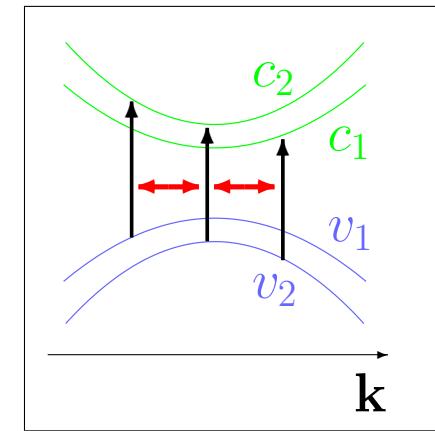
\Rightarrow Free interband transitions; $\Omega_S = (\varepsilon_{c\mathbf{k}+\mathbf{Q}}^{\text{QP}} - \varepsilon_{v\mathbf{k}}^{\text{QP}})$



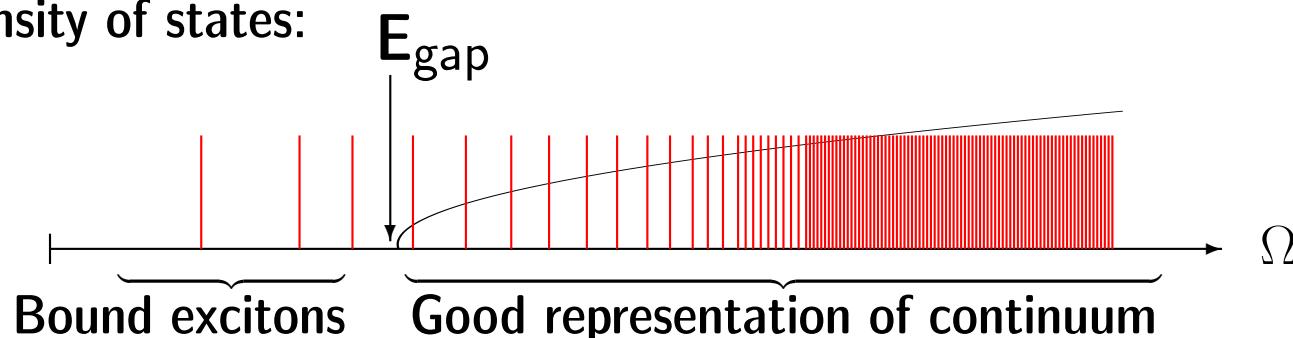
- With interaction:

$$(\varepsilon_{c\mathbf{k}+\mathbf{Q}}^{\text{QP}} - \varepsilon_{v\mathbf{k}}^{\text{QP}}) A_{vc\mathbf{k}}^S + \sum_{v'c'\mathbf{k}'} \langle v\mathbf{c}\mathbf{k} | K^{eh} | v'\mathbf{c}'\mathbf{k}' \rangle A_{v'\mathbf{c}'\mathbf{k}'}^S = \Omega_S A_{vc\mathbf{k}}^S$$

All electron-hole excitation energies



\Rightarrow Density of states:



Effect of the electron-hole interaction on the optical spectrum:

- Free interband transitions / no interaction:

$$\epsilon_2(\omega) = \frac{4\pi e^2}{\omega^2} \sum_{vck} |M_{vck}|^2 \delta(\omega - (\epsilon_{ck}^{\text{QP}} - \epsilon_{vk}^{\text{QP}}))$$

$$M_{vck} = \vec{\lambda} \cdot \langle v\mathbf{k} | \vec{V} | c\mathbf{k} \rangle$$

$\vec{\lambda}$ Polarisation of the light

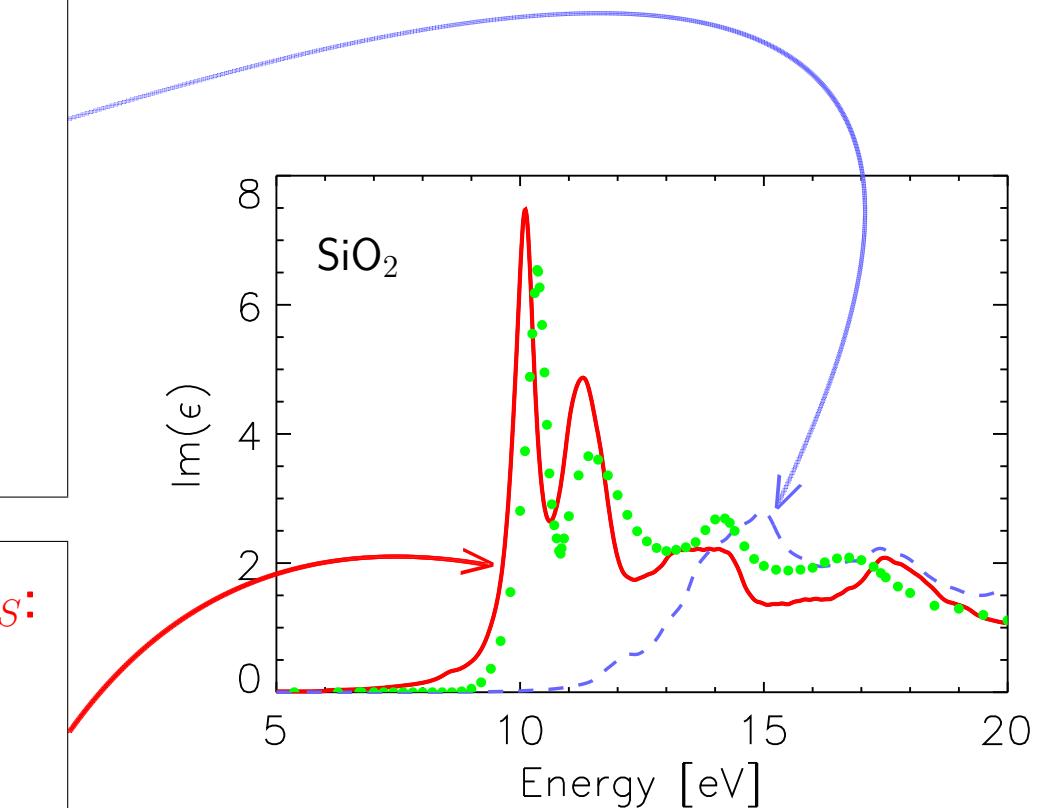
\vec{V} Velocity operator

- Coupled excitations / $(\Delta\varepsilon + K^{eh})A_S = \Omega_S A_S$:

$$\epsilon_2(\omega) = \frac{4\pi e^2}{\omega^2} \sum_S |M_S|^2 \delta(\omega - \Omega_S)$$

$$M_S = \vec{\lambda} \cdot \langle 0 | \vec{V} | S \rangle = \sum_{vck} A_{vck}^S M_{vck}$$

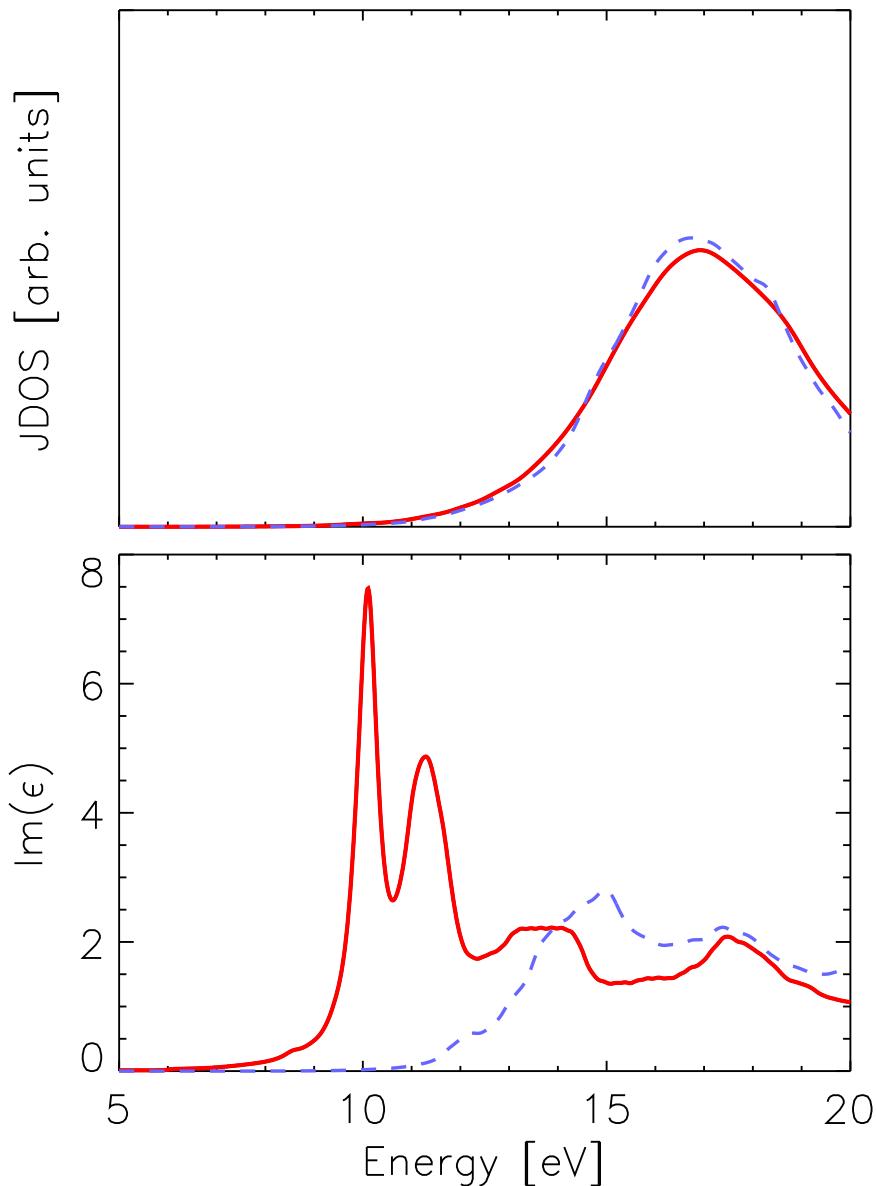
Coherent superposition $M_{vck} \rightarrow M_S$;
phase sensitive → constructive / destructive!



• ○ ○ Exp.:

H.R. Philipp, Sol.St.Comm. 4, 73 (1966).

SiO_2 : $\epsilon_2(\omega) \longleftrightarrow \text{JDOS}$



- - - Without interaction:
Free interband transitions $|vck\rangle$

— With interaction:
Coupled Excitations $|S\rangle$

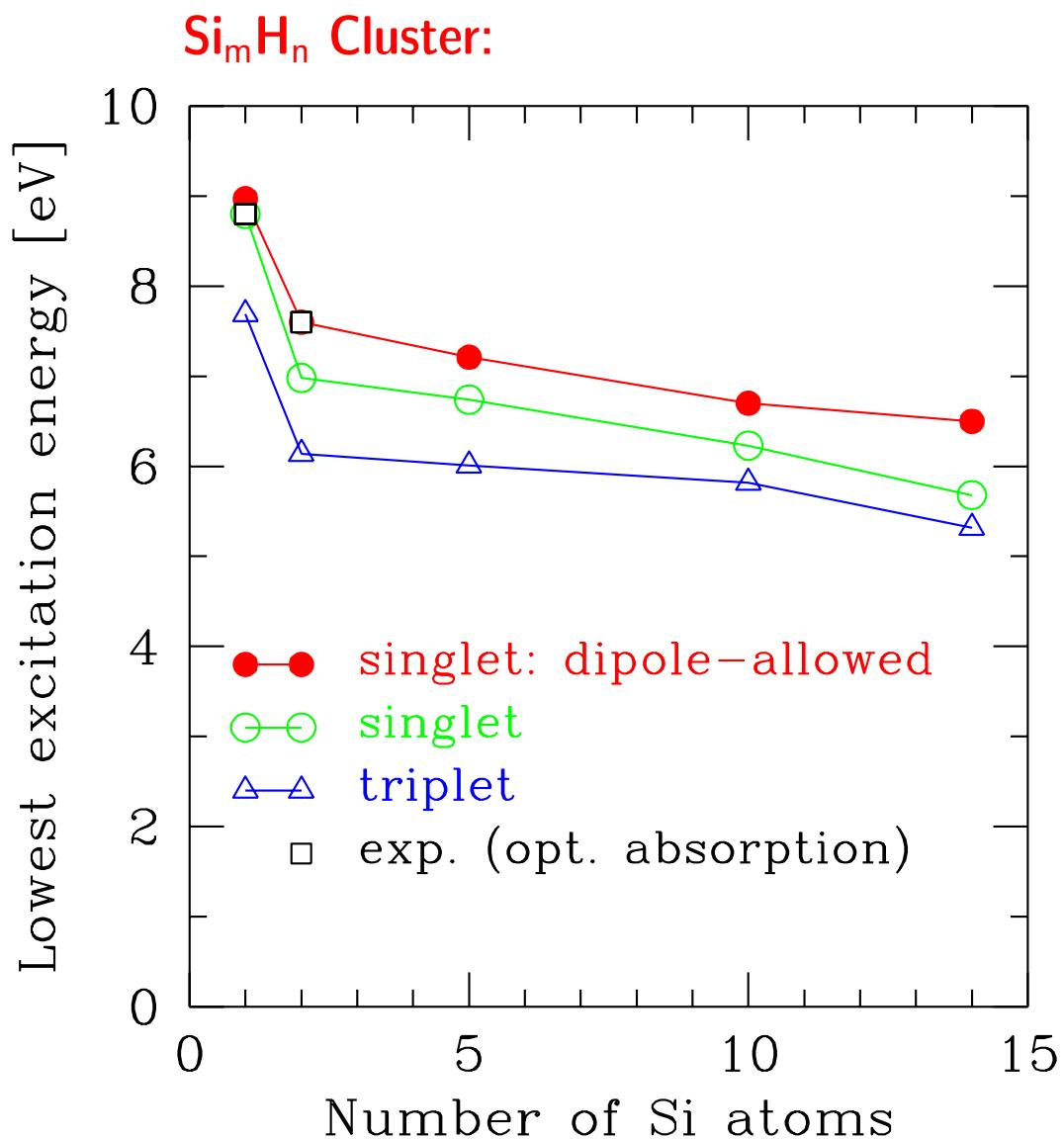
- The density of states is only weakly modified by K^{eh} .
- The modifications in the SPECTRUM result from INTERFERENCE:

$$M_S = \sum_{vck} A_{vck}^S M_{vck}$$

Excitation energies of atoms and molecules

	Spin singlet		Spin triplet	
[eV]	BSE	Exp. ^a	BSE	Exp. ^a
He	20.75	20.615	19.81	19.818
Ne	16.95	16.848	16.71	16.668
Ar	11.99	11.827	11.76	11.631
CO	7.9	8.03	5.5	6.01
HCl	7.5	8.1	6.9	
CH ₄	8.6	8.52	8.2	
C ₂ H ₄	7.0	7.11	3.7	4.36

^aExp.: Landolt-Börnstein, Vol. I-1 (1950);
 G. Herzberg, *Molecular Spectra* (1966);
 L. Serrano-Andres *et al.*, JCP 98, 3151 (1993).



Exp.: U. Itoh *et al.*, J.Chem.Phys. 85, 4867 (1986).

LiF: Bulk Band Structure

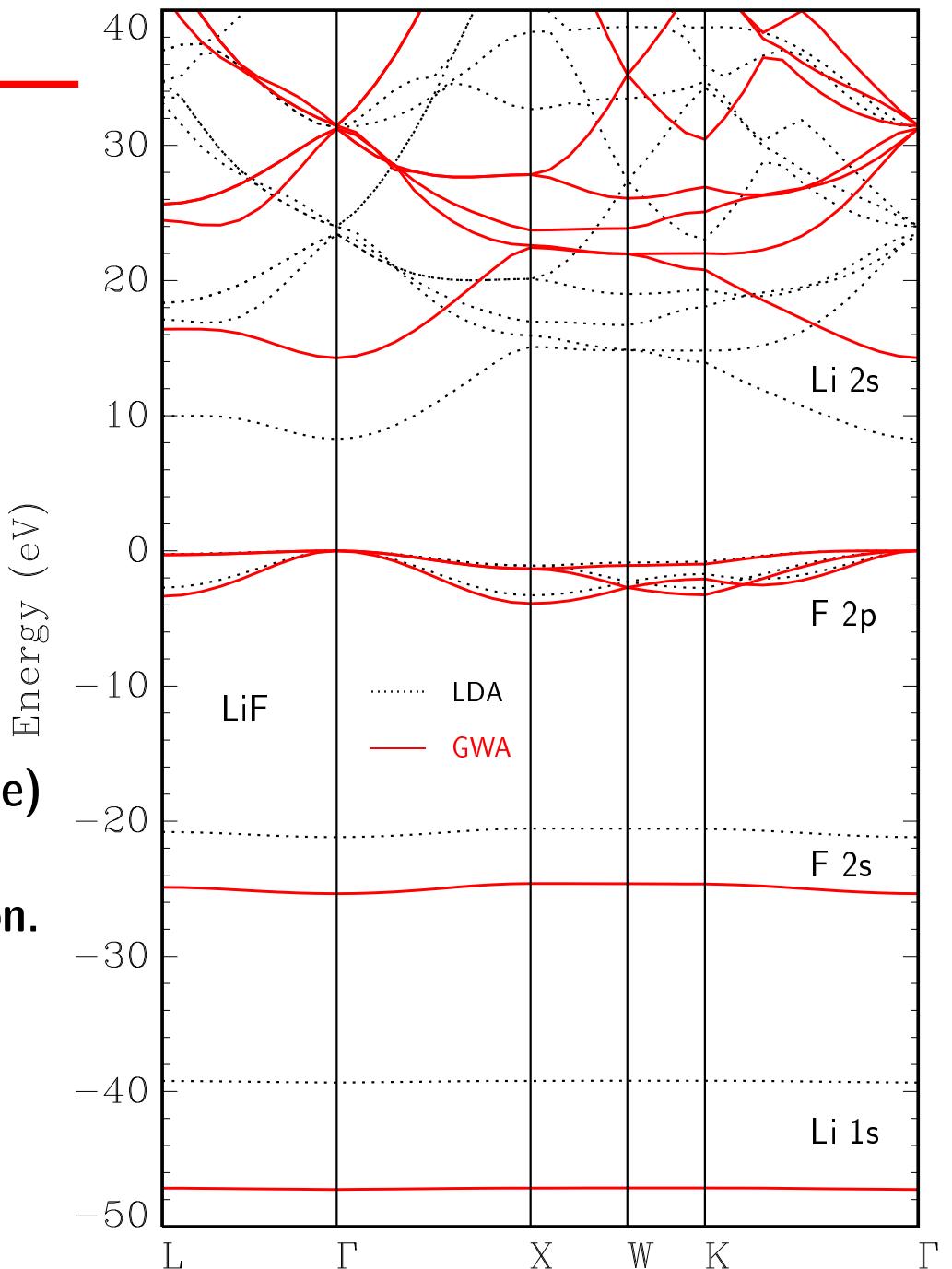
- Dielectric constant: $\epsilon_{\infty}^{\text{RPA}} = 1.8$

- Some band-structure energies:

[eV]	LDA	Δ^{QP}	QP	Exp. ¹
Li 1s	-39.3	-7.9	-47.2	-49.8
F 2s	-20.8	-4.2	-25.0	-23.9
F 2p	-2.7..0	-0.6..0	-3.3..0	
"Li 2s"	8.3...	+6.0...	14.3...	14.4...
E _{gap}	8.3	+6.0	14.3	14.4...

- A dielectric model function (Hybertsen/Louie) yields the same band structure near the gap
 ⇒ We use this model function from now on.

Exp.: L.I. Johansson and S.B.M. Hagström,
 Physica Scripta 14, 55 (1976);
 D.M. Roessler and W.C. Walker,
 J.Opt.Soc.Am. 57, 835 (1967).



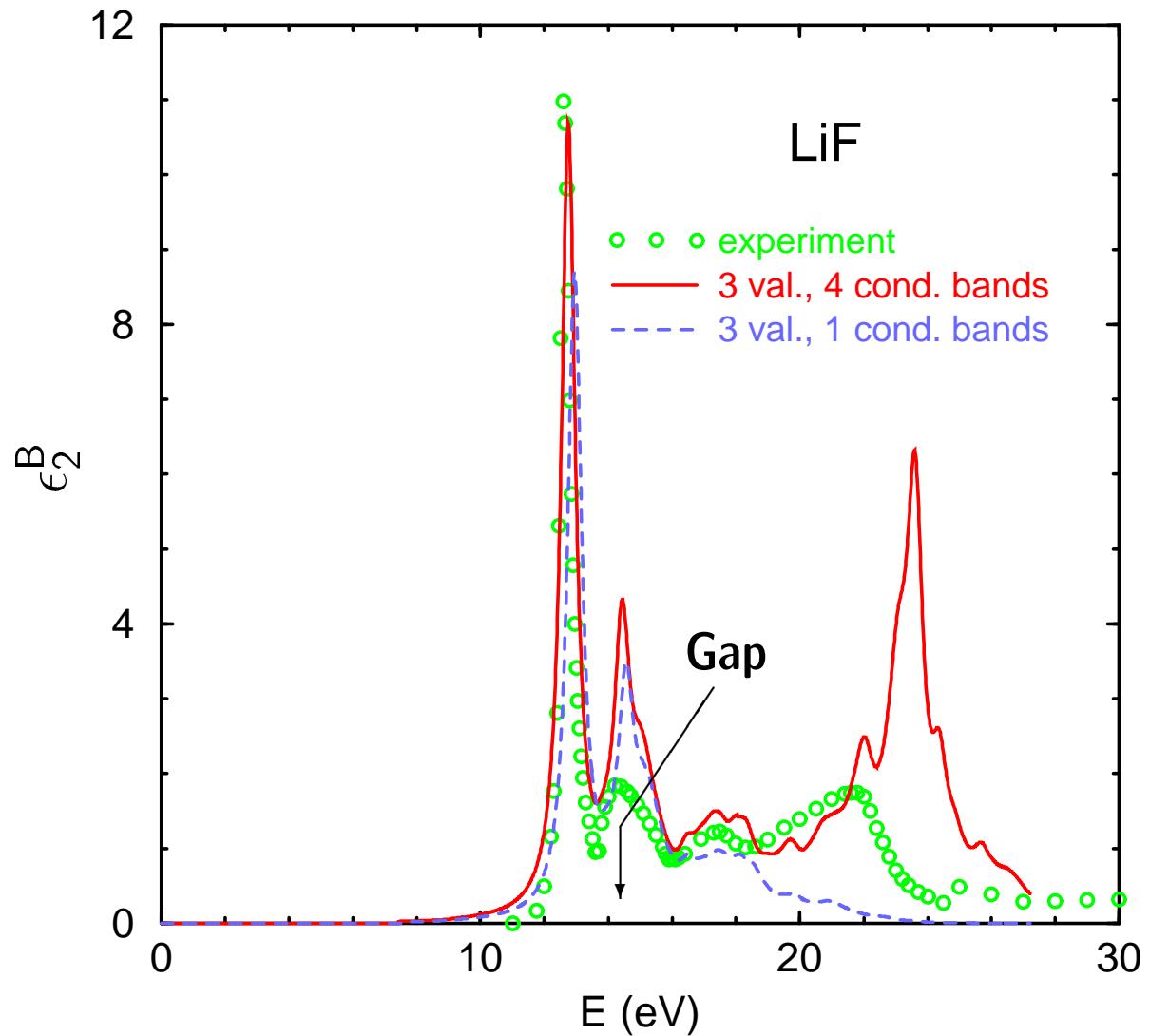
LiF: Bulk Absorption Spectrum

- Exciton at 12.8 eV
(1.5 eV below $E_{\text{gap}}=14.3$ eV)

- Matrix size of $H^{e,h}$:

	$\text{F } 2p \rightarrow$ "Li 2s+p"	$\text{F } 2p \rightarrow$ "Li 2s"
Bulk (256 k)	$3 \times 4 \times 256 = 3072$	$3 \times 1 \times 256 = 768$
6-L Slab (64 k_{\parallel})	$18 \times 24 \times 64 = 27648$	$18 \times 6 \times 64 = 6912$

- $\text{F } 2p \rightarrow$ "Li 2s":
Good representation of the exciton
AND convenient matrix size.



LiF Exciton: Electron-Hole Correlation

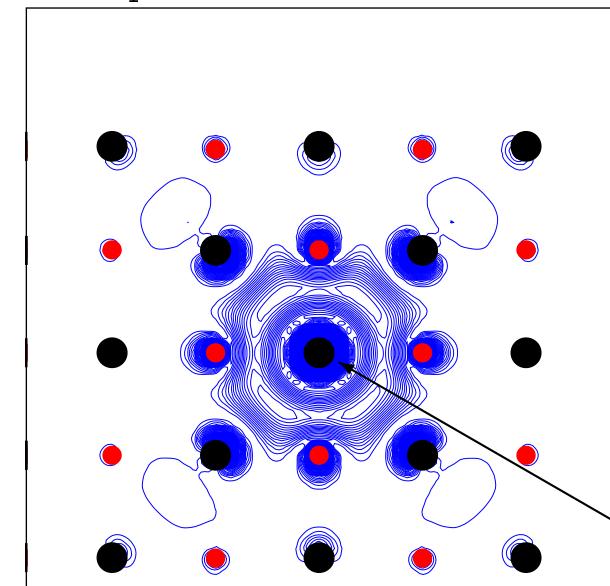
Exciton wave function (2-particle correlation function):

$$\chi_S(\mathbf{r}_h, \mathbf{r}_e) = \sum_{\mathbf{k}} \sum_v^{\text{hole}} \sum_c^{\text{elec}} A_{vck}^S \psi_{v\mathbf{k}}^*(\mathbf{r}_h) \psi_{c\mathbf{k}+\mathbf{Q}}(\mathbf{r}_e)$$

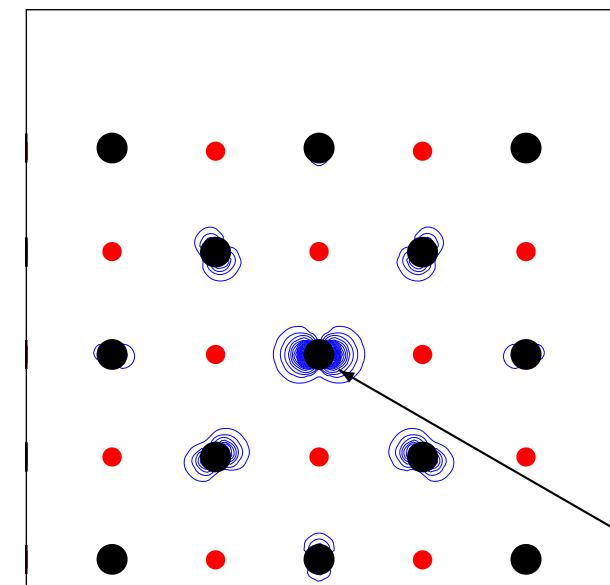
\mathbf{r}_h fixed $\Rightarrow \chi_S(., \mathbf{r}_e) = \text{Distribution of the electron}$

\mathbf{r}_e fixed $\Rightarrow \chi_S(\mathbf{r}_h, .) = \text{Distribution of the hole}$

B₁: electron distribution



B₁: hole distribution



Li
F

Hole

Electron

Si(111)-(2×1): Surface Exciton

Experiment (Internal Reflection; DRS):

G. Chiarotti,, PRB 4, 3398 (1971);

P. Chiaradia et al., PRL 52, 1145 (1984);

F. Ciccaci et al., PRL 56, 2411 (1986).

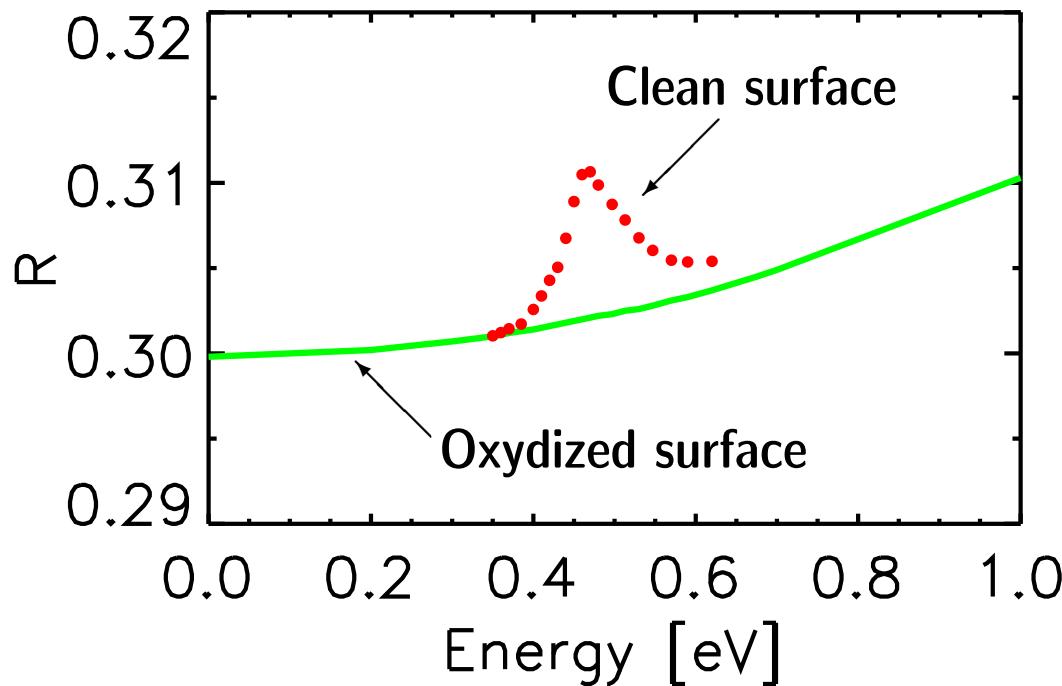
Theory: J.E. Northrup et al., PRL 66, 500 (1991);

L. Reining and R. Del Sole, PRL 67, 3816 (1991);

M. Rohlfing and S.G. Louie, PRL 83, 856 (1999).

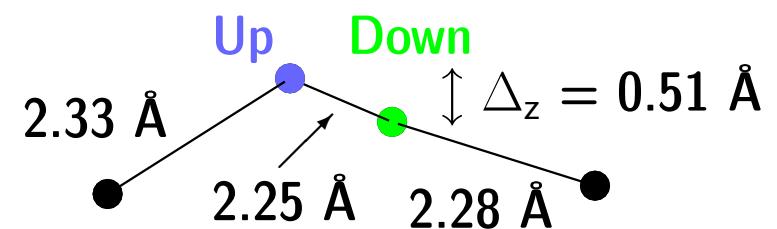
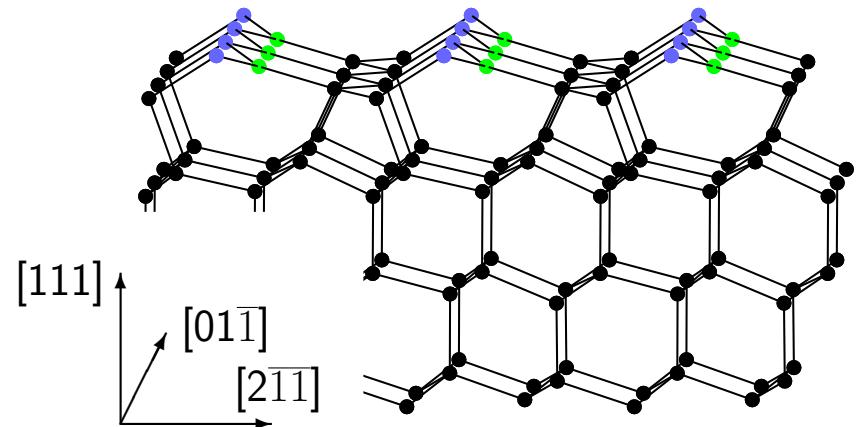
M. Rohlfing and J. Pollmann, PRL 88, 176801 (2002).

Reflectivity (exp.):

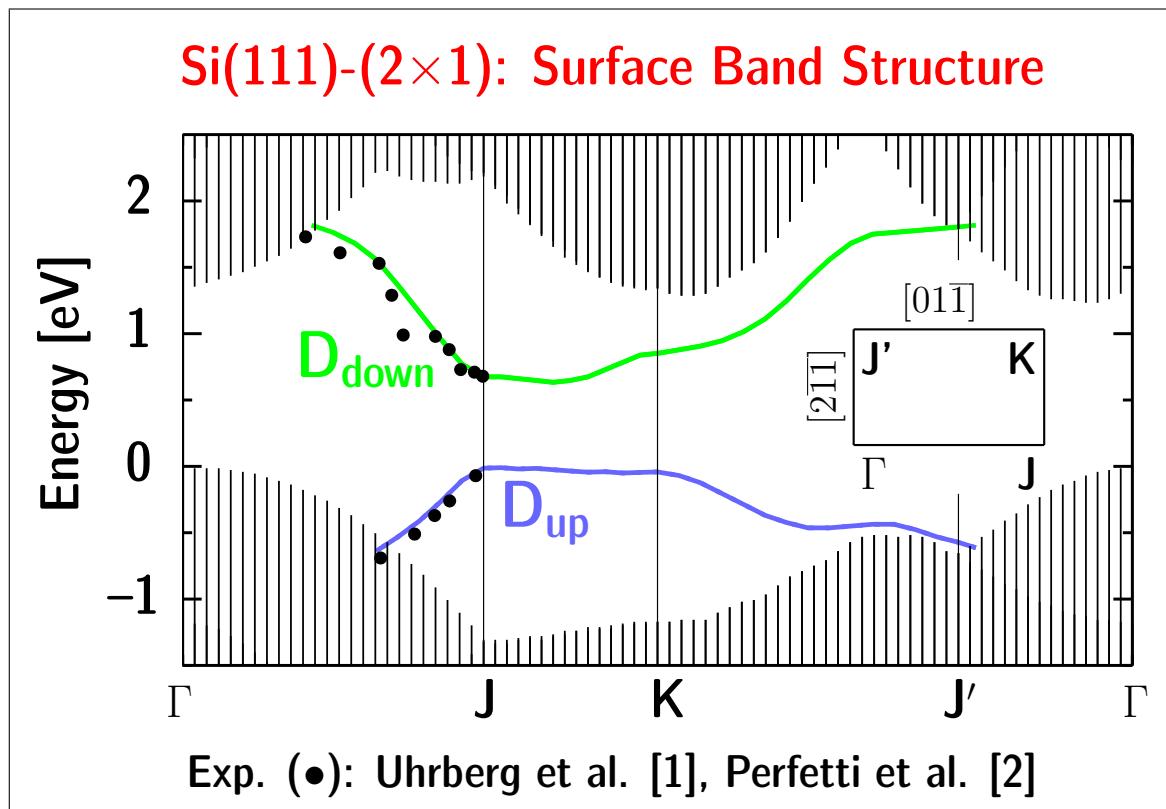


⇒ Surface Exciton at 0.45 eV

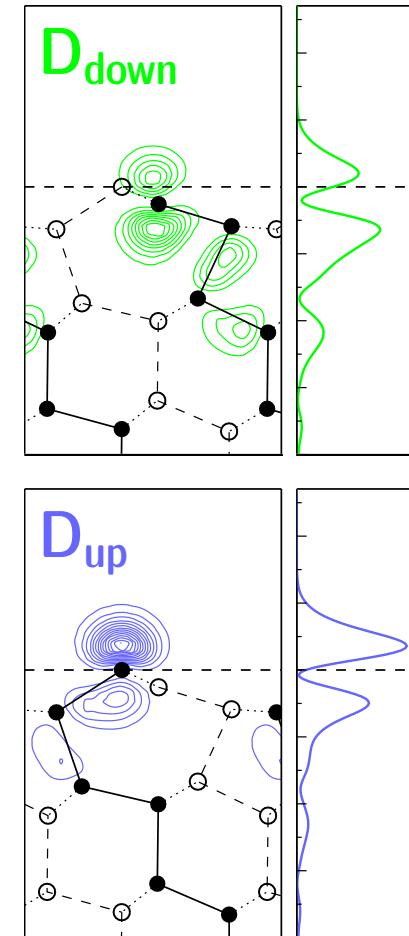
Structure of the surface: Pandey chains



Si(111)-(2×1): Surface Band Structure



Dangling-Bond States (at J):



Surface Gap [eV]:

This work	0.69	[GWA]
ARPES+ARIPES	0.75	[1, 2]
ARPES (<i>n</i> -Si)	0.5	[3]
STM	0.6	[4]

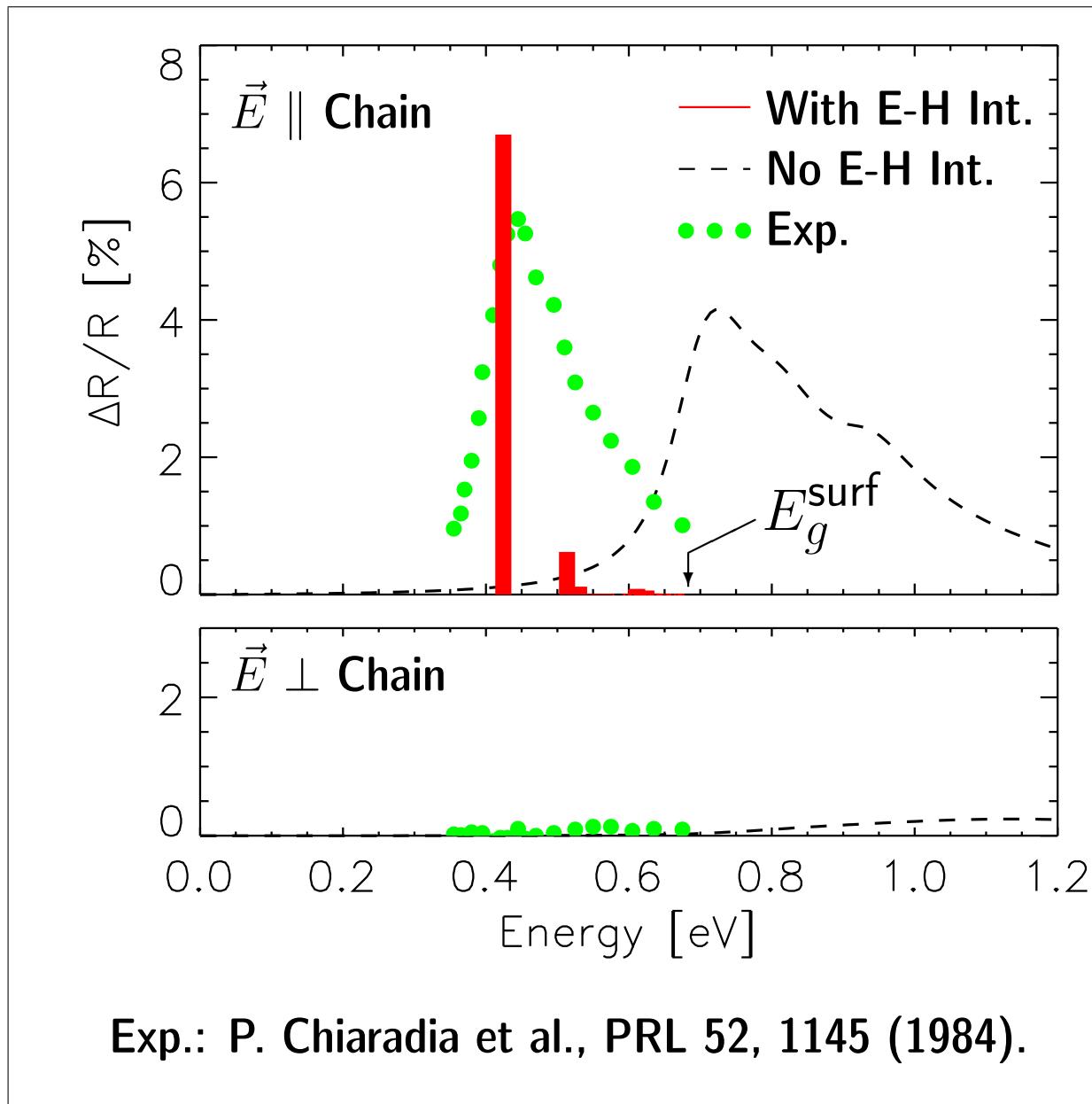
[1] R.I.G. Uhrberg et al.
[PRL 48, 1032 (1982)]

[2] P. Perfetti et al.
[PRB 36, 6160 (1987)];
A. Cricenti et al.
[PRB 41, 12908 (1990)]

[3] P. Martensson et al. [PRB 32, 6959 (1985)]

[4] R.M. Feenstra et al. [PRL 56, 608 (1986); PRL 57, 2579 (1986).]

Si(111)-(2×1): Differential Reflectivity Spectrum (DRS)

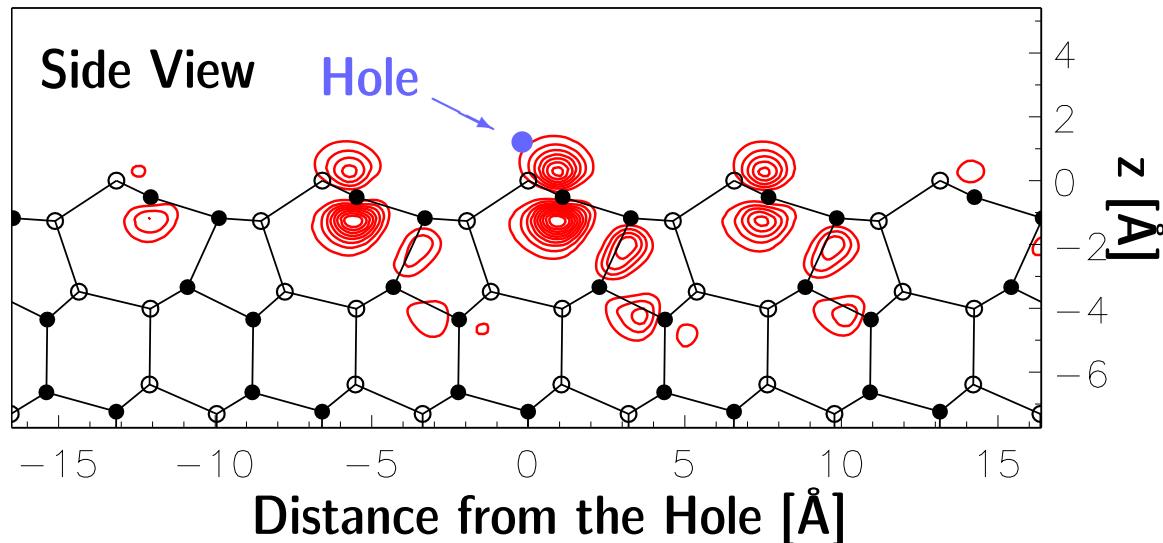


$$\frac{\Delta R}{R}(\omega) = 4 \frac{\omega \Im(d \cdot \epsilon^{\text{surf}}(\omega))}{c (\epsilon_b - 1)}$$

[S. Selci *et al.*, J.V.Sc.Tc. A 5, 327 (1987)]

- Discrete exciton spectrum.
- The optical spectrum is dominated by the lowest exciton at 0.43 eV.
- Binding energy: 0.26 eV ($\gg E_B^{\text{bulk}} = 0.015$ eV).

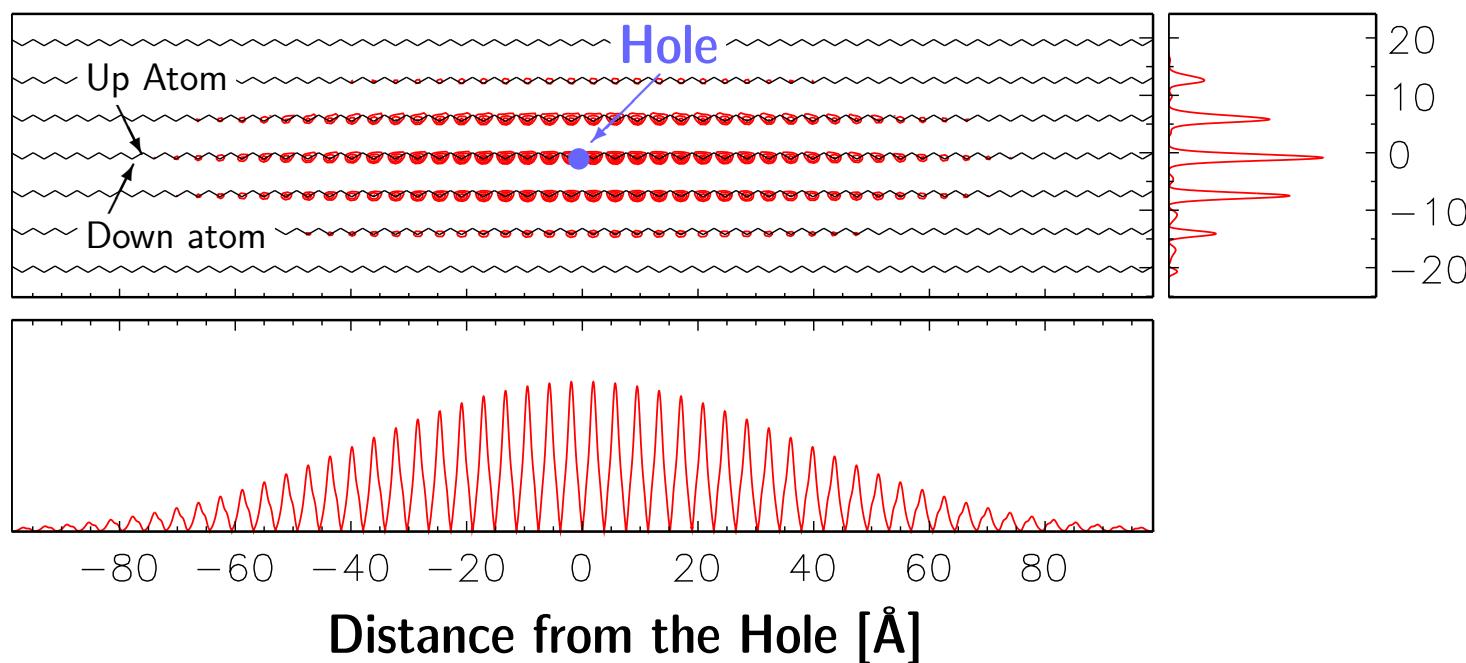
Exciton Wave Function: Electron Distribution Relative to the Hole



$$\chi^S(\mathbf{r}_h, \mathbf{r}_e) = \sum_{vck} A_{vck}^S \psi_{v\mathbf{k}}^*(\mathbf{r}_h) \psi_{c\mathbf{k}+\mathbf{Q}}(\mathbf{r}_e)$$

[\mathbf{r}_h = coordinates of Hole, \mathbf{r}_e = Electron]

Top View:



**Mean distance
Electron—Hole:**

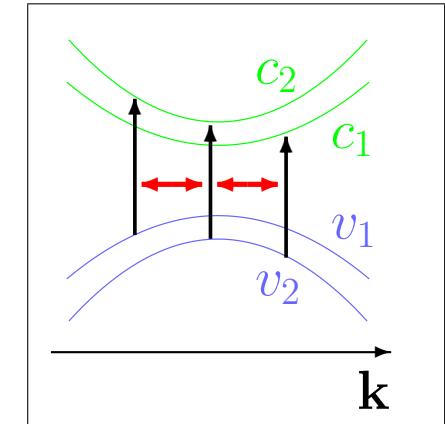
- along chain: **40 Å**
(~ Wannier Exciton)
- perpendicular: **8 Å**
(~ Frenkel Exciton)

Coupled electron-hole excitations

- Expansion of the excitations:

$$|S\rangle = \sum_v^{\text{hole}} \sum_c^{\text{elec}} \sum_{\mathbf{k}} A_{vck}^S |vck\rangle$$

$|vck\rangle := \hat{a}_{vk}^\dagger \hat{b}_{ck+Q}^\dagger |0\rangle$ free electron-hole interband transition



- Bethe-Salpeter equation for G_2 or, resp., $|S\rangle$:

$$(E_{ck+Q}^{\text{QP}} - E_{vk}^{\text{QP}}) A_{vck}^S + \sum_{v'c'\mathbf{k}'} \langle vck | K^{eh} | v'c'\mathbf{k}' \rangle A_{v'c'\mathbf{k}'}^S = \Omega_S A_{vck}^S$$

E_{vk}^{QP} , E_{ck+Q}^{QP} QP energies of the single-particle states
 K^{eh} Electron-hole interaction
 Ω_S Excitation energy

\sim Single-excit. CI
 with W instead of v
 \implies Correlation
 in the interaction

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