

## Pseudoelectron Excitation in the Luttinger Liquid

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The concept of spin-charge separation plays a central role in describing the low-energy physics near Fermi points in one-dimensional (1D) interacting electron gas, a typical example of the Luttinger liquid (LL). This concept may be confirmed in real materials by the recent high-resolution angular resolved photoemission spectroscopy (ARPES) in which the one-electron spectral function  $A(p,\omega)$  can be directly measured in the wide range of momentum  $p$  and energy  $\omega$ .

If  $|p|$  is not restricted to the region near the Fermi momentum  $p_F$ , the linear dispersion approximation, usually adopted in the LL theory, is not sufficient in appropriately obtaining  $A(p,\omega)$ . In fact, the effect of the nonlinearity in the electron dispersion on  $A(p,\omega)$  has been intensively studied in recent years [1]. According to those studies on integrable systems,  $A(p,\omega)$  has singularities for arbitrary  $p$  at  $\omega=\varepsilon_s(p)$  and  $\omega=\varepsilon_c(p)$ , where  $\varepsilon_s(p)$  and  $\varepsilon_c(p)$  are the spectra of the spinon and holon modes, respectively. In the usual LL theory, the exponents characterizing these singularities are independent of  $p$ , but the nonlinearity in  $\varepsilon_s(p)$  and  $\varepsilon_c(p)$  makes them depend on  $p$ . Their actual  $p$  dependences can be calculated with use of the Bethe-ansatz method. For nonintegrable systems, the threshold singularity, which is located at  $\omega=\varepsilon_s(p)$ , remains intact, but the singularity at  $\omega=\varepsilon_c(p)$  is smeared into a broad peak.

In those preceding works, only the singularities at  $\omega=\varepsilon_s(p)$  and  $\omega=\varepsilon_c(p)$  are discussed on the belief that the electron nature will not sustain in the spin-charge separated system. For  $p$  far away from  $p_F$ , however, the effect of interactions becomes so weak that we would naively expect that the nature of an injected electron to measure  $A(p,\omega)$  manifests itself as a main peak in  $A(p,\omega)$ . Then a natural question arises: *Does an electron-like excitation mode actually exist in the 1D interacting electron gas for  $|p|\gg p_F$ ? If yes, a related and more intriguing question is: How does the electron-like mode reconcile with the physics of spin-charge separation for  $|p|\sim p_F$ ?*

In this study, by employing both analytical and numerical methods, we have reexamined  $A(p,\omega)$  in the Luttinger liquid with both linear and nonlinear electron dispersion at zero temperature. Irrespective of dispersion, we have discovered a conspicuous peak structure between two singularities at  $\omega=\varepsilon_s(p)$  and  $\omega=\varepsilon_c(p)$ . The structure is featured by a decay mode of an electron-like particle (*pseudoelectron*) which may be regarded as an electron dressed with a "cloud" of spinons and holons. This pseudoelectron exists even for  $|p|\sim p_F$ , is considered as the 1D counterpart of the Landau's quasiparticle in higher dimensions, and appears in  $A(p,\omega)$  as a peak with a cusp in addition to the well-known double-peak divergences representing the low-energy spin-charge separation. For  $|p|\gg p_F$ , on the other hand, it becomes a main contribution to  $A(p,\omega)$  and behaves as a nearly free electron. This 1D nearly free electron does not, however, manifest itself in  $A(p,\omega)$  as a Lorentzian (as in the case of higher dimensions) but divergent peak. Those results answer the questions posed in the previous paragraph. In our talk, we also give a full picture of  $A(p,\omega)$  in 1D interacting electron gas, including a transition between the spin-charge separated and nearly-free electron regimes.

[1] For a review, see A. Imambekov T. L. Schmidt, and L. I. Glazman, *Rev. Mod. Phys.* **84**, 1253 (2012).