

# Field Induced Reentrant Phase Transitions in Quasi-One-Dimensional Random Quantum Heisenberg Chains Saitama University Kazuo HIDA

## 1 Introduction

### 1.1 Quantum Magnets in the Magnetic Field

#### 1. Magnetization Plateau

Field induced spin gap = Disordered Phase

#### 2. Field Induced Antiferromagnetic Phase

##### (a) Quasi-One-Dimensional System

Magnetic Field  $>$  Spin Gap  $\Rightarrow$  Tomonaga-Luttinger Liquid  
+ Weak interchain interaction



Antiferromagnetic Long Range Order

## Disorder effects on the quantum spin systems in the magnetic field

- Revived spins  $\Rightarrow$  **magnetic long range order**
- Localized singlet pairs = plateau formation  $\Rightarrow$  **Suppress the magnetic order**

Competition



Field induced reentrant transition

## 2 Model

Quasi-1-dimensional Random  $S = 1/2$  Heisenberg model with bond alternation

$$H = \sum_j \left\{ \sum_{i=1}^{N/2} J \mathbf{S}_{2i-1,j} \mathbf{S}_{2i,j} + \sum_{i=1}^{N/2} J_{ij} \mathbf{S}_{2i,j} \mathbf{S}_{2i+1,j} \right\}$$

intrachain

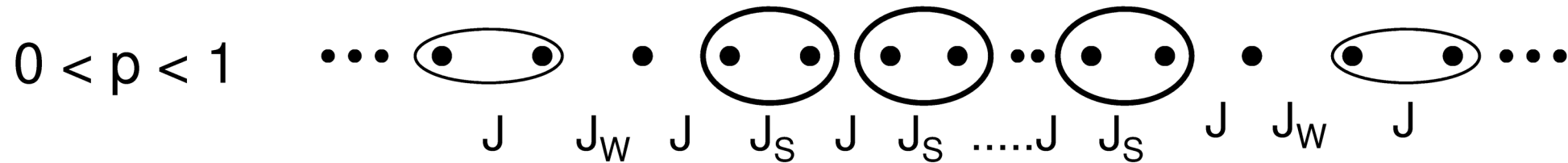
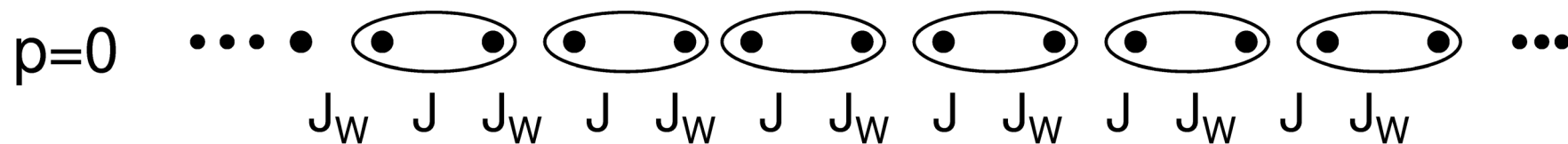
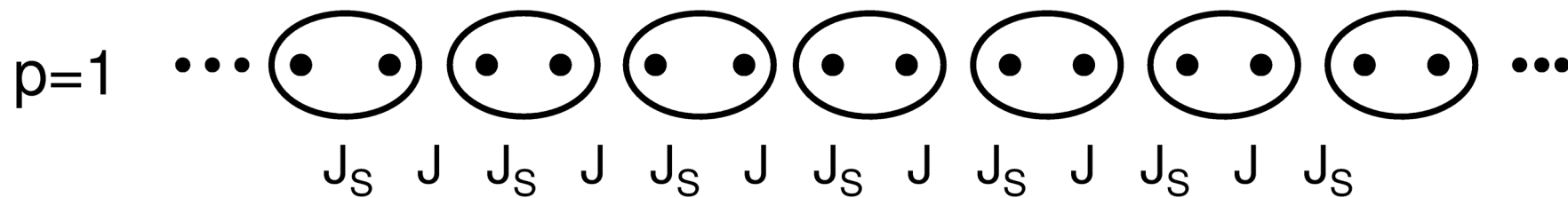
$$+ \sum_{i=1}^N \sum_{\langle j,j' \rangle} J_{\text{int}} \mathbf{S}_{i,j} \mathbf{S}_{i,j'} \quad \text{interchain}$$

- Intrachain Interaction:

$$J = 1 \quad J_{i,j} = \begin{cases} J_S & \text{probability } p \\ J_W & \text{probability } 1 - p \end{cases}$$
$$J_S > J > J_W > 0$$

Method DMRG: Keeping 60 ~ 160 states.

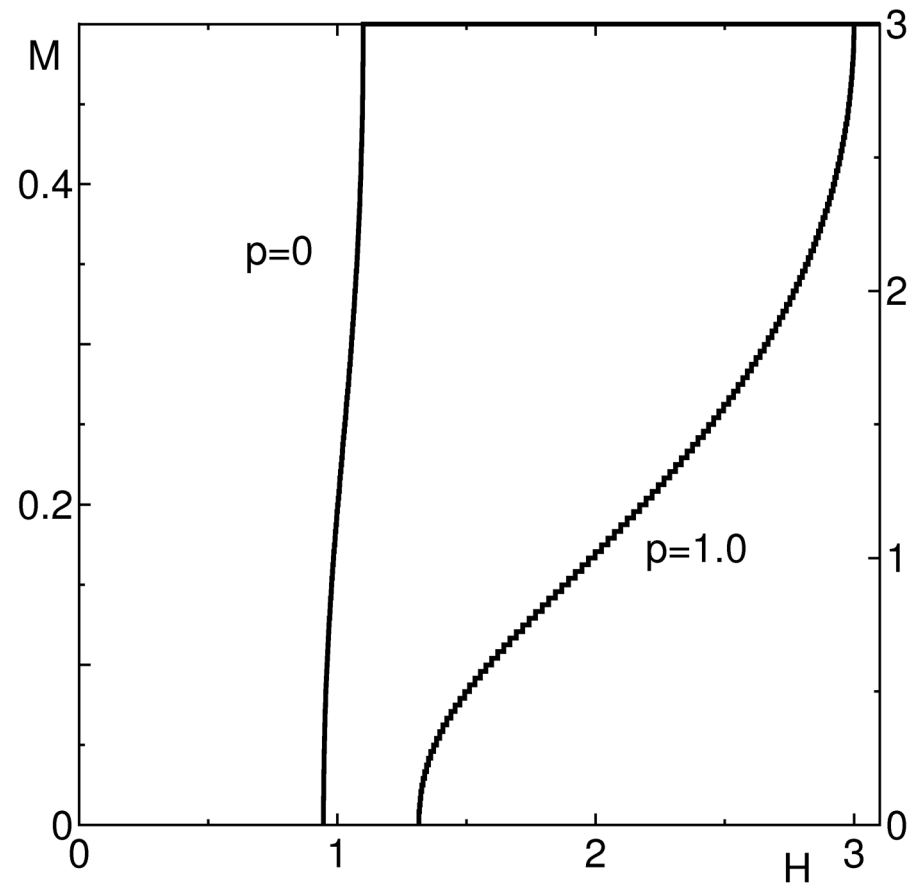
# [Bond Configuration]



### 3 Magnetization Curve of an Isolated Chain at $T = 0$ .

$$J_S = 2 \quad J_W = 0.1 \quad J = 1$$

#### 1. Uniform Chain - DMRG





## 4 Effect of Interchain Interaction

### Mean Field Approximation for the Interchain Interaction

$$\langle S_{i,j}^x \rangle = \begin{cases} (-1)^i m & J_{\text{int}} < 0 \text{ Interchain ferromagnetic interaction} \\ (-1)^i P_j m & J_{\text{int}} > 0 \text{ Interchain antiferromagnetic interaction} \end{cases}$$

for  $J_{\text{int}} > 0$

$$P_j = +1 \quad j \in \text{A-sublattice}$$

$$P_j = -1 \quad j \in \text{B-sublattice}$$

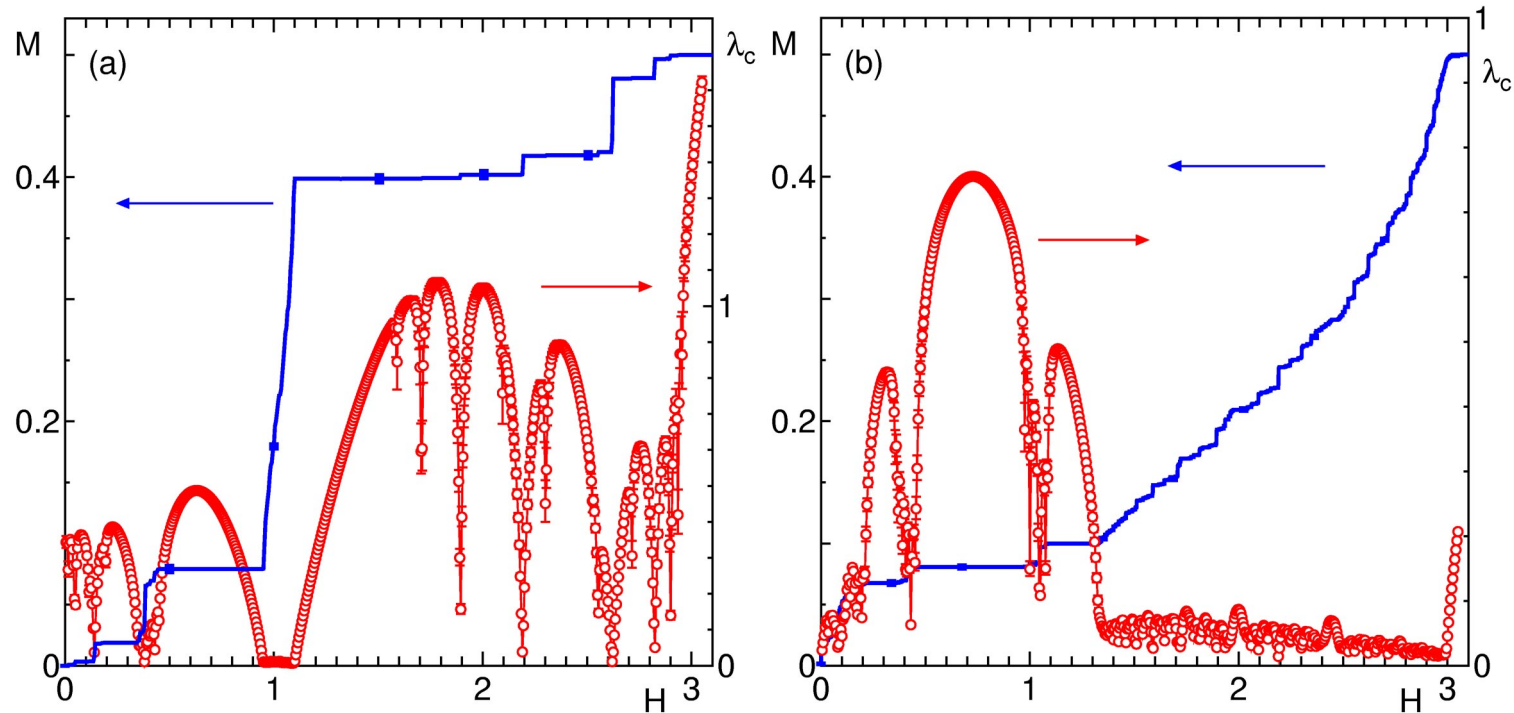
### Interchain mean field Hamiltonian

$$H^{\text{IMF}} = \sum_{i=1}^N J \mathbf{S}_{2i-1} \mathbf{S}_{2i} + \sum_{i=1}^N J_i \mathbf{S}_{2i} \mathbf{S}_{2i+1} - H_{\text{st}} \sum_{i=1}^N (-1)^i S_i^x$$

$$H_{\text{st}} = \lambda m(H_{\text{st}}) \quad \text{Self-consistent equation} \quad \lambda \equiv z |J_{\text{int}}|$$

$$\lambda_c = \lim_{H_{\text{st}} \rightarrow 0} H_{\text{st}} / m(H_{\text{st}})$$

$\lambda_c$ : critical interchain interaction  $\lambda$



$$p = 0.2$$

$$H_{\text{st}} = 0.0005$$

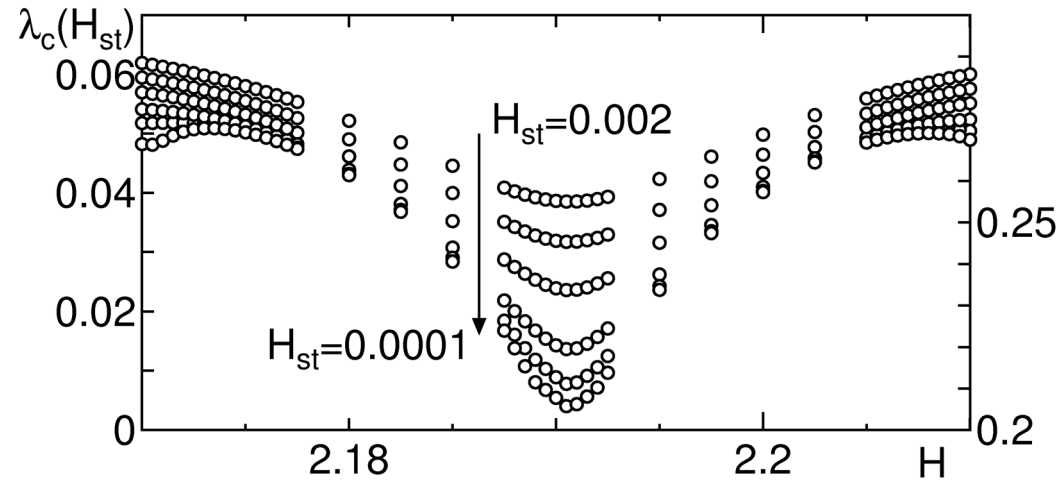
$$p = 0.8$$

Multiple reentrant transition

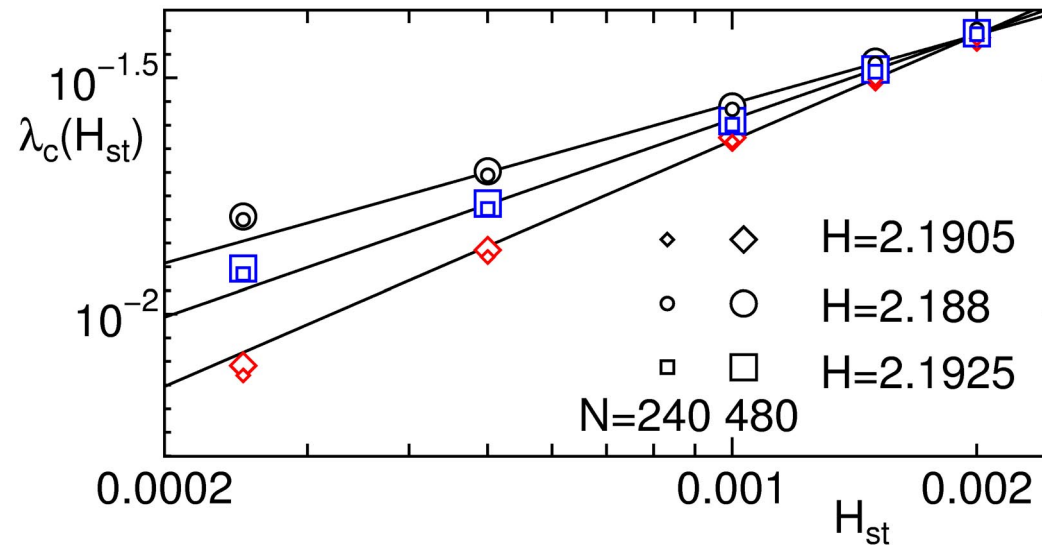
$N = 120$ . Averaged over 512 samples.



- Fine peak structure



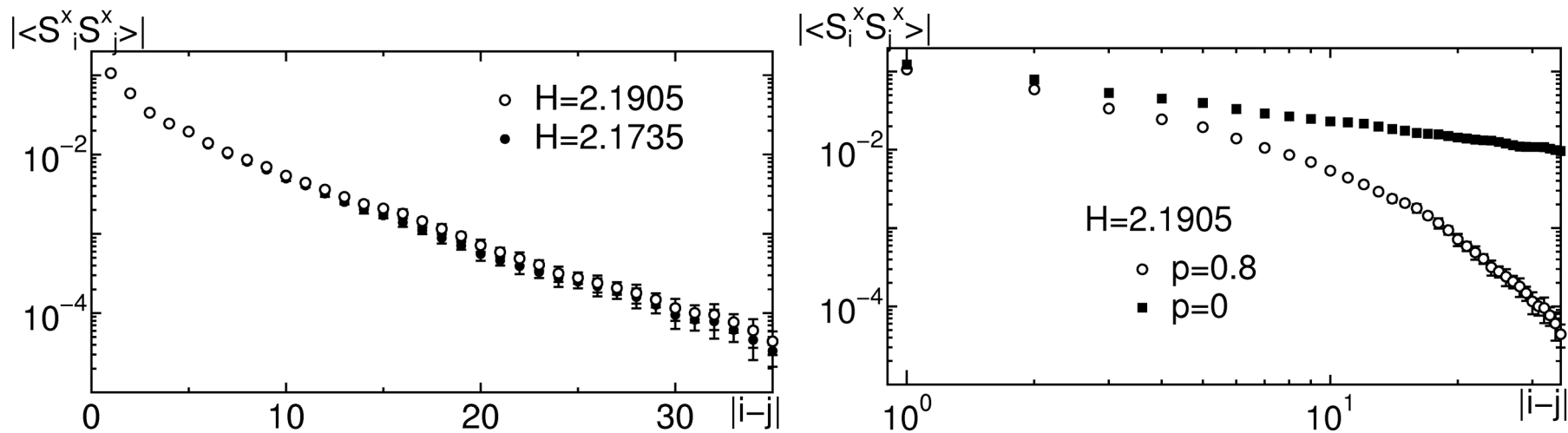
- $H_c$ -dependence of  $\lambda_c$



$N = 240,480$  Averaged over 256 samples (middle 240 sites).

$\lambda_c = 0$  only for discrete points?

## 5 Intrachain spin-spin correlation function



Averaged over 512 samples for  $N = 240$

**Exponential decay even for  $\lambda_c = 0$**

Non plateau state : Dense excited states near the ground state

Staggered mean field mixes the excited state into the ground state



Divergence of  $\chi_{st}$



Long range order with weak interchain interaction

## 6 Summary

1. Quasi-one-dimensional random alternating bond  $S = 1/2$  Heisenberg model exhibits multiple reentrant transitions in the magnetic field.

Method : DMRG+Interchain mean field approximation

2. In the absence of interchain coupling, the spin-spin correlation function decays exponentially even in the non-plateau regime.

Mix up the low energy excited states by interchain interaction.

⇒ Long range order

3. Bose glass phase(Nohadani *et al*) is not found

Limitation of interchain mean field approximation

4. ● Reentrant transition in 3-D random dimer system

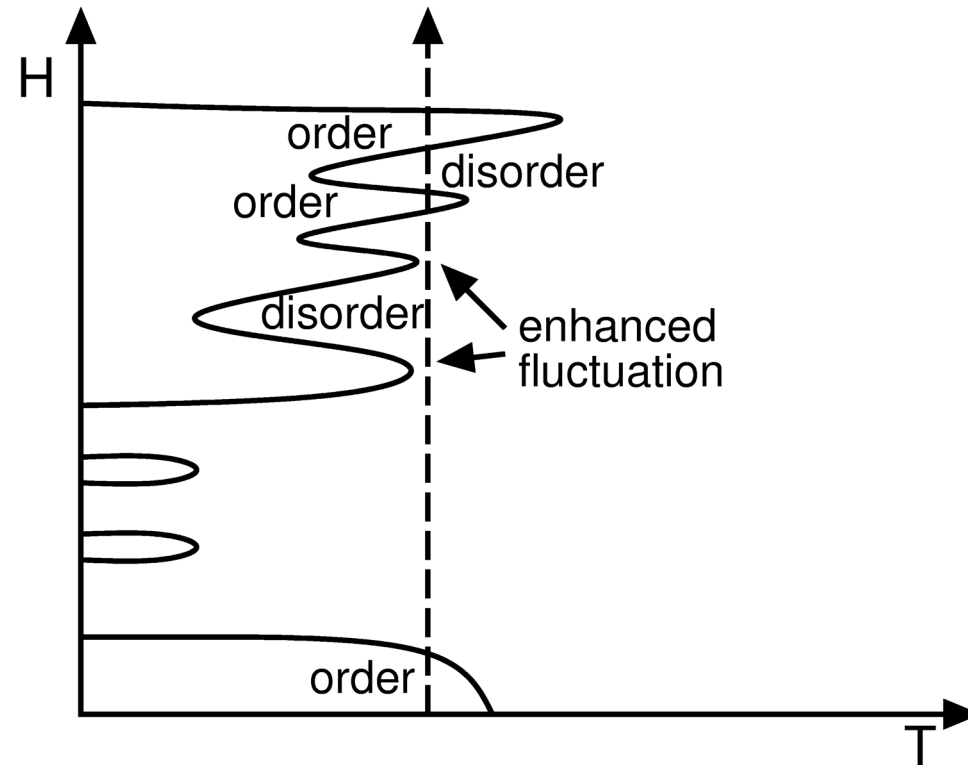
Random destruction of singlet dimer ⇒ Local moment

- Quasi-One-dimensional random alternating bond system

Competition of 2 types of dimer patterns ⇒ local moment

More complicated features = Multiple reentrant transition

## 5. Speculated finite temperature phase diagram



6. Possibility of experimental observation : random substitution of anions on the superexchange path

Ref. K. Hida : J. Phys. Soc. Jpn. 75, 074709 (2006). cond-mat/0602016;