

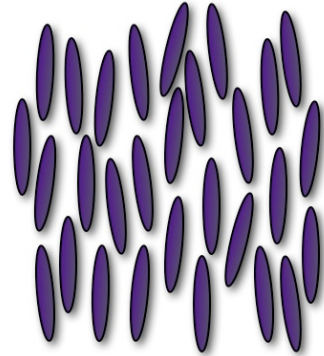
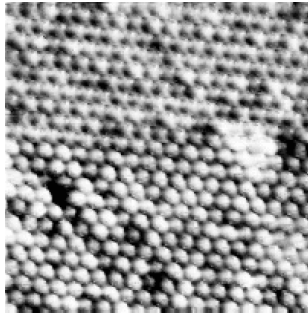
Symmetry Protected and Symmetry Enriched Topological Phases

Lukasz Fidkowski

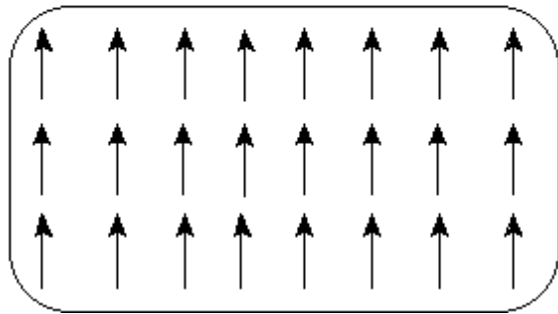
with N. Lindner, A. Kitaev

F. Burnell, X. Chen, A. Vishwanath

Symmetry breaking classification:



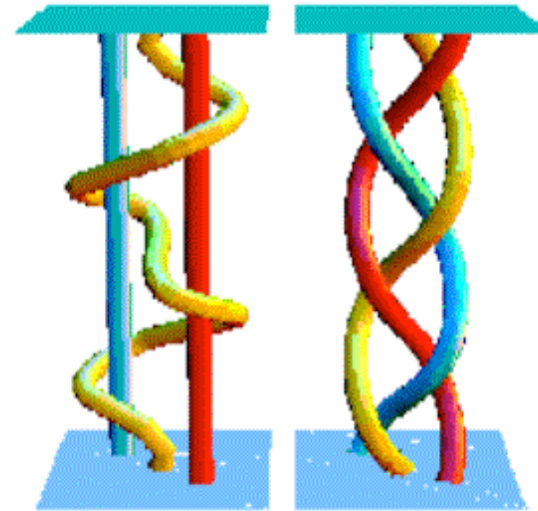
Nematic



- Ginzburg-Landau-Wilson
- local order parameters

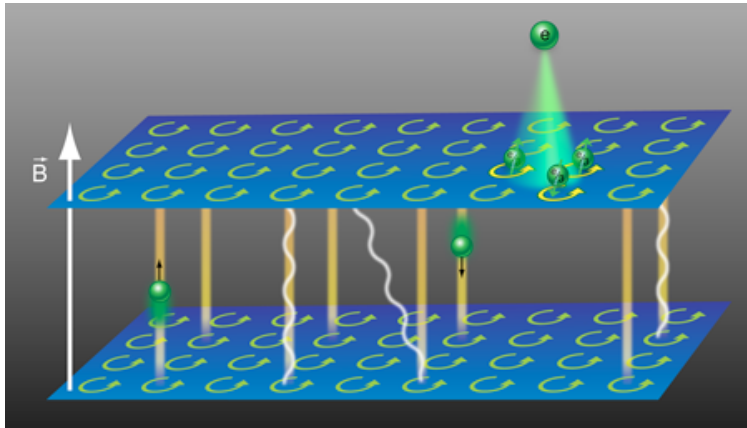
Topological phases:

2d gapped

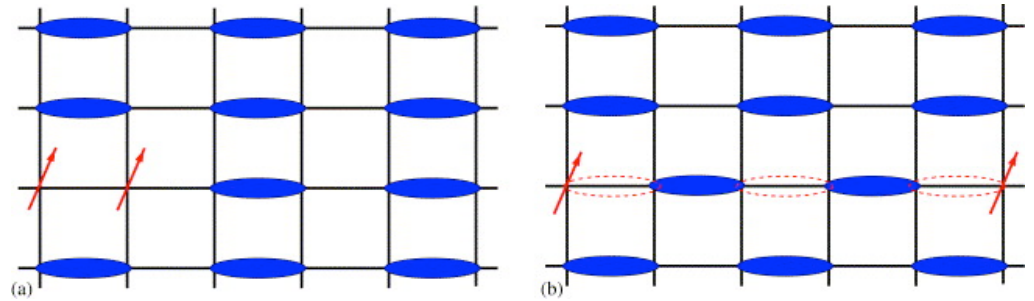


- deconfined fractionalized quasiparticles

Fractionalized:

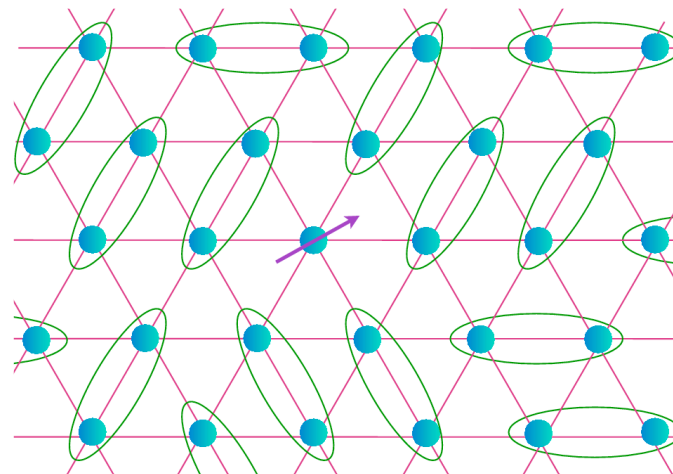


Laughlin quasiparticles



spinons (not quite)

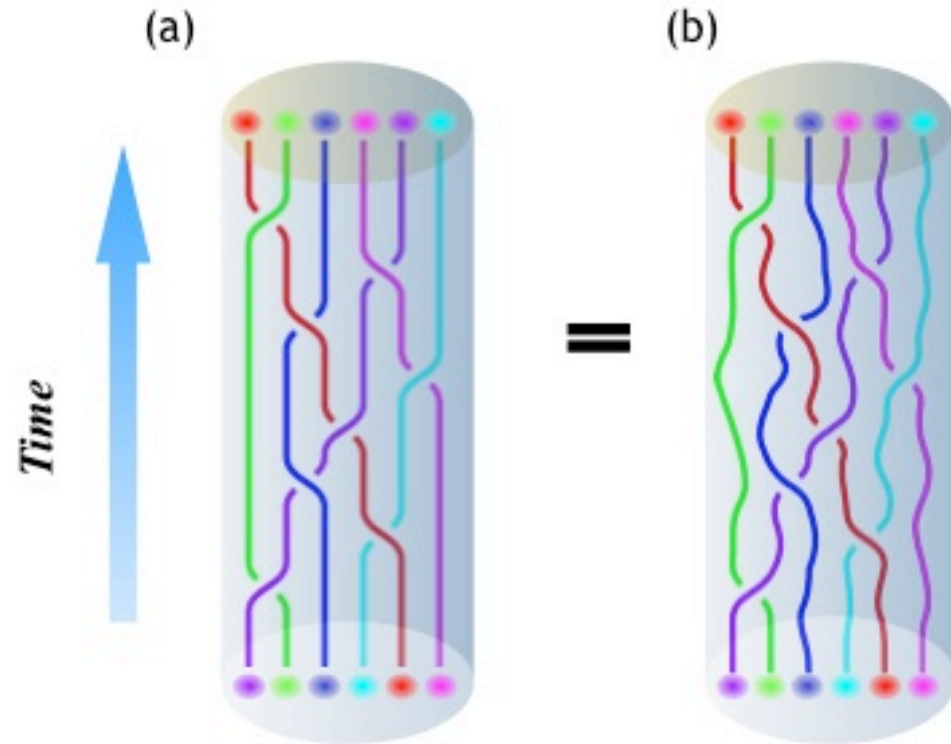
Deconfined:



deconfined spinon

Braiding statistics:

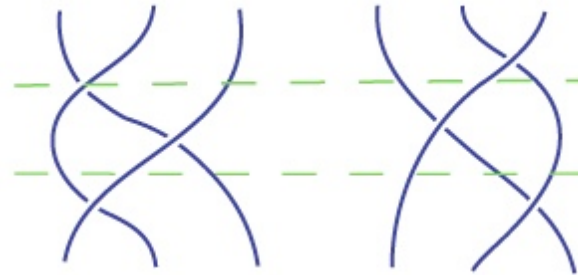
- more than fermions / bosons
- statistics can be a more general complex phase, or even matrix
- “anyons” (Wilczek)



Rigidity:

-“anyon” is a bit of a misnomer: quasiparticle braiding structure is very rigid.

- consistency conditions, eg.:

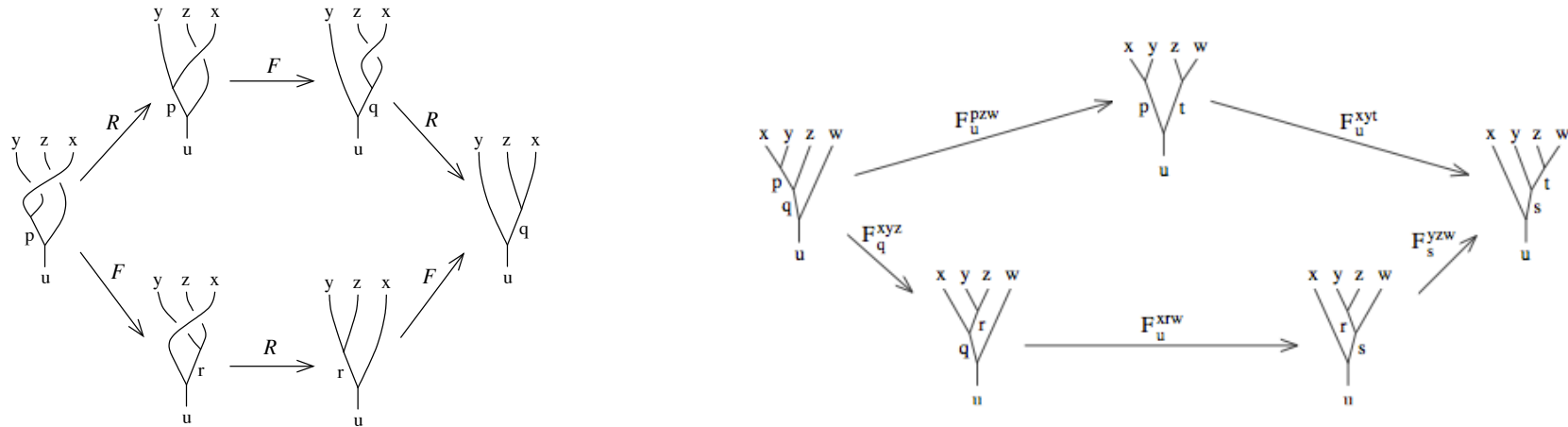


$$R_{12} R_{13} R_{23} = R_{23} R_{13} R_{12},$$

- structure encoded in numbers (or matrices):

$$R_{ij} \text{ and } F_{ijk}$$

- resulting structure: “Unitary Modular Tensor Category” (UMTC):



- the solutions are discrete, so the particular solution is itself a discrete invariant of the gapped phase
- works for “bosonic” systems - i.e. spin systems (can be extended to fermions)
- Hence, can classify bosonic “intrinsic topological order” by UMTCs. (Moore & Read, Wen, etc.)

Symmetries:

- symmetry group G : $[H, U_g] = 0$
- “Symmetry Protected Topological Order”: no symmetry breaking, no deconfined anyons
- classified by group cohomology $H^{d+1}(G, U(1))$

(1d: AKLT; Chen, Gu, Wen; Fidkowski, Kitaev; Pollman, Berg, Turner, Oshikawa; Hatsugai et al.; 2d: Chen, Gu, Liu, Wen; Levin, Gu)

Remainder of talk:

- only gapped 2d bosonic (i.e. spin) systems
- intrinsic topological order + symmetries = “symmetry enriched topological order”
 - classification and physical picture
(Essin, Hermele; Mesaros, Ran; Lu, Vishwanath, etc.)
- Subtle obstructions: some seemingly OK states are physically inconsistent

Eg. $G=SO(3)$ (i.e. integral spins)

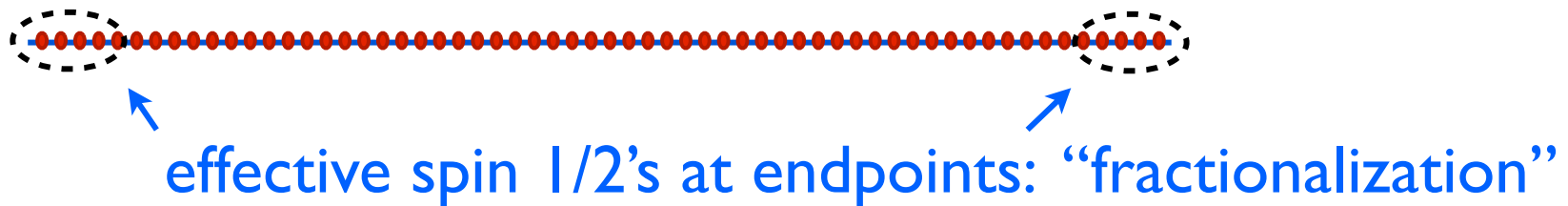
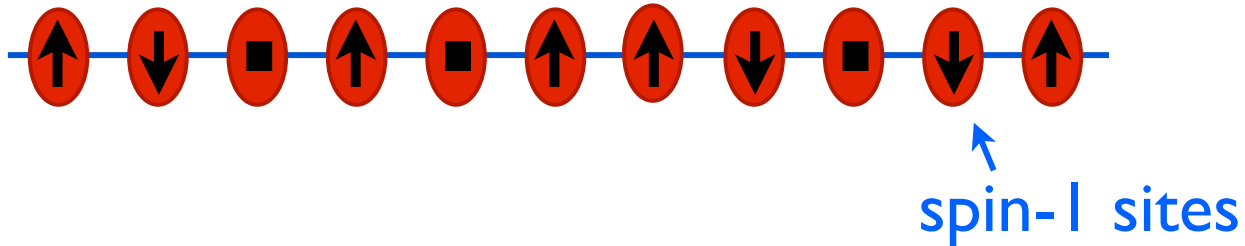
- \mathbb{Z}_2 gauge theory with spin-1/2 spinon **OK**

- Chiral spin liquid ($\nu = 1/2$ bosonic FQH) with spin-1/2 semion.

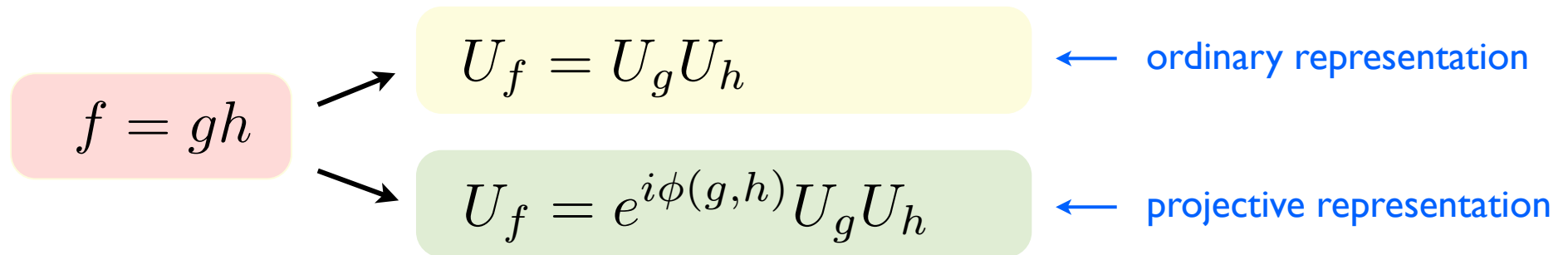
INCONSISTENT

1 dimension:

$G=SO(3)$, Haldane chain: $H = J \sum_j \vec{S}_j \cdot \vec{S}_{j+1}$



Projective representations of G

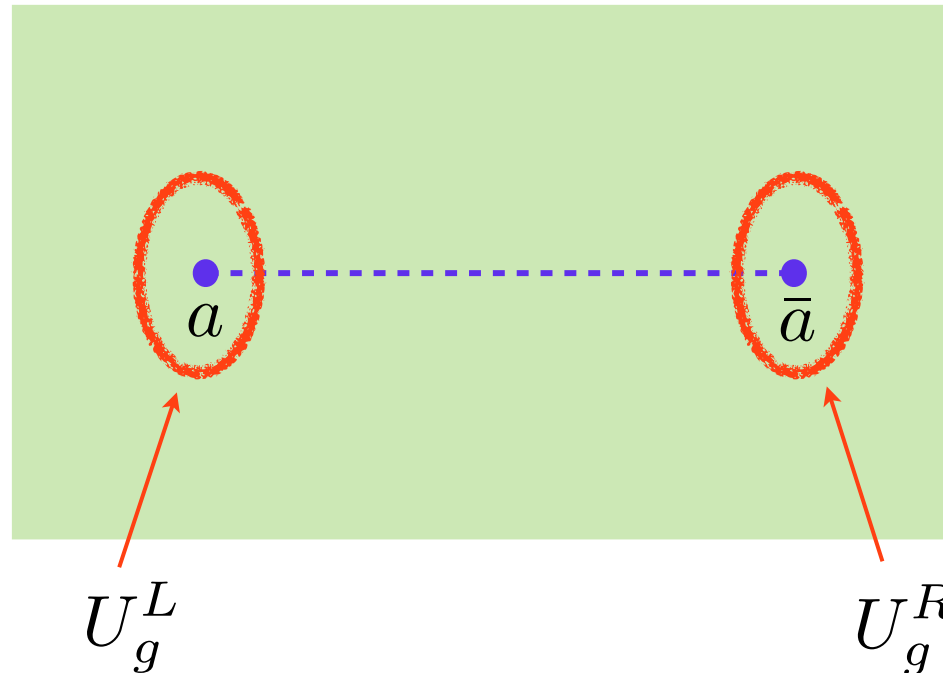


$\phi(g, h)$: function of 2 group variables \Rightarrow

$$H^2(G, U(1))$$

2 dimensions:

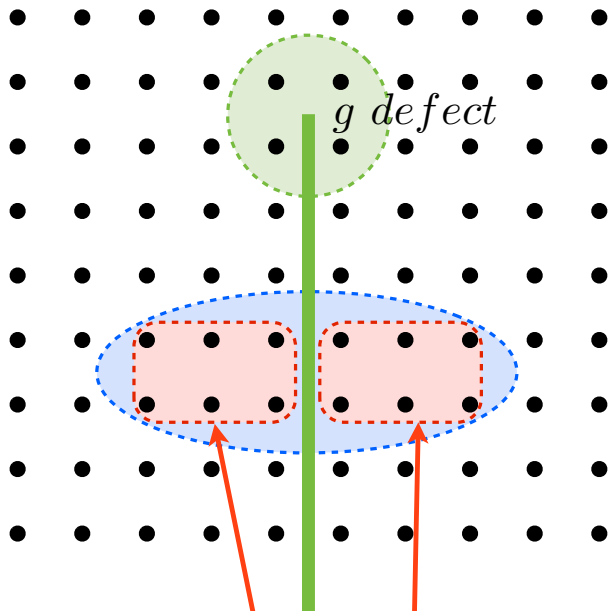
- anyons can carry fractional/projective quantum numbers of G :



$$U_{gh}^a = e^{i\phi_a(g,h)} U_g^a U_h^a$$

constraint: $\phi_a(g, h) + \phi_b(g, h) = \phi_c(g, h) \Rightarrow H^2(G, \text{abelian anyons})$
 $a \times b = c$

G-defects:



$$U_{fh}^{(g)} = e^{i\phi_g(f,h)} U_f^{(g)} U_h^{(g)}$$

↓ leads to

$$H^3(G, U(1))$$

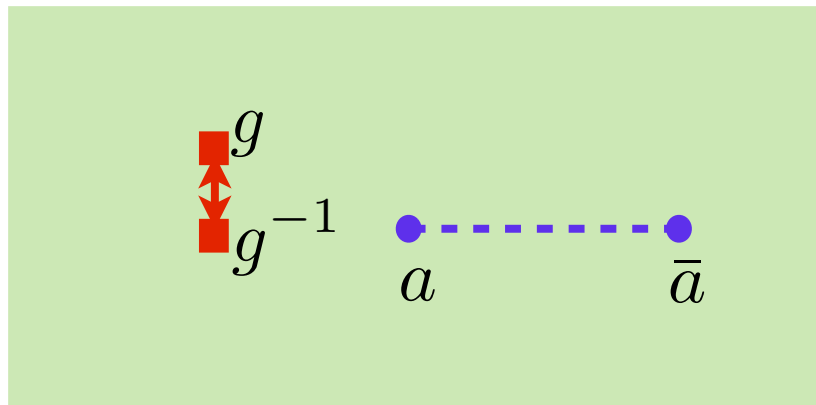
$$H_i = \sum_{\alpha} H_i^{L,\alpha} H_i^{R,\alpha}$$

↓

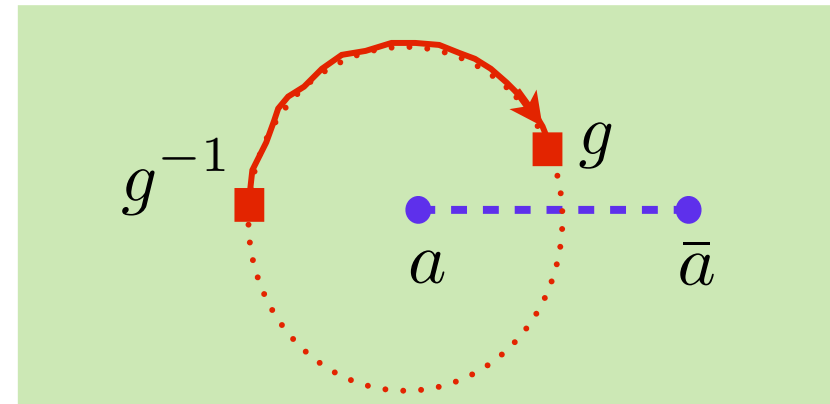
$$H_i^{defect} = \sum H_i^{L,\alpha} \left(\mathcal{U}_g^{-1} H_i^{R,\alpha} \mathcal{U}_g \right)$$

Local action of G via adiabatic process:

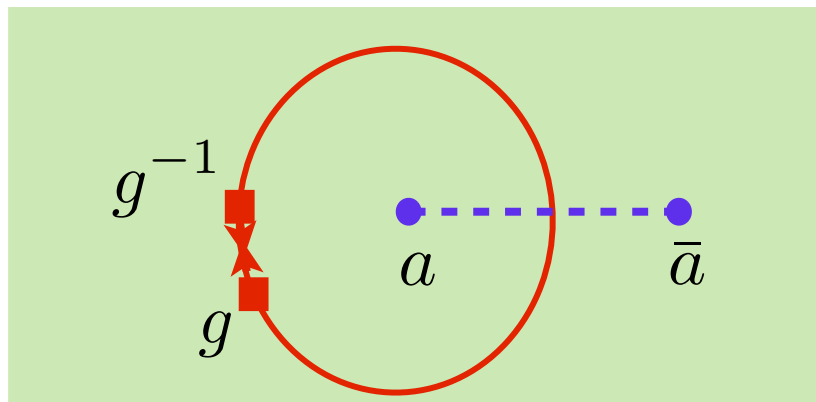
1. nucleate defect / anti-defect pair:



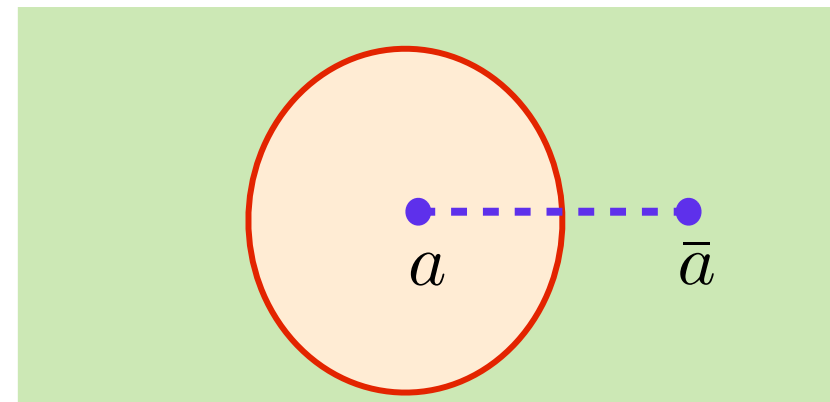
2. take one defect around anyon:



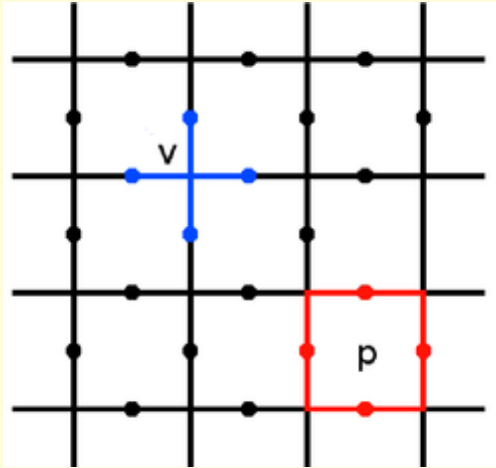
3. Annihilate defects:



3. Act with g on spins inside defect loop to remove branch cut



Example: “deformed” toric code: (= \mathbb{Z}_2 gauge theory)



$$A_v = \prod_{i \in v} \sigma_i^x, \quad B_p = \prod_{i \in p} \sigma_i^z$$

$$H = -J \sum_v A_v + J \sum_p B_p, \quad J > 0$$

note plus sign

$G = \text{spatial translations} = \mathbb{Z} \times \mathbb{Z}$

generators: T_x, T_y

Ground state of “deformed” toric code:

- ordinary toric code:

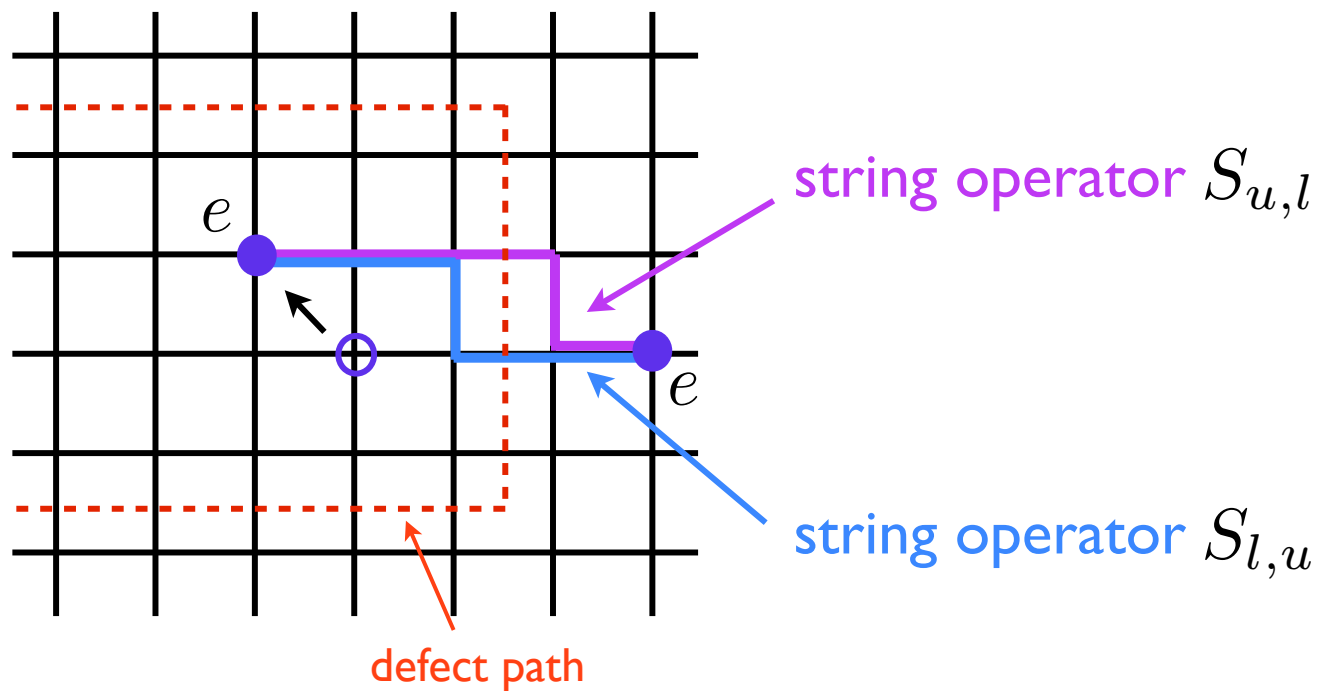
$$|\Psi\rangle = \sum_{\text{loop configs } L} |L\rangle$$

- “deformed” toric code:

$$|\Psi\rangle = \sum_{\text{loop configs } L} (-1)^{|L|} |L\rangle$$

area of loops in L (well defined with periodic boundary conditions and even # of sites)

Electric charge is projective under $\mathbb{Z} \times \mathbb{Z}$



$$U_{T_x}^L U_{T_y}^L |\Psi\rangle = S_{l,u} |0\rangle$$

$$U_{T_y}^L U_{T_x}^L |\Psi\rangle = S_{u,l} |0\rangle$$

but $S_{u,l} |0\rangle = -S_{l,u} |0\rangle$

Example 2:

$G = \mathbb{Z}_n$: fractional charges

$$U_{\frac{[k]}{n}} |a\rangle = e^{2\pi i \frac{[k]}{n} q} |a\rangle \leftarrow \text{a has charge } q$$

$[k]$ means $(k \bmod n)$

$$U_{\frac{[k+l]}{n}} |a\rangle = e^{2\pi i \frac{[k+l]}{n} q} |a\rangle$$

$$U_{\frac{[k]}{n}} U_{\frac{[l]}{n}} |a\rangle = e^{2\pi i \frac{[k]+[l]}{n} q} |a\rangle$$

phase mismatch when q is fractional

Gauge G:

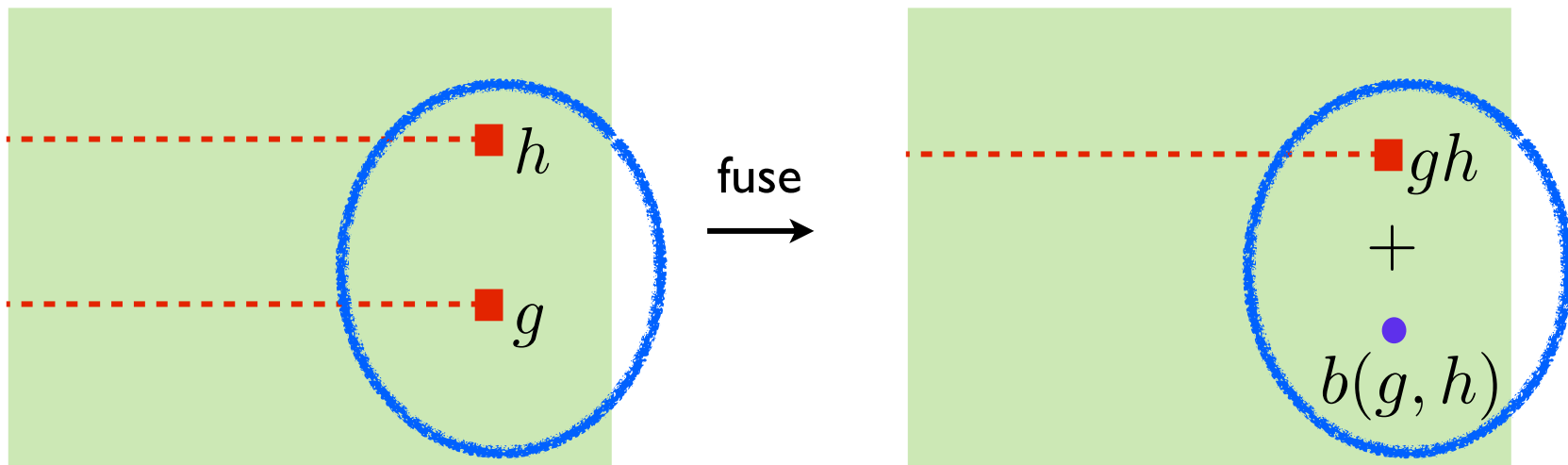
G-fluxes, G-charges become deconfined excitations:

quasiparticles = {original anyons, G-fluxes, G-charges}

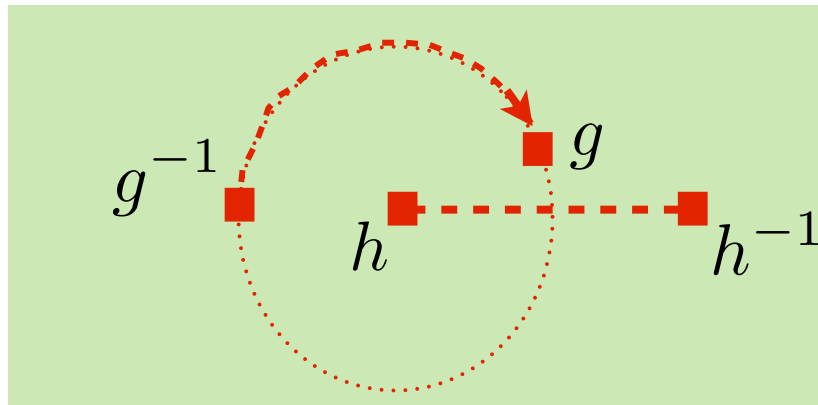
Then the projective character of the anyons and defects is reflected in the fusion/braiding structure of this enlarged theory.

(see also Levin & Gu, arXiv:1202.3120)

1) projective/fractional anyons \Rightarrow defect fusion rules



2) projective/fractional defects => defect braiding rules



Any consistent theory for fusion/braiding in enlarged theory must come from some some choice of these two invariants

(Etingof et al., Kirillov, Mueger)

(can also handle case when G changes anyon superselection sectors)

Connection to 3D SPT's:

Is every choice of $H^2(G, A)$ and $H^3(G, U(1))$ realized in some 2d theory?

No: sometimes there is an obstruction to a consistent choice of braiding and fusion rules in the enlarged theory

Example:

anyons = $\{1, a\}$ (a =semion)

abelian theory, $K=(2)$

$\nu = 1/2$ bosonic FQHE

$G = \mathbb{Z}_2 \times \mathbb{Z}_2$, a is projective under G

$G = SO(3)$ works too

Although this theory cannot be realized in 2d, we have an exactly solvable model which realizes it at the surface of a 3d SPT.

(with F. Burnell, X. Chen, A. Vishwanath)

In general, the obstruction is in $H^4(G, U(1))$, which classifies 3d SPT's.

Conclusions:

- Classified symmetry enriched topological phases (SETs), given data of UMTC and symmetry group G
- Resulting SETs are in one to one correspondence with fusion/braiding structures for anyons + G -defects
- These SETs are parametrized by data that describes fractional/projective character of anyons and defects
- Not all choices of this data lead to consistent SETs: interpretation as “anomalous” realization of G at the surface of 3d SPT.