

Realistic many-body calculations with spatial correlations and for systems with molecular orbitals

Harald O. Jeschke

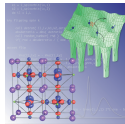
Johannes Ferber, Hunpyo Lee, Kateryna Foyevtsova, Roser Valentí

Institut für Theoretische Physik, Universität Frankfurt

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Emergent Quantum Phases in Condensed Matter - from topological to first principles approaches

June 10, 2013



Acknowledgments

Frankfurt:

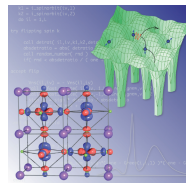
Hunpyo Lee, Kateryna Foyevtsova, Johannes Ferber, Milan Tomić, Roser Valentí.

Würzburg (experiment):

Andreas Ruff, Michael Sing, Ralph Claessen.

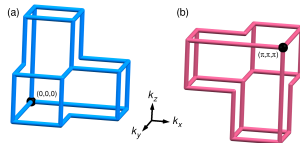
Funding:

- DFG FOR 1346 “Dynamical Mean Field Approach with Predictive Power for Strongly Correlated Materials”
- DFG TR 49 “Condensed Matter Systems with Variable Many-Body Interactions”

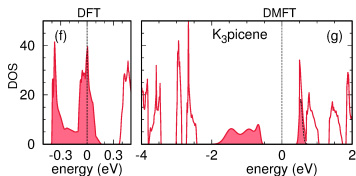


Topics

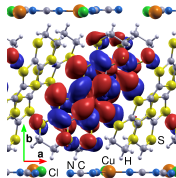
Introducing spatial correlations into realistic DMFT studies



Potassium doped picene: metallic or not?



Projectors for molecular orbitals: LDA+DMFT description of charge transfer salts



Introducing spatial correlations into realistic DMFT studies

Wanted: Momentum dependent selfenergy in realistic manybody calculations

- LDA+single site DMFT derives k dependence only from DFT;
 $\Sigma(\mathbf{k}, \omega) \equiv \Sigma(\omega)$
- Many systems with dimension < 3 show phenomena that are not captured by single site DMFT (cuprates, organic charge transfer salts)

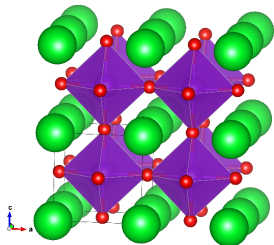
Hope comes from various recent developments:

- 1 dual fermions (DFT+DF) (Lichtenstein, Katsnelson, Hafermann, ...)
- 2 variational cluster approach (DFT+VCA) (Aichhorn, ...)
- 3 DFT+DVA (Held, ...)
- 4 GW+DMFT (Biermann, Georges, Held, ...)

Here, we try to combine DFT with DCA.

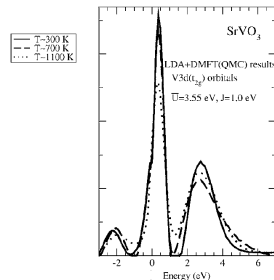
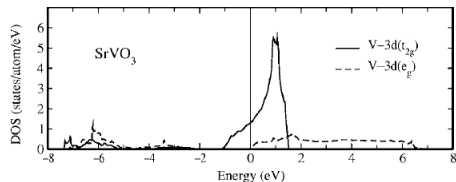
SrVO₃ as a test system

- SrVO₃ is a cubic perovskite



- V t_{2g} with $1/6$ filling
- correlated metal studied many times with LDA+DMFT
- intensively investigated experimentally; latest addition Aizaki *et al.*, PRL **109**, 056401 (2012).

Typical DFT, LDA+DMFT results:



Nekrasov *et al.*, Phys. Rev. B **72**, 155106 (2005).

LDA+DCA Method

Projective Wannier functions within FLAPW (Aichhorn et al., PRB **80**, 085101 (2009)):

$$|\chi_{\mathbf{k},m}^{\alpha,\sigma}\rangle = \sum_{\nu \in W} \langle \psi_{\mathbf{k},\nu}^{\sigma} | \chi_m^{\alpha,\sigma} \rangle | \psi_{\mathbf{k},\nu}^{\sigma} \rangle,$$

with atom α , band ν , spin σ , orbital m , Bloch eigenfunction $|\psi_{\mathbf{k},\nu}^{\sigma}\rangle$, correlated orbital

$|\chi_m^{\alpha,\sigma}\rangle = |u_l^{\alpha,\sigma}(E_l)Y_m^l\rangle$, linearization energies E_l , radial wave function $u_l^{\alpha,\sigma}$, and spherical harmonic function Y_m^l .

Orthonormalized projectors:

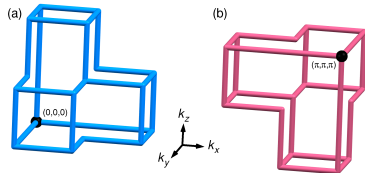
$$P_{m,\nu}^{\alpha,\sigma}(\mathbf{k}) = \sum_{\alpha',m'} \langle u_l^{\alpha',\sigma}(E_l)Y_m^l | \psi_{\mathbf{k},\nu}^{\sigma} \rangle [O(\mathbf{k},\sigma)^{-1/2}]_{m,m'}^{\alpha,\alpha'}$$

LDA+DCA lattice Greens function (cluster momenta \mathbf{K})

$$G_{\nu,\nu'}^{\sigma}(\mathbf{K} + \tilde{\mathbf{k}}, i\omega_n) = \left[i\omega_n + \mu - \epsilon_{\mathbf{K} + \tilde{\mathbf{k}},\nu}^{\sigma} - \Sigma_{\nu,\nu'}^{\sigma}(\mathbf{K} + \tilde{\mathbf{k}}, i\omega_n) \right]^{-1}$$

Lattice self-energy

$$\begin{aligned} & \Sigma_{\nu,\nu'}^{\sigma}(\mathbf{K} + \tilde{\mathbf{k}}, i\omega_n) \\ &= \sum_{\alpha,m,m'} P_{\nu,m}^{\alpha,\sigma*}(\mathbf{K} + \tilde{\mathbf{k}}) \Delta \Sigma_{m,m'}^{\sigma,\text{imp}}(\mathbf{K}, i\omega_n) P_{m',\nu'}^{\alpha,\sigma}(\mathbf{K} + \tilde{\mathbf{k}}). \end{aligned}$$



Brillouin zone sectors

Double counting correction

$$\Delta \Sigma_{m,m'}^{\sigma,\text{imp}}(\mathbf{K}, i\omega_n) = \Sigma_{m,m'}^{\sigma,\text{imp}}(\mathbf{K}, i\omega_n) - \Sigma_{m,m'}^{\text{dc}}$$

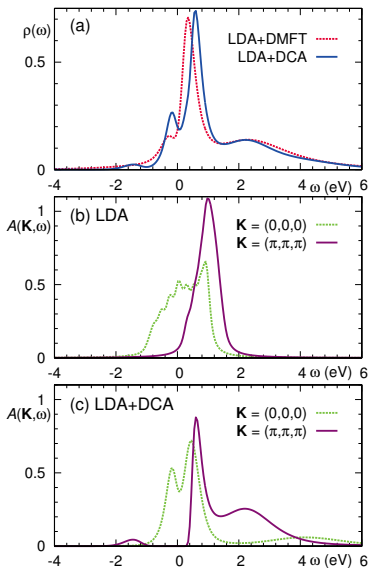
Local cluster Greens function

$$\begin{aligned} G_{m,m'}^{\sigma,\text{loc}}(\mathbf{K}, i\omega_n) &= \sum_{\tilde{\mathbf{k}},\nu,\nu'} P_{m,\nu}^{\alpha,\sigma}(\mathbf{K} + \tilde{\mathbf{k}}) \\ &\quad \times G_{\nu,\nu'}^{\sigma}(\mathbf{K} + \tilde{\mathbf{k}}, i\omega_n) P_{\nu',m'}^{\alpha,\sigma*}(\mathbf{K} + \tilde{\mathbf{k}}), \end{aligned}$$

Weiss field update

$$\begin{aligned} & [G_{m,m'}^{\sigma,0}(\mathbf{K}, i\omega_n)]^{-1} \\ &= \Sigma_{m,m'}^{\sigma,\text{imp}}(\mathbf{K}, i\omega_n) + [G_{m,m'}^{\sigma,\text{loc}}(\mathbf{K}, i\omega_n)]^{-1}. \end{aligned}$$

Comparison between LDA+DMFT and LDA+DCA



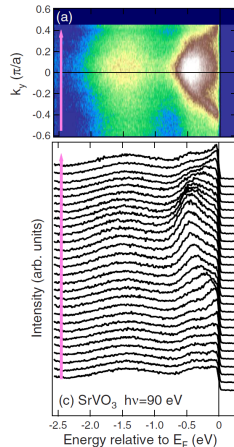
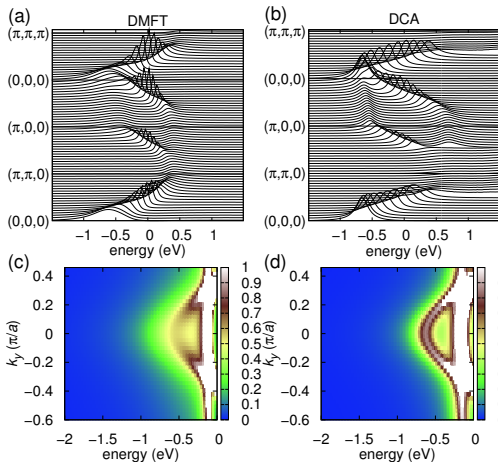
Interaction Hamiltonian to be solved:

$$H_I = U \sum_m n_{m\uparrow} n_{m\downarrow} + \sum_{m < n, \sigma} [U' n_{m\sigma} n_{n\bar{\sigma}} + (U' - J) n_{m\sigma} n_{n\sigma}],$$

- LDA+DMFT, calculated with hybridization expansion CT-QMC (from ALPS), agrees well with result of Aichhorn *et al.*, PRB **80**, 085101 (2009).
- LDA+DCA, calculated with interaction expansion CT-QMC, shows some extra structure.
- Understandable from \mathbf{K} sector resolved spectral functions.
- $\Sigma_{m,m'}^{\sigma, \text{imp}}(\mathbf{K}, i\omega_n)$ different for the two sectors

Compare also Zhang, Imada, PRB **76**, 045108 (2007);
Lin, Millis, PRB **79**, 205109 (2009).

Comparison to experiment



LDA+DMFT (left) and LDA+DCA (right)
comparison; (c), (d) at $k_z = 0.32\pi/a$

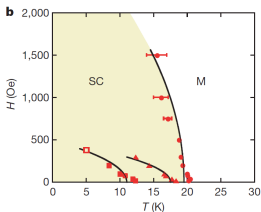
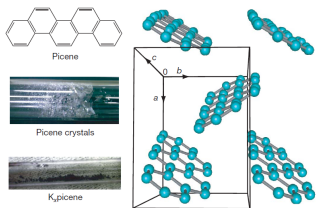
Lee, Foyevtsova, Ferber, Aichhorn, Jeschke, Valentí, Phys. Rev. B **85**, 165103 (2012).

Yoshida *et al.*, Phys. Rev. B
82, 085119 (2010).

Potassium doped picene: metallic or not?

Doping of organic molecular crystals

Superconductivity in K_3 picene with $T_c = 18$ K:

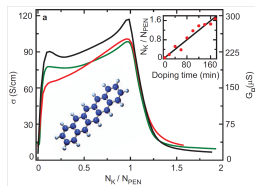
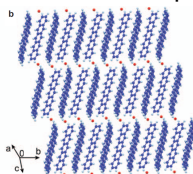


Mitsuhashi *et al.*, Nature **464**, 76 (2010)

- SC subsequently found in phenanthrene (5 K) and dibenzopentacene (33 K)
- Indication for Mott physics in doped organics: compare case of pentacene

- unexpected
- observed through Meissner effect
- recently: zero conductivity at 7 K (PRB 2013)

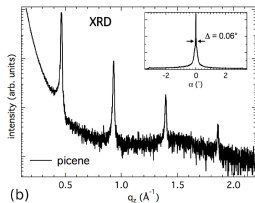
K intercalated pentacene:



Craciun *et al.*, Phys. Rev. B **79**, 125116 (2009).

Problem for theory: Explain metallic state

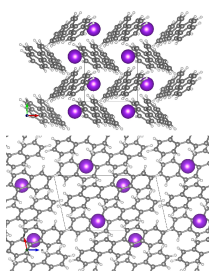
K_x picene structure and electronic structure



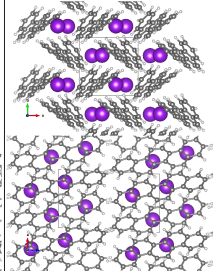
(b) Structure not known experimentally

→ need to simulate

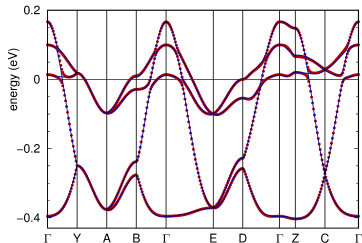
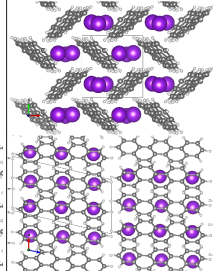
(a) K_1 picene



(b) K_2 picene



(c) K_3 picene

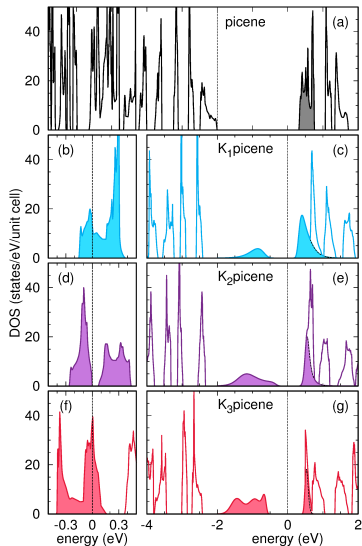
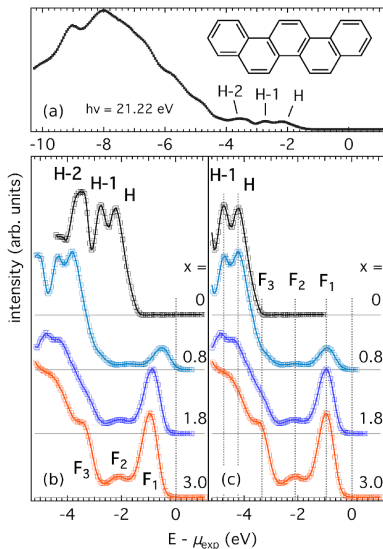


Hamiltonian we solve is

$$H_K = - \sum_{ijm\sigma} t_m c_{mj\sigma}^\dagger c_{mi\sigma} - t_\perp \sum_{i\sigma} c_{1i\sigma}^\dagger c_{2i\sigma}$$

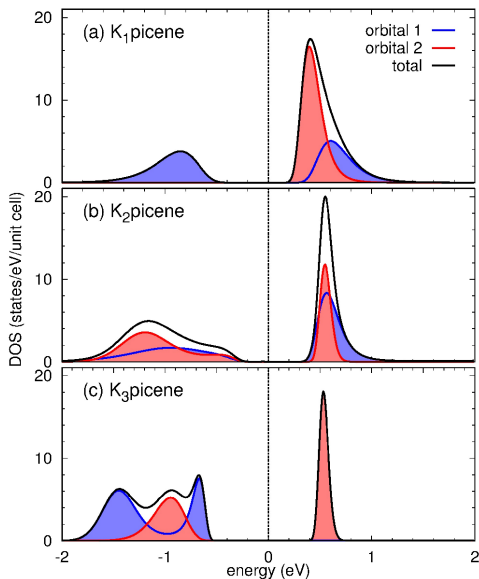
$$H_I = U \sum_{mi} n_{mi\uparrow} n_{mi\downarrow} + \sum_{\sigma} \left[(U - 2J_z) n_{1i\sigma} n_{2i\bar{\sigma}} + (U - 3J_z) n_{1i\sigma} n_{2i\sigma} \right].$$

K_x picene: Photoemission and DFT/TB+DMFT



Ruff, Sing, Claessen, Lee, Tomić, Jeschke, Valentí, PRL **110**, 216403 (2013).

K_x picene: DMFT for two orbital model

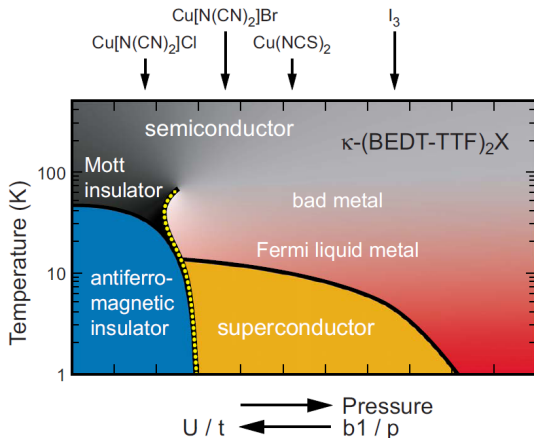


- essential ingredient: two orbital model with hybridization between orbitals
- off-diagonal Greens functions
→ use of interaction expansion CT-QMC
- both bands gapped in K_2 picene
- no metallic state for commensurate doping
- source of superconductivity unclear

Ruff, Sing, Claessen, Lee, Tomić, Jeschke, Valentí, PRL **110**, 216403 (2013).

Projectors for molecular orbitals: LDA+DMFT description of charge transfer salts

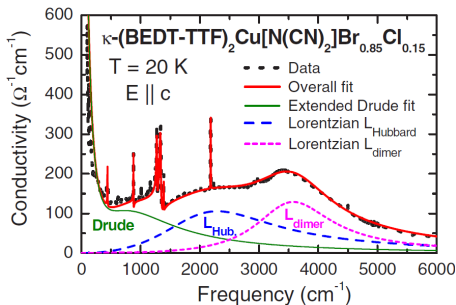
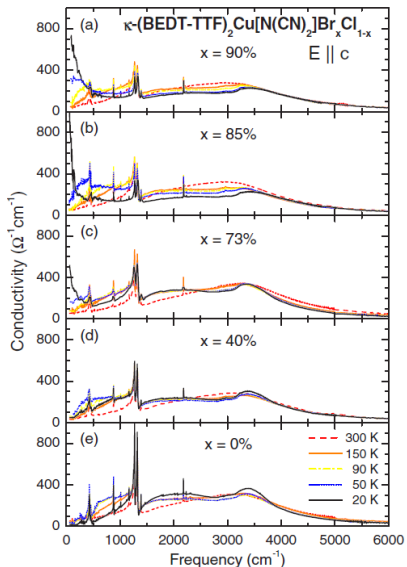
κ -phase BEDT-TTF based charge transfer salts



Dumm *et al.*, Phys. Rev. B **79**, 195106 (2009).

- Rich phase diagram accessible by chemical or physical pressure
- Extensively studied by theory
- So far: Model calculations, generic or based on Hückel/DFT extracted parameters, solved with myriad manybody methods.
- **Plan: apply full LDA+DMFT method.**

κ -(BEDT-TTF)₂-Cu[N(CN)₂]Br_xCl_{1-x} optical conductivity



Faltermeier *et al.*, PRB **76**, 165113 (2007).

- Drude peak in κ -Br at low T , disappears at high T
- absence of Drude peak in κ -Cl
- conjecture about composition of midinfrared spectrum based on ED model calc. (Merino, McKenzie)

LDA+DMFT for molecular orbitals

Idea for projection on molecular orbital: **Diagonalize occupation matrix in the basis of atomic orbitals within the correlated subspace.**

Starting point: FLAPW projectors $\tilde{P}_{m\nu}^{\alpha,\sigma}(\mathbf{k})$

$$O_{m,m'}^{\alpha,\alpha'}(\sigma) = \sum_{\nu \in W'} \tilde{P}_{m\nu}^{\alpha,\sigma}(0) \tilde{P}_{m'\nu}^{\alpha',\sigma*}(0).$$

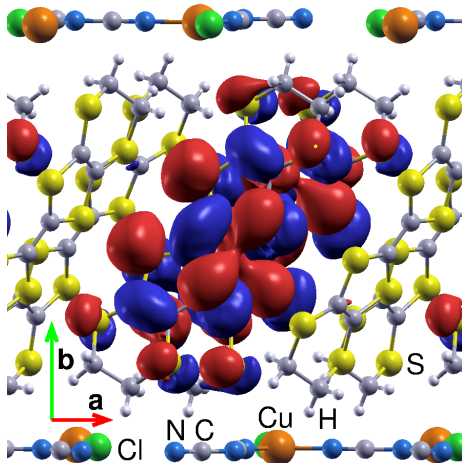
Narrow energy window W' , $\mathbf{k} = 0$, atomic orbitals $\{\alpha, m\}$.

Eigenvectors of $O_{m,m'}^{\alpha,\alpha'}(\sigma)$ corresponding to largest eigenvalues now indicate composition of molecular orbitals in terms of $\{\alpha, m\}$.

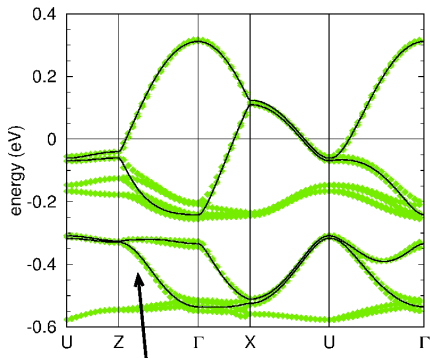
If $U_m^\alpha(\sigma)$ is such an eigenvalue,

$$\tilde{P}_{M\nu}^\sigma(\mathbf{k}) = \sum_{\alpha,m} U_m^\alpha(\sigma) \tilde{P}_{m\nu}^{\alpha,\sigma}(\mathbf{k})$$

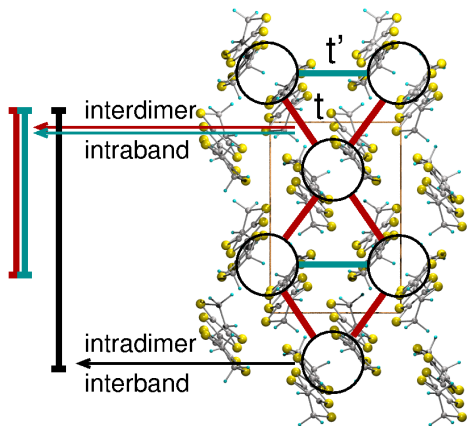
is a molecular projector for molecular orbital M.



κ -(BEDT-TTF)₂-Cu[N(CN)₂]Cl,
BEDT-TTF dimer Wannier function

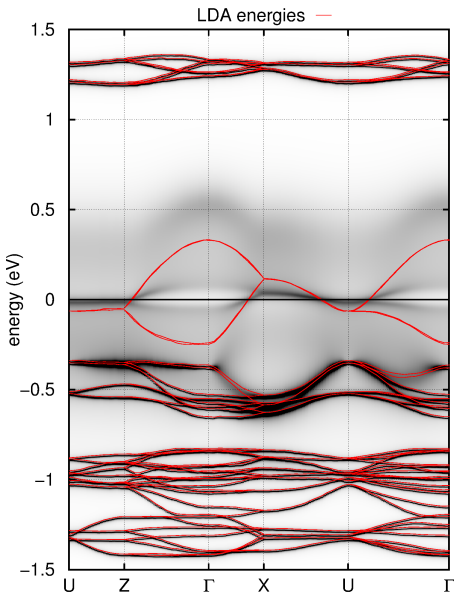
κ -(BEDT-TTF)₂X structure and bands κ -(ET)₂Cu[N(CN)₂]Br

black band count: 8 from ET HOMO
 or 4 from (ET)₂ HOMO-1, 4 from (ET)₂ HOMO



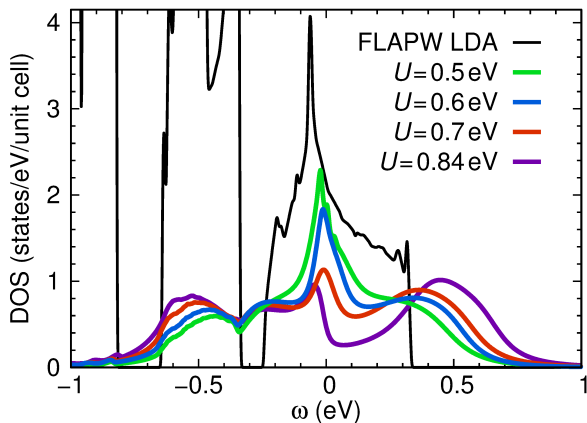
Kandpal, Opahle, Zhang, Jeschke, Valentí, PRL **103**, 067007 (2009).

κ -(ET)₂-Cu[N(CN)₂]Cl spectral function $A(\mathbf{k}, \omega)$



- Hybridization expansion CT-QMC impurity solver from ALPS project
- Temperature $T = 300$ K, $U = 0.6$ eV $\approx 10t$ (compare Nakamura *et al.*, J. Phys. Soc. Jpn. **78**, 083710 (2009)).
- Renormalized quasiparticle band and spectral weight transfer to lower and upper Hubbard bands

Ferber, Foyevtsova, Jeschke, Valentí,
arXiv:1209.4466.

κ -(ET)₂-Cu[N(CN)₂]Cl DOS and $A(\omega)$


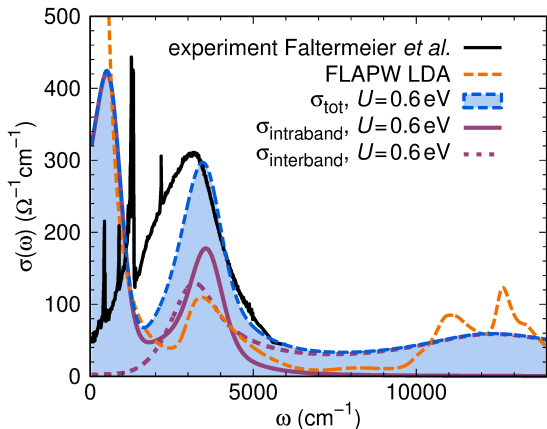
- evolution of quasiparticle peak and Hubbard bands with interaction strength U
- $U = 0.84$ eV close to insulating state
- some underestimation of critical U due to use of single site DMFT

Ferber, Foyevtsova, Jeschke, Valentí, arXiv:1209.4466.

Optical conductivity

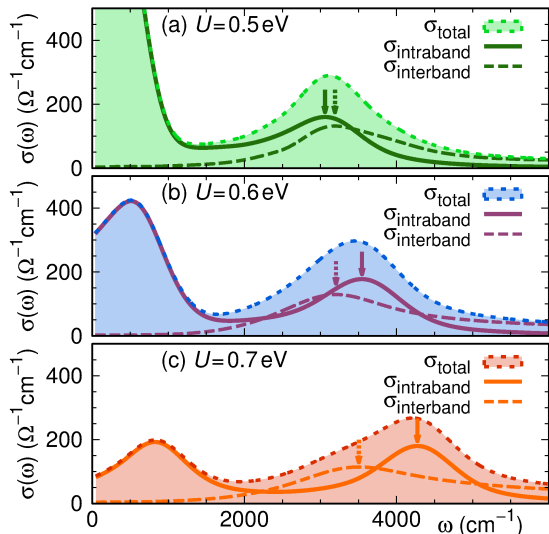
$$\sigma_{zz}(\tilde{\nu}_n) = \frac{e^2}{4\pi\epsilon_0 V \tilde{\nu}_n \beta} \sum_{\nu\nu'\nu''\nu''', \mathbf{k}, \sigma} v_{z,\mathbf{k}}^{\nu\nu'} v_{z,\mathbf{k}}^{\nu''\nu'''} \sum_{\omega_n} G_{\mathbf{k}}^{\nu'\nu''}(i\omega_n + i\tilde{\nu}_n) G_{\mathbf{k}}^{\nu'''\nu}(i\omega_n)$$

(lattice Greens function $G_{\mathbf{k}}(i\omega_n)$, bosonic (fermionic) Matsubara freq. $\tilde{\nu}_n$ (ω_n), Fermi velocity $v_{\alpha,\mathbf{k}}^{\nu\nu'}$).



- comparison to $T = 300 \text{ K}$ measured optical conductivity (Faltermeier *et al.*, Phys. Rev. B **76**, 165113 (2007)).
- interband (intradimer) contribution nearly similar to DMFT result
- DMFT shifts spectral weight from Drude peak into intraband (“Hubbard”) peak

Influence of the interaction strength



- Drude peak suppressed as interaction strength increases
- position of intraband (interdimer) peak shifts like Hubbard bands, proportional to $0.75U$ due to contributions from U and $U/2$
- position of interband (intradimer) peak nearly constant
- conjecture of Faltermeier *et al.* confirmed from first principles

Conclusions

LDA+DCA summary (Lee *et al.*, Phys. Rev. B **85**, 165103 (2012).)

- LDA+DCA introduces some momentum dependence into realistic calculations for strongly correlated materials.
- Low energy features seem to be described better in SrVO₃.

Theory&PES for K_xpicene summary (Ruff *et al.*, PRL **110**, 216403 (2013)).

- Insulating state found in photoemission and DFT/TB+DMFT theory
- Importance of multiorbital nature

LDA+DMFT for CT salts summary (Ferber *et al.*, arXiv:1209.4466).

- Projector scheme for molecular orbitals within FLAPW allows LDA+DMFT calculation for Mott insulator κ -(ET)₂-Cu[N(CN)₂]Cl.
- Interband and intraband contributions to optical conductivity described at the same level.