

# Magnetic-field-induced phase transitions in a Si/SiGe hole system

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Seminal theoretical work on the integer quantum Hall effect describes a two-parameter scaling [1] in which both the dissipative conductance  $\sigma_{xx}$  and the Hall conductance  $\sigma_{xy}$  vary with the sample length  $L$ . The renormalisation-group function can be illustrated by flow lines. After successive length-scale transformation the flow lines are directed towards fixed points  $(\sigma_{xx}, \sigma_{xy}) = (0, n)$  (in units of  $e^2/h$ ) in which  $n = 0, 1, 2, \dots$  is the Landau level index. In addition to these “localization fixed points” which describe the localisation of the electron wavefunctions near the Fermi energy, there are “intermediate-coupling fixed points” on  $\sigma_{xy} = n + 1/2$  which are related to transitions between quantum Hall states. Note that this theory was developed for zero temperature  $T = 0$  and the length scale transformations are accomplished by varying the sample length  $L$ . It has been shown that, in practice, the effective sample size can be varied by changing the temperature [2]. The studies of the “temperature-driven flow lines” in GaAs systems and more recently in Ge/SiGe systems [3] support the two-parameter scaling theory of Pruisken [1]

Recently there has been a great deal of interest in Si/SiGe hole systems. For example, at a Landau level filling factor  $\nu = 1.5$ , an insulating phase observed in such a system is not fully understood at present. It is called a “Hall insulator” since although  $\rho_{xx}$  approaches infinity, the Hall resistivity remains finite at approximately  $h/(2e^2)$ . This magneto-driven transition within the global phase scheme of Kivelson, Lee and Zhang. Moreover, in a Si/SiGe hole gas, the observed quantum Hall states are off odd filling factors at low magnetic fields, indicating that spin-splitting is comparable with the spacing between adjacent Landau levels. This provides an interesting system for the study of magnetic-field-induced transitions in the quantum Hall effect. In this paper, we present low-temperature transport measurements on a Si/SiGe hole system. In particular we report a study of the temperature-driven flow lines in the vicinity of the “Hall insulator” regime. We also observe magnetic-field-induced phase transitions which correspond to the temperature-independent points in  $\sigma_{xx}$  as shown in figure 1.

The temperature-driven flow lines are complicated thus we divided our results into four parts as shown in Fig. 2(a)-(d). It is evident that  $(\sigma_{xx}, \sigma_{xy}) = (1, 0)$  for  $\nu \rightarrow 1$ , as illustrated by the  $T$ -driven flow lines as shown in Fig. 2(a). Figure 2 (b) and 2 (c) show the  $T$ -driven flow diagrams near the two critical points  $B_{B1} = 8.1$  T and  $B_{B2} = 9.8$  T, respectively. It is evident that  $(\sigma_{xx}, \sigma_{xy}) = (0, 0)$  as  $T \rightarrow 0$  for the “Hall insulator” since  $\rho_{xx}$  approaches infinity whereas  $\rho_{xy}$  remains finite. We can clearly see that the directions of the flow lines in the “Hall insulator” are different from those near a quantum Hall state.  $B_{B2}$ , labelled as full squares, is, in fact, the boundary of the “Hall insulator” as well as the intermediate-coupling fixed point. It is worth mentioning that near  $B_{B1}$ ,  $\rho_{xy}$  slowly varies at different temperatures. Thus, although  $B_{B1}$  corresponds to the boundary of the “Hall insulator” regime, the intermediate-coupling fixed point is estimated to be at around  $B = 8.2$  T,  $\sim 0.1$  T higher than  $B_{B1}$ . To the right of this fixed point, the flow lines all show a “kink”. In contrast, to the left of the fixed point, the lines flow directly  $(0, 0)$  as  $T \rightarrow 0$ . The main finding is that, in contrast to a *single* intermediate-coupling fixed point between neighboring quantum Hall states previously predicted and observed, there are *two* intermediate-coupling points in a SiGe hole system. The reason for this is the existence of the “anomalous Hall insulator” near  $\nu = 1.5$ . Figure 2(d) shows the  $T$ -driven flow lines for  $2 < \nu < 3$ . At around  $A_2$ , the directions of the  $T$ -driven flow lines seem ambiguous. At present, the physical origin of this effect is not fully understood and awaits further experimental and theoretical investigations. Our new experimental results on  $T$ -driven flow lines, together with the pioneering work on the “Hall insulator” urge further investigations in order to understand the physical origin of this “anomalous insulating state”.

References:

- [1] A. M. M. Pruisken, Phys. Rev. B 32, 2636 (1985).
- [2] H. P. Wei *et al.*, Phys. Rev. B 33, 1488 (1986).
- [3] M. Hilke *et al.*, Nature 400, 735 (1999).

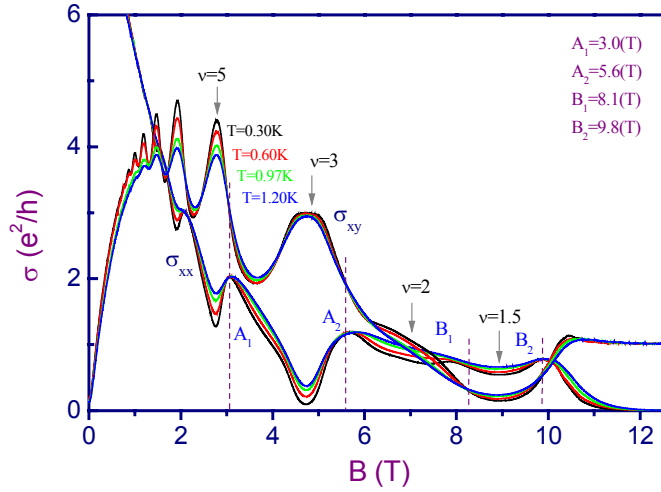


Figure 1. Conductivity  $\sigma_{xx}$  and  $\sigma_{xy}$  at various temperatures.

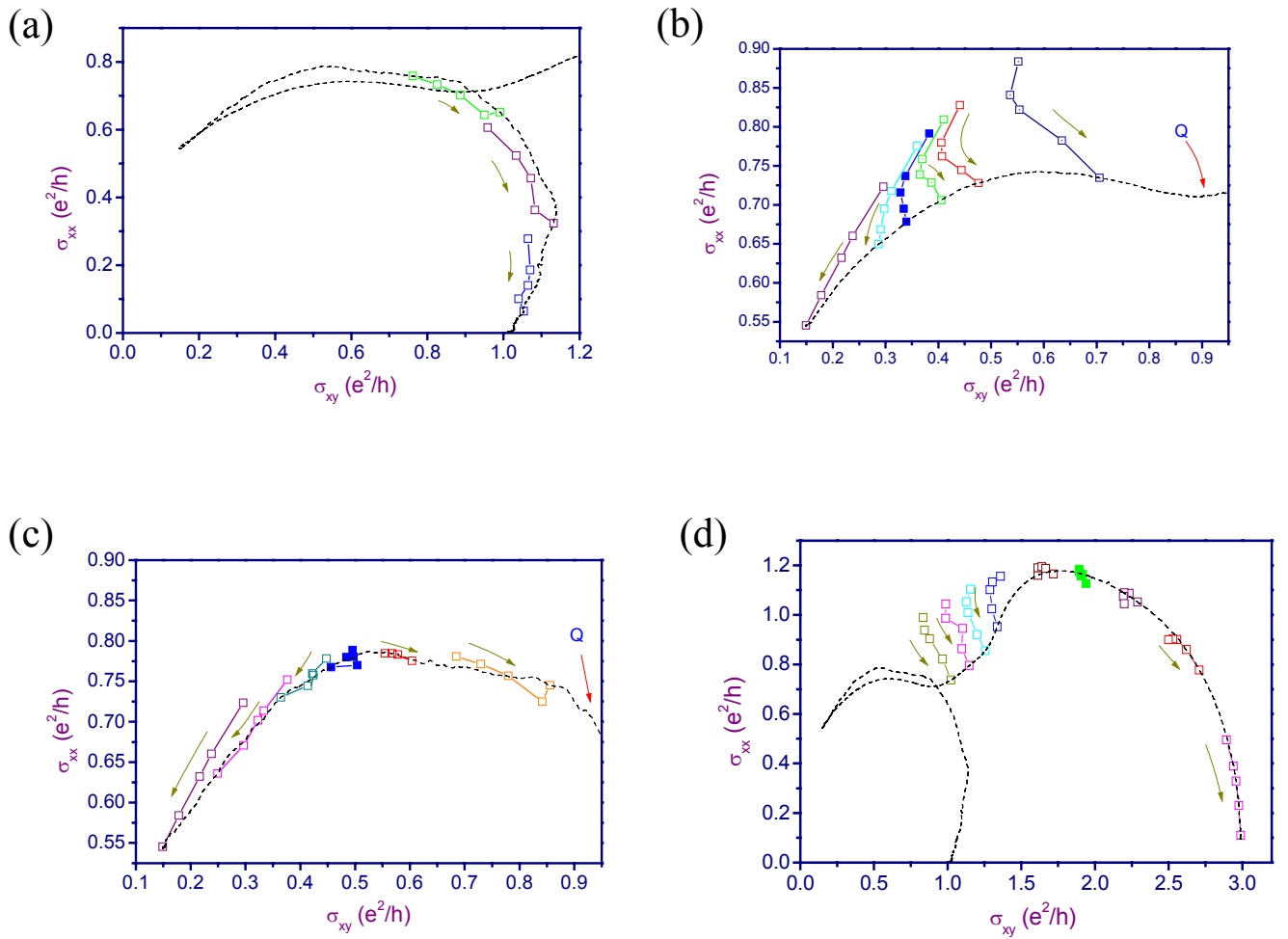


Figure 2 (a)-(d) Temperature-driven flow lines between  $1 < \nu < 3$ . The arrows indicate the directions of the flow lines (From high-T to low-T) and the full squares correspond to intermediate-coupling fixed points. The dotted line represents  $\sigma_{xx}(\sigma_{xy})$  at  $T = 0.3$  K.